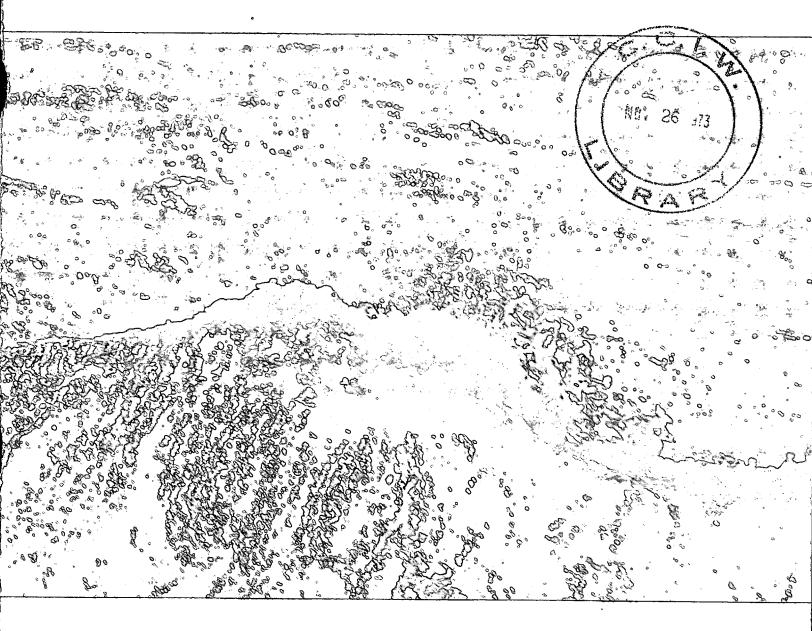
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Environment Environnement Canada Canada The Effect of Impurities on the Mechanical Properties of Ice Single Crystals

T. Nekamura and S. J. Jones



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INLAND WATERS DIRECTORATE, WATER RESOURCES BRANCH, OTTAWA, CANADA, 1973.



Environnement Canada

The Effect of Impurities on the Mechanical Properties of Ice Single Crystals

T. Nakamura and S. J. Jones

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Abstract

Constant strain-rate tensile tests were carried out on single crystals of ice doped with various impurities such as HF, HCI, HBr, NH₄OH, NH₄F, NaF, KF, NaOH, H₂O₂ and He. The tests were carried out in the temperature range of -5°C to -53°C. HF-doped crystals showed a softening effect in this temperature range as well as in the lower temperature range previously reported. HCI-doped crystals showed a less pronounced softening effect than HF. $H_2 O_2$ doped crystals showed a softening effect at -26°C, one He-doped crystal showed a hardening effect at -15° C. and the other impurities had no apparent effect. An activation energy for pure ice of 18±2 kcal/mole was deduced; for HCI-doped ice an activation energy of 12±3 kcal/mole and 13±2 kcal/mole was found using two methods of analysis. A power law relationship between stress and strain-rate was found to hold for pure and HCI-doped crystals with $n = 1.73\pm0.07$. The distribution of fluorine and chlorine ions in the crystals showed marked increases in ion concentration towards the surface.

Effect of Impurities on the Mechanical Properties of Ice Single Crystals

T. Nakamura and S. J. Jones

INTRODUCTION

This work was done between 1968 and 1970 when one of the authors (Nakamura) was on leave from the Institute of Snow and Ice Studies, National Research Center for Disaster Prevention, Nagaoka, Japan and held a National Research Council of Canada Post Doctorate Fellowship with the Inland Waters Directorate. This report contains experimental details, all the basic data obtained in the experiments, and the resultant conclusions. Preliminary results have been published (Nakamura and Jones, 1970). In addition, a paper summarizing the conclusions but of necessity omitting much of the data, was presented at a recent conference (Nakamura and Jones, 1973, in press).

The effect of HF, NH₄ OH and NH₄ F on the mechanical properties of ice single crystals has been reported by Jones and Glen (1969a). These impurities were chosen due to their effect on the electrical properties of ice; they were thought to enter the ice lattice substitutionally (Steinemann, 1957; Gränicher, 1963; Glen, 1969). It was found that at temperatures around -60° C, HF had a softening effect in that the creep rate was greater for HF-doped ice than for pure ice, while NH₄ OH had a slight hardening effect and NH₄ F had no observable effect. The object of the present work was to investigate the effect of HF, NH₃ and NH₄ F at higher temperatures than those used by Jones and Glen (1969a) and to ascertain if any other impurities, which may or may not dissolve substitutionally in ice, had a similar effect.

Ida and others (1966) have reported experiments on the dielectric dispersion of impure ice below -100° C. They found such dispersion only in KOH, NaOH, LiOH, H₂O₂, NH₄ F or HF-doped ice, and not in Ba(OH)₂, NH₄OH, HCI or KCI-doped ice. Some of these impurities were chosen to determine if there was a corresponding grouping in terms of mechanical properties.

From dielectric measurements, Young and Salomon (1968) concluded that HCl enters the ice lattice substitutionally, at least at small concentrations. This was another reason for investigating the mechanical properties of HCldoped ice. The majority of the results reported here are concerned with HCl-doped ice.

While this work was in progress, Kahane and others (1969) showed that the inert gases He and Ne could be dissolved in ice, presumably as an interstitial atom. We have attempted, therefore, to determine the effect helium might have on ice's mechanical behaviour. Results are given in the temperature range -5° C to -46° C for the following impurities: HF, NH₄OH, NH₄F, HCl, HBr, NaF, KF, NaOH, H₂O₂ and He.

EXPERIMENTAL

Preparation of Crystals

The single crystals were grown in a manner similar to that described previously by Glen & Jones (1967) except that a small immersion heater was inserted in the bowl to prevent the ice from freezing over completely and thus causing a pressure to build up during freezing. A seed crystal was also placed in the water to prevent supercooling which otherwise would sometimes occur. The water used was de-ionized water of greater than $10M\Omega$ -cm resistivity, de-aired by boiling and cooled under vacuum. Except for Helium (the procedure for He is described in the next section), the impurity was added to the water just before growth of the water crystals started. In order to obtain ice crystals with the angle between the tensile and optic axes close to 45°, all glass tubes in the bowl were placed at an angle of 45° to the vertical. After growth of the crystals, the optic axis angle was checked by a polarized light method. Almost all the crystals had c-axis orientations within 6° of 45° to the tensile axis.

Procedure

A crystal was mounted in a metal grip and frozen in place. The grip and crystal were then placed in the upper jaw of an Instron table model mechanical tester, and the crystal frozen into the lower metal grip. All this was done in a cold room at -10° C. A polystyrene box was made to fit around the crystal thus reducing the temperature fluctuations from the cold room. The lowest temperature obtainable using the cold room alone was about -25° C; to obtain lower temperatures, methanol, cooled in a separate cold bath (model Tenney Engineering Inc., MR-9), was passed through coils around the crystal. In this case the temperature was controlled with a platinum resistance thermometer and controller to an accuracy of $\pm 0.05^{\circ}$ C. The lowest temperature obtained was -53° C.

All the tests reported here were conducted at a constant velocity of deformation. Since, however, an Instron is a 'hard' machine and since the total strain of a sample was always small – typically less than 5% – this is equivalent to a constant strain-rate.

The concentration of the various ions in the crystals was usually determined by melting the sample after a test and analyzing the water with a pH/mV meter (Orion model 801) and various specific ion electrodes. In the case of NH₄ OH and NH₄ F, a colorimetric method was used with Nessler's reagent since no specific ion electrode existed for NH₃. In the case of Helium a different method altogether was used.

The He-doped crystals were prepared by placing a pure ice crystal in a vessel containing Helium gas at a pressure of 125 bars. The crystal was left for at least 24 hours, and then removed and mechanically tested at atmospheric pressure. After the test, the crystal was melted and the amount of Helium gas given off was measured with a microburet system.

RESULTS

General Results

All the crystals used for the tensile tests were good single crystals as determined optically. However, from X-ray topographic studies, it was found that sub-grain boundaries were often present. The sub-grains were typically a few mm in diameter and the boundaries represented misorientations of approximately one minute of an arc. The dislocation density within these sub-grains was at least 10^5 cm⁻².

Table 1 shows the concentration of impurities in the mother solution and in the ice crystals formed from the solutions. Relatively large amounts of HF, HCl and NH₄F would enter the ice and it was usually easy to grow single crystals with these impurities. NaF and KF-doped single crystals were very difficult to grow and very little of the impurity was found in the ice. NH₄ OH and NaOH-doped crystals were quite easy to grow, but again very little of the impurity had entered the ice. The distribution of the impurity in the ice is discussed in a later section.

The crystals exhibited slip on the basal (0001) plane as shown by one example of slip lines on an NaOH-doped

Table 1. The concentration of impurity in the ice from
given mother solutions

Impurity	Concentration of mother solution ppm	Concentration in ice ppm
HF	100	5
HCI	_° 200	2
HBr	60	0.4
NH4OH	100	0.1
NH4F	100	10
NaF	20	0.1
КF	~20	0.1
NaOH	100	~0.1
H ₂ O ₂	200	-

crystal in Figure 1. This is the normal mode of deformation in ice, as observed by previous workers (McConnel, 1891; Glen and Perutz, 1954; Higashi *et al.*, 1964; Glen and Jones, 1967).

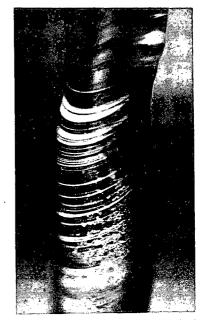


Figure 1. Basal Slip lines on an NaOH-doped single crystal of ice deformed 20% at -26°C.

Resolved shear stresses and resolved shear strains used in this work are defined by Schmid & Boas (1968, p. 60).

Unfortunately, the results obtained from a number of nominally pure crystals all tested at the same strain-rate and temperature were not entirely reproducible, as will be seen when studying the various figures. For example, Figure 8 shows results from three pure crystals tested at a strain-rate of $1.6 \times 10^{-6} \text{ s}^{-1}$ at -26° C. As can be seen, the curves vary in stress by a factor of 1.3. It is believed that this variation is due to differences in the initial number of dislocations present in the crystal.

HF-doped Crystals

Figure 2 shows typical stress-strain curves for pure and HF-doped crystals at -11° C for a number of different strain-rates as marked on the figure. The curves show an initial increase in stress up to a maximum or 'yield' stress followed by a continuous decrease in stress, as found by previous workers. No work hardening was observed. Figure 2 shows that the HF-doped crystals were somewhat softer than pure ice in that the stress level required to maintain the constant strain-rate was lower for the HF-doped crystals. This is particularly true of the maximum stress reached. If one plots the maximum resolved shear stress against the resolved strain-rate, one obtains such curves as are shown for -11° C in Figure 3. Figure 3 shows that, for a given stress, the strain-rate of the HF-doped ice is 6 times that of the pure ice. Similar results at -30° C (Fig. 4 - 7) indicate an increase in strain-rate of a factor of 7. Some additional results at -26° C are shown in Figure 8. There is a suggestion in Figures 3 and 7 that the effect of the HF is smaller at higher strain-rates but this may be an artifact due to the greater amount of data collected at the higher strain-rates. As has been mentioned there is some scatter in the data but as far as could be detected there was no systematic dependence on the amount of impurity.

HCI-doped Crystals

Preliminary results for HCI-doped ice have been published (Nakamura and Jones, 1970). Figures 9-12 show the stress-strain curves obtained at -26°C for pure and HCI-doped ice at various strain-rates and Figure 13 shows the maximum shear stress plotted as a function of strain-rate. Least squares straight lines have been fitted to the pure data and the HCI-doped data on Figure 13. The HCI-doped crystals have a strain-rate that is about a factor of 3 greater than the pure ice. Figures 14 and 15 show the data collected at -46°C; Figures 16, 17 and 18 show the data collected at -15°C and Figures 19, 20 and 21 show the data collected at -5.6°C. Figure 18 contains also two points which refer to He-doped ice; these will be discussed in the He section on doped crystals. Comparing Figures 13, 15, 18 and 21 there seems to be a significant softening at -46°C and -26°C but very little, if any, at -15°C and -5.6°C. This suggests that the activation energy of the HCI-doped ice is less than that for pure ice. A direct determination of this activation energy also gives a value (12±3 kcal/mole) less than for pure ice (18±2 kcal/mole) as shown in the section headed 'Activation Energy of Pure and HCI-doped Crystals".

H_2O_2 -doped Crystals

Results for 2 crystals of H_2O_2 -doped ice are shown in Figure 22, for a temperature of -26°C. Although the data is limited there appears to be a softening effect since the stress level for the doped ice is about a factor of three less than for the pure crystals.

NH₄OH-doped Crystals

Jones and Glen (1969a) showed that ammoniumdoped ice crystals had a slight hardening effect at about -60° C. This effect was examined at higher temperatures and the results at -11° C are shown in Figures 23 and 24. A slight softening is apparent although there is some scatter in the data. At -26° C there was no effect of the NH₄OH (Figure 25).

HBr-doped Crystals

While it was not expected that HBr would enter the ice lattice substitutionally because of the radius of the Br atom (1.96 Å) as compared to that of the oxygen atom (1.38 Å), two specimens were tested at -26° C and it was observed that there was no significant difference between these crystals and pure ice (see Fig. 26).

NH₄ F-doped Crystals

Only two crystals were tested, one at -11° C and one at -26° C, and in both cases there was no significant difference from pure ice. This is in agreement with the data of Jones and Glen (1969a) at lower temperatures.

NaF-doped Crystals

Three crystals were tested at -26° C and no significant effect of the NaF was observed (Fig. 27).

KF-doped Crystals

Only one crystal was tested at -33° C as shown in Figure 28, and no effect was observed.

NaOH-doped Crystals

Five crystals were tested at -26° C and the results are shown in Figure 29. No difference was observed between these and pure crystals (Fig. 30) but only small amounts of NaOH entered the ice samples (≤ 0.1 ppm) and at this level our method of determining the concentration was not reliable.

He-doped Crystals

The results from only two crystals are shown in Figure 31 and the maximum stress values are plotted on Figure 18. Apparently the crystal containing \approx 400 ppm. Helium was harder than pure ice, but a second crystal prepared in the same manner contained zero Helium when it was analyzed. This was because the Helium was diffusing out of the specimen in a shorter time than it took to do the mechanical test. The diffusion coefficient of Helium in ice has recently been measured by two groups of workers. Davy and Miller (1970) have reported a value of 10^{-6} cm² s⁻¹ at -10° C, and Haas and others (1971) give D $\approx 10^{-5}$ at -10° C with a strong anisotropy.

Dependence of Strain-rate on Stress

Previous workers (Higashi, 1969; Readey & Kingery, 1964; Jones and Glen, 1969b) have found that for pure ice the strain-rate, $\dot{\epsilon}$ can be expressed as a function of stress, τ , as $\dot{\epsilon} \propto \tau^n$ where n is a constant or a function of stress (Jones and Glen, 1969b).

There are two ways in which this relationship can be investigated; different samples can be tested at different strain-rates and the stresses recorded, or one sample can be stressed at different strain-rates. The first method leads to such curves as Figure 3 etc. from which n can be determined from the slope of the lines. The second method avoids the lack of reproducibility due to different crystals; with this method a crystal is strained at one strain-rate and then the strain-rate is changed to a higher (or lower) value. The two values of stress, immediately before and after the change in strain-rate, are noted and then a curve of stress against strain-rate is plotted. Such a curve is shown in Figure 32 and from the slope of the straight lines the value of n is determined. Table 2 summarizes all the n values we have obtained using both methods of determination for various pure and doped samples. A mean $n = 1.73 \pm 0.07$ is obtained for all the crystals.

Table 2.	Values of n determined for pure and
	doped ice single crystals

Temperature	n	Mean n
°c		
- 16	1.7	
	+	
•	1.6	
	1.5	
-15	1.7	
-26	1.4	
- 30	2.0	
-46	2.5	
		1.73 ±0.07
-6	1.9	
-53	1.6	
-26		
-6		
-15	1.6	
-26		
-11		
-46	1.8	
_5	22	
	$^{\circ}C$ -16 -12 -6 -11 -15 -26 -30 -46 -6 -53 -26 -6 -15 -26 -11	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$

Activation Energy of Pure and HCI-doped Crystals

It is generally accepted that the dependence of creep deformation on temperature is given by an Arrhenius equation:

$$\dot{\epsilon} \propto \tau^{n} \exp \left(-Q/RT\right)$$
 (2)

where \dot{e} is strain-rate, τ is stress, n is a constant, Q is an activation energy, (strictly an activation enthalpy) R is the gas constant and T is the absolute temperature. This equation can be re-written as:

$$\tau \propto \dot{\epsilon}^{1/n} \exp (+Q/nRT)$$
 (3)

so that for a constant strain-rate the dependence of τ on T will give Q/n, (Higashi, 1969). If n is known from other tests as in the section "Dependence of Strain-rate on Stress" then Q can be determined. We attempted to measure Q in this way for HCI-doped ice. The tests were done by taking one specimen and straining it at a constant rate and temperature, and then changing the temperature and observing the new stress required to maintain the same strain-rate. The results are shown in

Figures 33 to 36. The time between the end of one test and the beginning of the new one at the new temperature varied from 1-10 hours. The stress immediately before the temperature change is compared with the stress immediately afterwards and these results are shown plotted against 1/T in Figure 37. Using the value of 1.73 for n given in the previous section, values of Q were obtained (see Table 3). There is quite a scatter in the results but a mean activation energy is 12 ± 3 kcal/mole for HCI-doped ice.

A second method, which avoided the need to know n, was also used to determine the activation energy of both pure and HCI-doped ice. From the data given in Figs. 1-21 the strain-rate required to give a certain maximum resolved shear stress at different temperatures was determined and plotted as shown in Figures 38 and 39. Figure 38 shows data for HCI-doped ice at three different stresses and Figure 39 shows data for pure ice at two different stresses. The slope of the lines is proportional to the activation energy and is found to be 13±2 kcal mole⁻¹ for HCI-doped ice and 18±2 kcal mole⁻¹ for pure ice. The two values of activation energy for HCI-doped ice obtained by different methods are in good agreement. The value for pure ice of 18 kcal mole⁻¹ is based on less data than the HCI-doped result and is rather high compared to previous workers. Nevertheless, using these two values of activation energy we have reduced all our stress/strain-rate data for pure and HCI-doped ice to one temperature, -26°C. This temperature was chosen because a lot of our data was obtained at that temperature. Pure ice data was reduced using an activation energy of 18 kcal mole⁻¹ and HCI-doped ice data was reduced using an activation energy of 13 kcal mole⁻¹. The result is shown in Figure 40. The line drawn on the graph indicates a demarcation between the pure and HCI data. To the left of the line lie 92% of the HCI-doped data points and to the right of the line 75% of the pure points. This shows a softening effect of HCI despite the scatter in the data.

Distribution of lons in Single Crystals

The distribution of HF and HCl ions in the doped crystals was examined. A doped crystal was allowed to melt in successive stages and two or three water samples were collected. Figures 41 and 42 show HF results while

Table 3.	Activation	energy	of	HCl-do	ped	ice	crystals

Sample	Concentration HCl ppm	Mean activation energy of sample kcal. mole ⁻¹
HC1-A	2.0	8.9
HCI-C	2.0	8.7
HC1-D1	2.5	12.5
HC1-D ₂	2,5	16.6

Mean activation energy = 12 ± 3 kcal. mole⁻¹.

Figures 43-45 give the HCl results. All show higher ion concentration at the surface than at the centre of the specimen. There were two variations of the experiment, wiping a thin layer off the crystal with a tissue (Figs. 41, 43, 45) or not wiping as the case may be (Figs. 42, 44). For the tensile tests the specimens were wiped with a tissue before testing and so the distribution of the ions in the specimen presumably corresponded to Figures 42, 44. As can be seen from these figures the distribution of HF was much more uniform than of HCI - reflecting the higher diffusion coefficient of HF (Krishnan and Salomon, 1969; Kopp and others, 1965; Haltenorth and Klinger 1969). The reason for such a concentration curve is thought to be that as the ice grows the ions are rejected to the melt preferentially at the centre of the crystal, but become trapped at the ice-glass interface.

CONCLUSIONS AND DISCUSSION

HF had a softening effect at high temperatures (-11°C) as well as low temperatures (-70°C); HCl had a smaller softening effect; HBr, NH₄F, NaF, KF, NaOH had no observable effect; NH₄OH showed a possible slight softening; H₂O₂ showed a softening effect. One Helium specimen showed a hardening effect.

The HF result was similar to that observed by Jones and Glen (1969a) and therefore can be explained by the same theory developed by Glen (1968) concerning the interaction of dislocations and point defects in ice. He has pointed out that due to the random positions of the hydrogen atoms in the ice lattice, a moving dislocation would create point defects of either the ionic or Bjerrum sort (Bjerrum 1951, 1952); he has calculated that this would require a stress of 4300 bar, a value far in excess of the stress needed (about 1 bar) to deform ice plastically. From this theory the result is a dislocation line held up at a number of points along its length until such time as one of the naturally occurring point defects passes and reorients the bond, thus allowing the dislocation line to move forward. When HF is dissolved in ice it is thought to enter the lattice substitutionally and by doing so it introduces Bjerrum and ionic defects (Gränicher 1963; Glen 1969). Kröger (1964) has shown that the number of L type Bjerrum defects increases with rising HF content and in the limit of high concentrations becomes equal to the HF concentration. Since the number of point defects has been greatly increased by the HF, the probability of reorientation of a bond which is holding up a dislocation is greatly increased. The dislocation will move faster and so the ice will be softer. Recently Jones and Gilra (1972) have shown that the diffusion of HF into ice causes a marked increase in dislocation density, which can account for the mechanical softness of HF-doped ice. This may be a result of dislocation multiplication brought about by the increased number of point defects introduced by the HF. If HCI goes into the ice lattice substitutionally as suggested by Young and Salomon (1968) and more recently by Seidensticker (1972), then it will presumably behave like HF and hence the softening effect of HCI can be explained in the same

way as HF. There is no evidence that HBr, NaF or KF enter the ice lattice substitutionally, although it might be expected that the F⁻ion could enter the lattice as it does with HF. The fact that we observed no effect of NaF or KF suggests that the F⁻ion does not in fact substitute for the O atom in these cases. Kelly and Salomon (1969) believe from dielectric measurements of polycrystalline ice that NaOH does not enter the lattice substitutionally either. The slight softening of NH₄OH doped specimens at -11°C is in apparent contradiction to Jones and Glen (1969a) who observed a hardening effect at -60° C. Our results, however, show a large scatter, and so perhaps it is not reasonable to propose a theory to explain this contradiction until the experimental evidence is clear. Nevertheless this apparent contradiction could be explained by postulating that the many D defects introduced by the NH₃ are mobile at the high temperature but immobile at the lower temperature. This implies that NH₄OH-doped ice would have a higher activation energy for creep than pure ice.

Foley and Giguère (1951) studied the phase equilibrium of the system $H_2O_2-H_2O$ and concluded that solid solutions are not formed. We had no means of testing for the presence of H_2O_2 in our samples but from comparison with other impurities we would expect there to be only a few ppm present. Foley & Giguère (1951) would not have detected such small quantities in their experiment. If H_2O_2 can enter the ice substitutionally in these small quantities it might introduce L-defects and so behave like HF. This would then account for the apparent softening.

Helium, being an inert gas, must enter the ice lattice as an interstitial atom. As such, it could act as a pinning point for dislocation and hence slow them down. This would give an apparent hardening.

Our value of n = 1.73 is close to the value, in the same temperature range, observed by previous workers. Higashi and others (1964) found n = 1.53 in tensile, constant strain-rate tests and 1.58 in bending creep tests (Higashi and others, 1965). Readey and Kingery (1964) obtained n = 2.0 and Muguruma (1969) found n = 1.3 and n = 1.7 for chemically polished and mechanically polished specimens respectively. Jones and Glen (1969b) found n = 2.3 at -70° C indicating that n increases with decreasing temperature.

The value of activation energy for pure ice obtained in this work, 18 ± 2 kcal/mole, is somewhat higher than that obtained by previous workers: Higashi and others (1964) obtained 15.8 kcal/mole in bending creep experiments, Readey and Kingery (1964) obtained 14.3 kcal/mole, Jones and Glen (1969) obtained 15.6 kcal/mole in tensile creep experiments above -50° C and 9.6 ± 0.5 kcal/mole below -50° C. Our value for HCI-doped ice is about 3 kcal/mole less than that obtained for pure ice by these workers.

ACKNOWLEDGMENTS

We are grateful to G.A. Barnett for technical assistance in the laboratory.

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Appendix

(Figures 2 to 45)

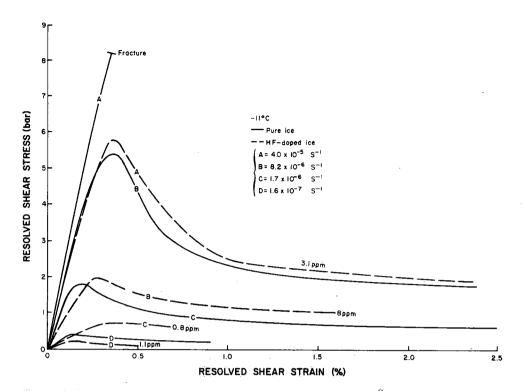
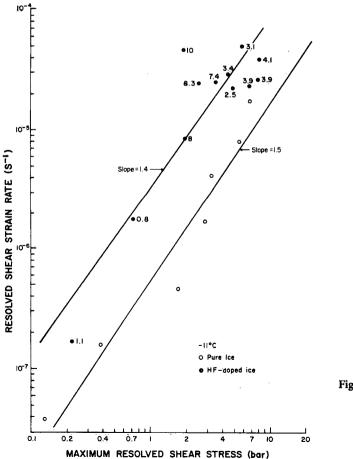
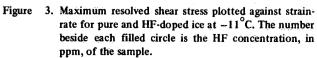


Figure 2. Stress-strain curves for pure and HF-doped ice deformed at $-11^{\circ}C$ and various strain-rates and dopant concentrations as marked on the figure.





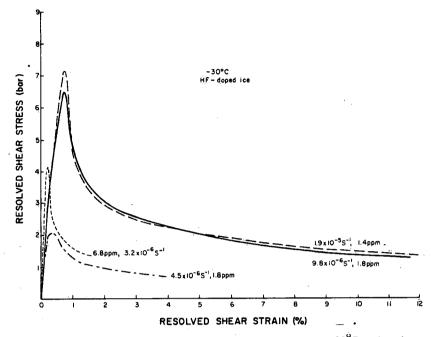


Figure 4. Stress-strain curves for HF-doped ice deformed at -30° C and various strain-rates as marked on the figure.

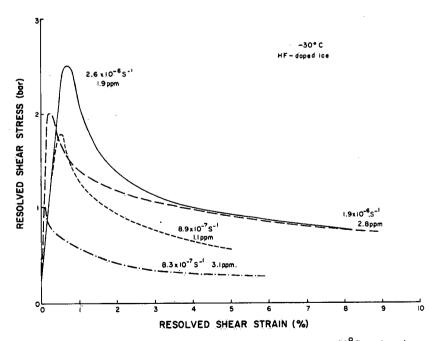
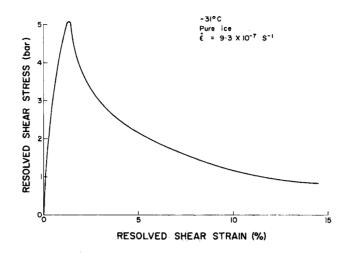
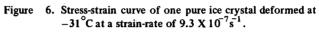


Figure 5. Stress-strain curves for HF-doped ice deformed at -30° C and various strain-rates as marked on the figure.





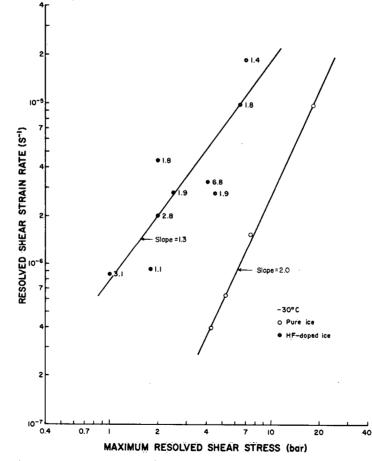


Figure 7. Maximum resolved shear stress plotted against strain-rate for pure and HF-doped ice deformed at -30° C. The number beside each filled circle is the HF concentration of the sample, in ppm.

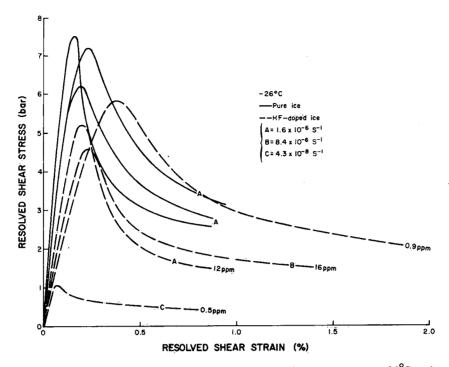


Figure 8. Stress-strain curves for pure and HF-doped ice deformed at $-26^{\circ}C$ and various strain-rates and dopant concentrations as marked on the figure.

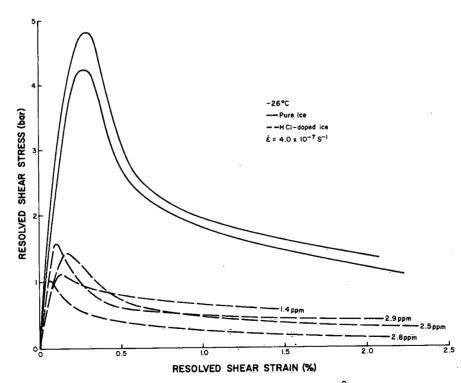
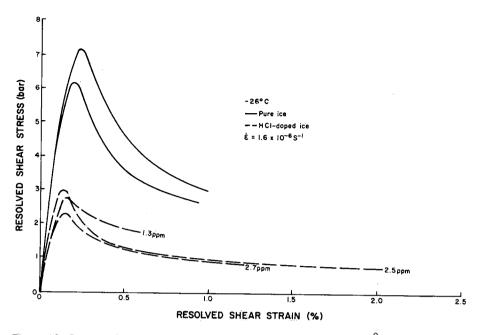
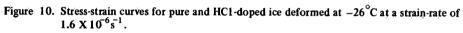


Figure 9. Stress-strain curves for pure and HC1-doped ice deformed at -26° C at a strain-rate of 4 X $10^{-7}s^{-1}$.





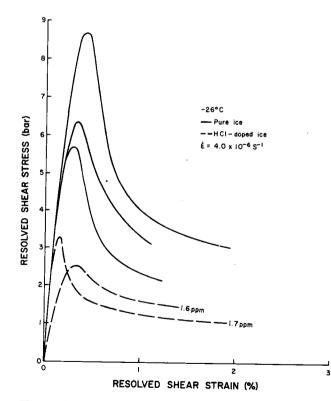


Figure 11. Stress-strain curves for pure and HC1-doped ice deformed at -26° C at a strain-rate of 4 X $10^{-6}s^{-1}$.

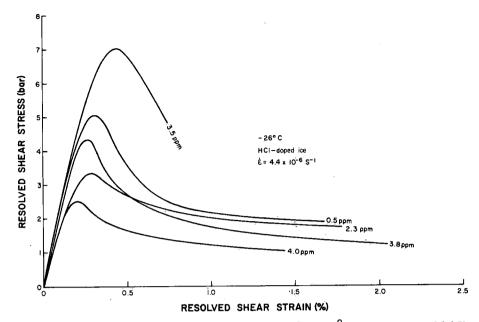


Figure 12. Stress-strain curve for HC1-doped ice deformed at -26° C at a strain-rate of 4.4 X $10^{-6}s^{-1}$.

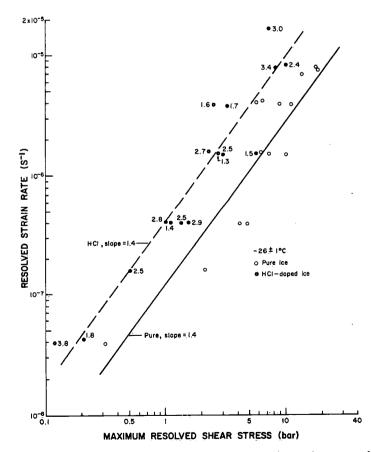


Figure 13. Maximum resolved shear stress plotted against strain-rate for pure and HC1-doped ice deformed at -26°C. The number beside each filled circle is the HC1 concentration of the sample in ppm.

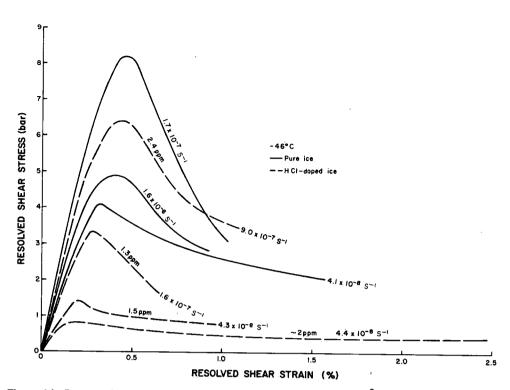
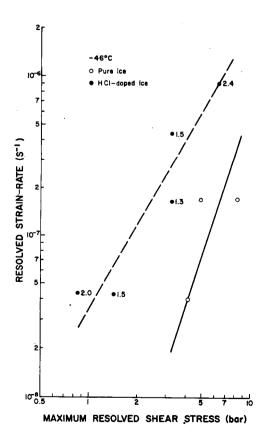
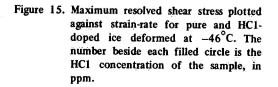


Figure 14. Stress-strain curves of pure and HC1-doped ice deformed at -46° C at various strain-rates as marked on the figure.





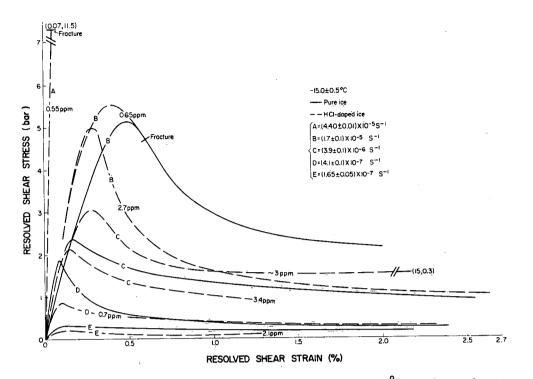
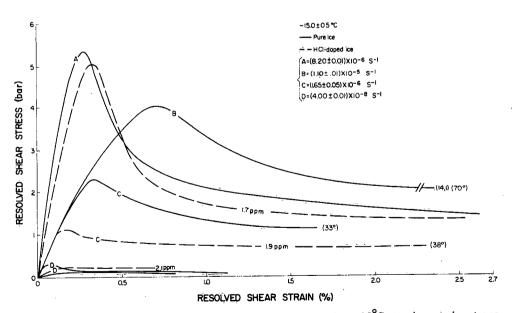


Figure 16. Stress-strain curves for pure and HC1-doped ice deformed at -15° C at various strain-rates and dopant concentrations as marked on the figure.



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Figure 17. Stress-strain curves for pure and HC1-doped ice deformed at -15°C at various strain-rates as marked on the figure.

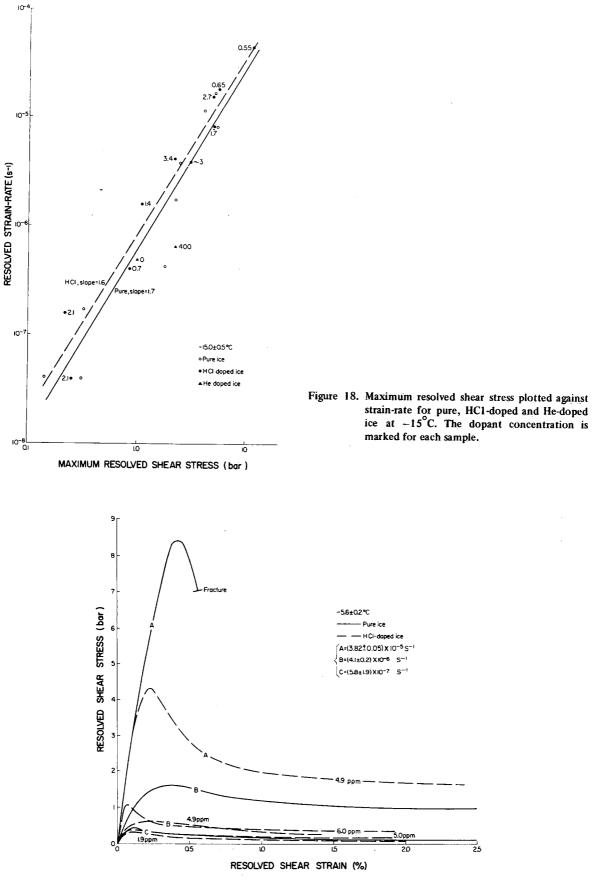


Figure 19. Stress-strain curves for pure and HC1-doped ice deformed at -5.6° C at various strain-rates as marked.

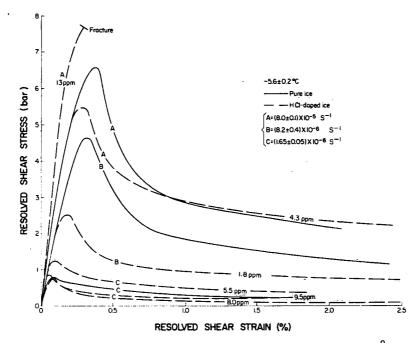


Figure 20. Stress-strain curves for pure and HC1-doped ice deformed at -5.6°C at various strain-rates as marked.

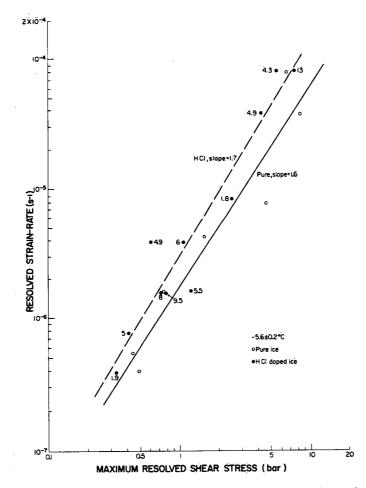
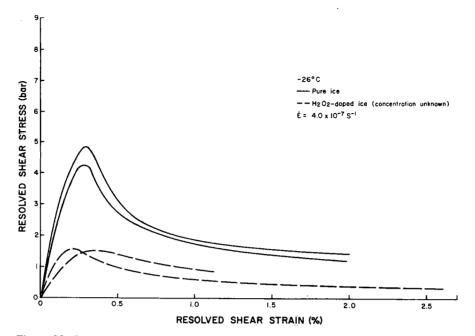


Figure 21. Maximum resolved shear stress plotted against strain-rate for pure and HC1-doped ice at -5.6°C. The number beside each filled circle is the HC1 concentration in the sample, in ppm.



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Figure 22. Stress-strain curves for two H_2O_2 -doped samples and two pure ice samples deformed at $-26^{\circ}C$ at a strain-rate of $4 \times 10^{-7} s^{-1}$.

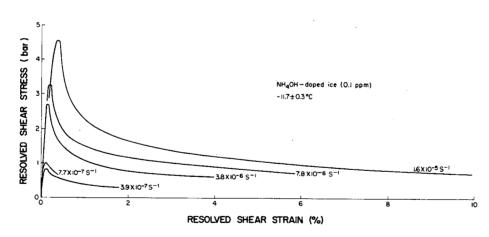


Figure 23. Stress-strain curves for NH₄OH-doped ice deformed at -11.7°C at various strain-rates as marked.

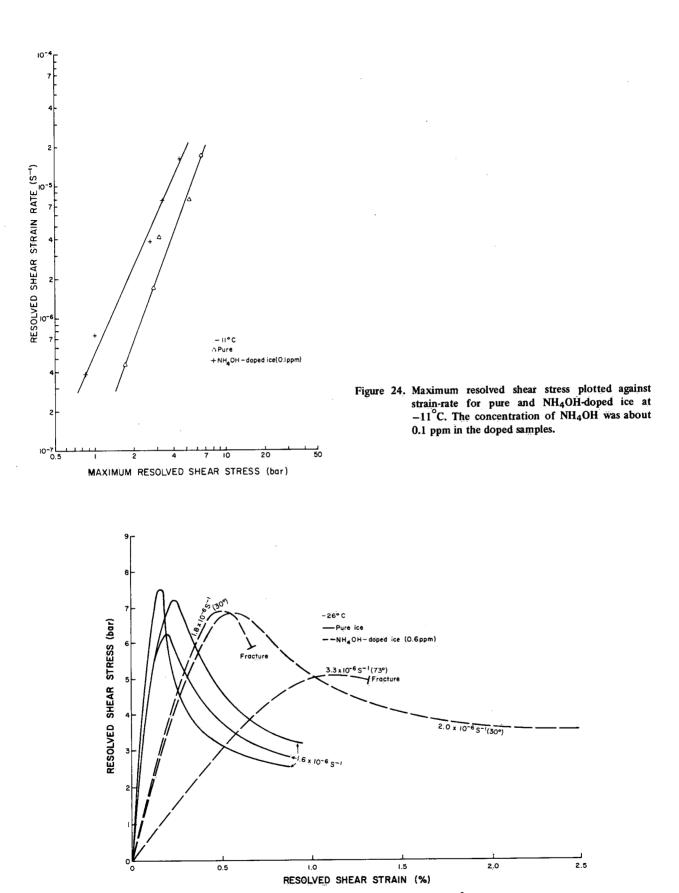


Figure 25. Stress-strain curves for NH₄OH and pure ice deformed at -26°C at various strain-rates as marked.

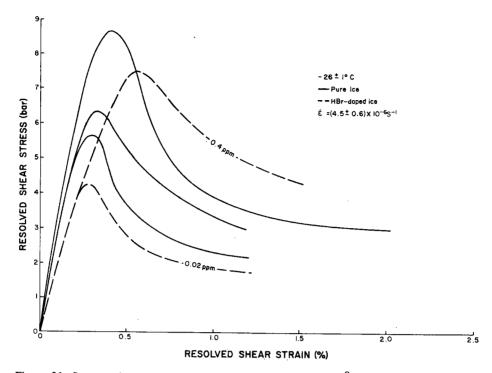


Figure 26. Stress-strain curves for HBr-doped ice deformed at $-26^{\circ}C$ at a strain-rate of 4.5 X $10^{\circ}s^{-1}$.

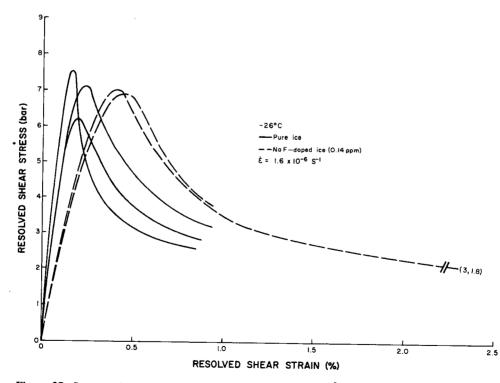


Figure 27. Stress-strain curves for pure and NaF-doped ice at -26° C at a strain-rate of 1.6 X 10^{-6} s⁻¹.

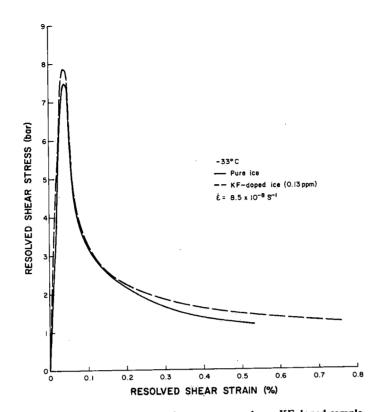


Figure 28. Stress-strain curves for one pure and one KF-doped sample deformed at -33° C and at a strain-rate of 8.5 X $10^{-8}s^{-1}$.

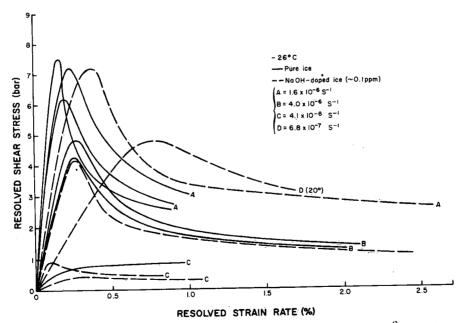
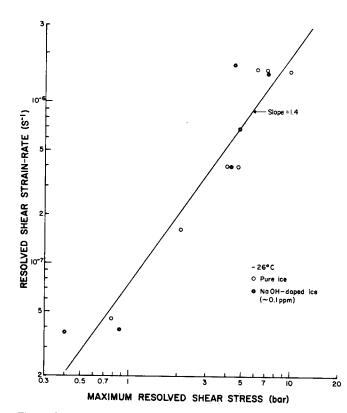


Figure 29. Stress-strain curves for pure and NaOH-doped ice deformed at -26°C at various strain-rates as marked on the figure. The curve marked D(20°) was a sample with an angle of 20° between its optic and tensile axis.



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Figure 30. Maximum resolved shear stress plotted against strain-rate for pure and NaOH-doped ice at -26° C.

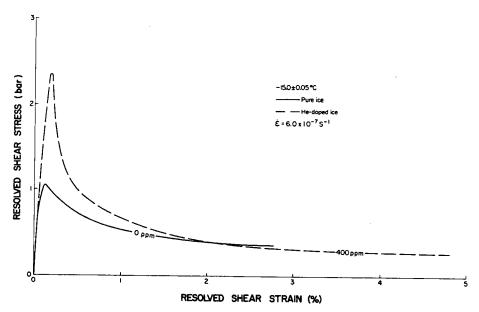
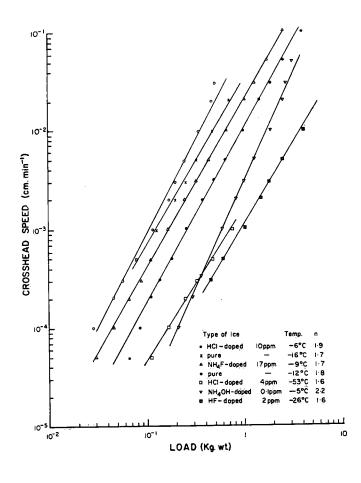
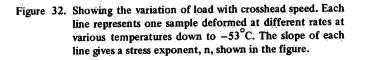


Figure 31. Stress-strain curve for two He-doped ice crystals deformed at -15°C. After testing, one sample had no He left in it.





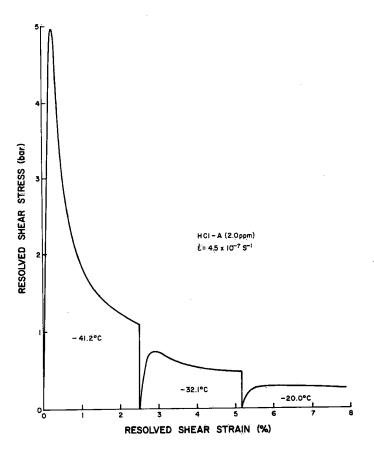


Figure 33. Stress-strain curve for one HC1-doped specimen deformed at different temperatures at a strain-rate of 4.5 X 10⁻⁷s⁻¹

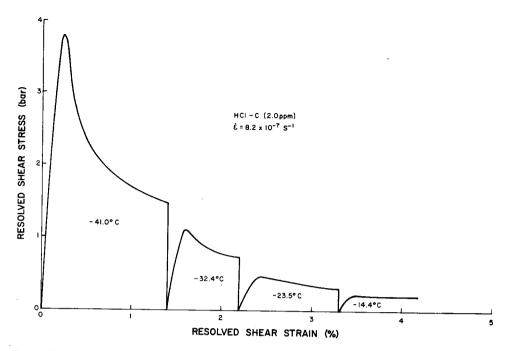


Figure 34. Stress-strain curve for one HC1-doped ice crystal deformed at different temperatures at a strain-rate of 8.2 X $10^{-7}s^{-1}$.

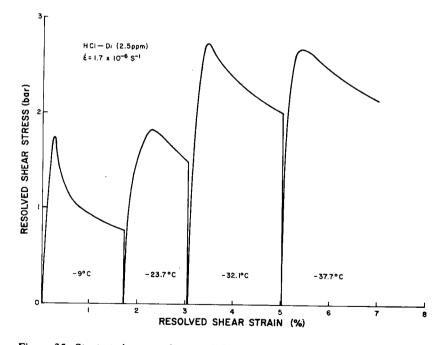


Figure 35. Stress-strain curve for one HC1-doped specimen deformed at different temperatures at a strain-rate of $1.7 \times 10^{-6} s^{-1}$.

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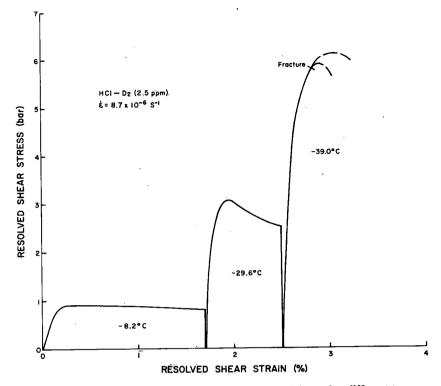
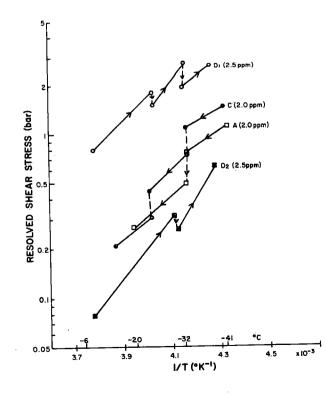
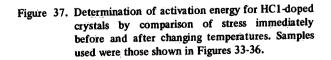
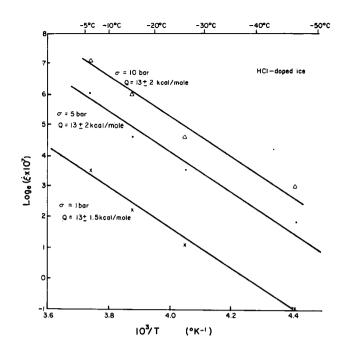
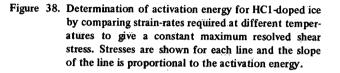


Figure 36. Stress-strain curve for HC1-doped specimen deformed at different temperatures at a strain-rate of 8.7 X $10^{-6}s^{-1}$. The sample was the same one as used for Figure 35.









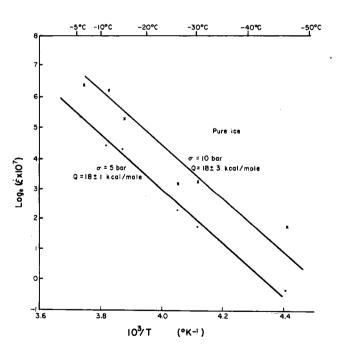


Figure 39. Determination of activation energy for pure ice by comparing strain-rates required at different temperatures to give a constant maximum resolved shear stress. Stresses are shown for each line and the slope of the line is proportional to the activation energy.

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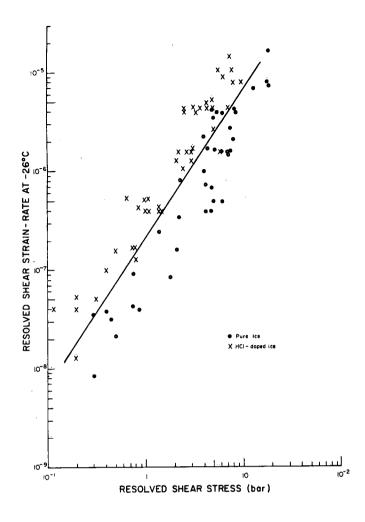


Figure 40. Maximum resolved shear stress for all pure and HC1doped samples plotted against strain-rate reduced to -26° C by using an activation energy of 18 kcal/mole for pure ice and 13 kcal/mole for HC1-doped ice.

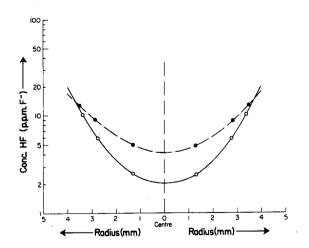


Figure 41. Radial distribution of fluoride in a cylindrical single crystal of HF-doped ice.

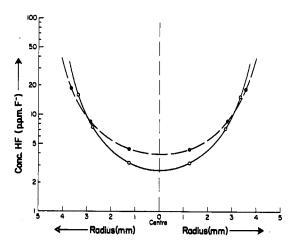


Figure 42. Radial distribution of fluoride ion in a cylindrical single crystal of HF-doped ice.

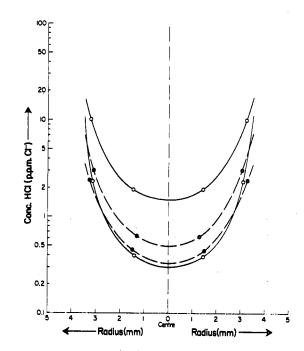


Figure 43. Radial distribution of chloride ion in a cylindrical single crystal of HC1-doped ice.

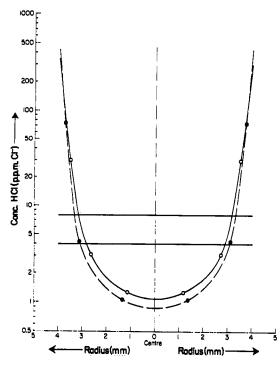


Figure 44. Radial distribution of chloride ion in a cylindrical single crystal of HC1-doped ice.

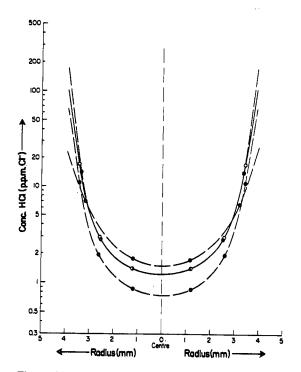


Figure 45. Radial distribution of chloride ion in a cylindrical single crystal of HC1-doped ice.



