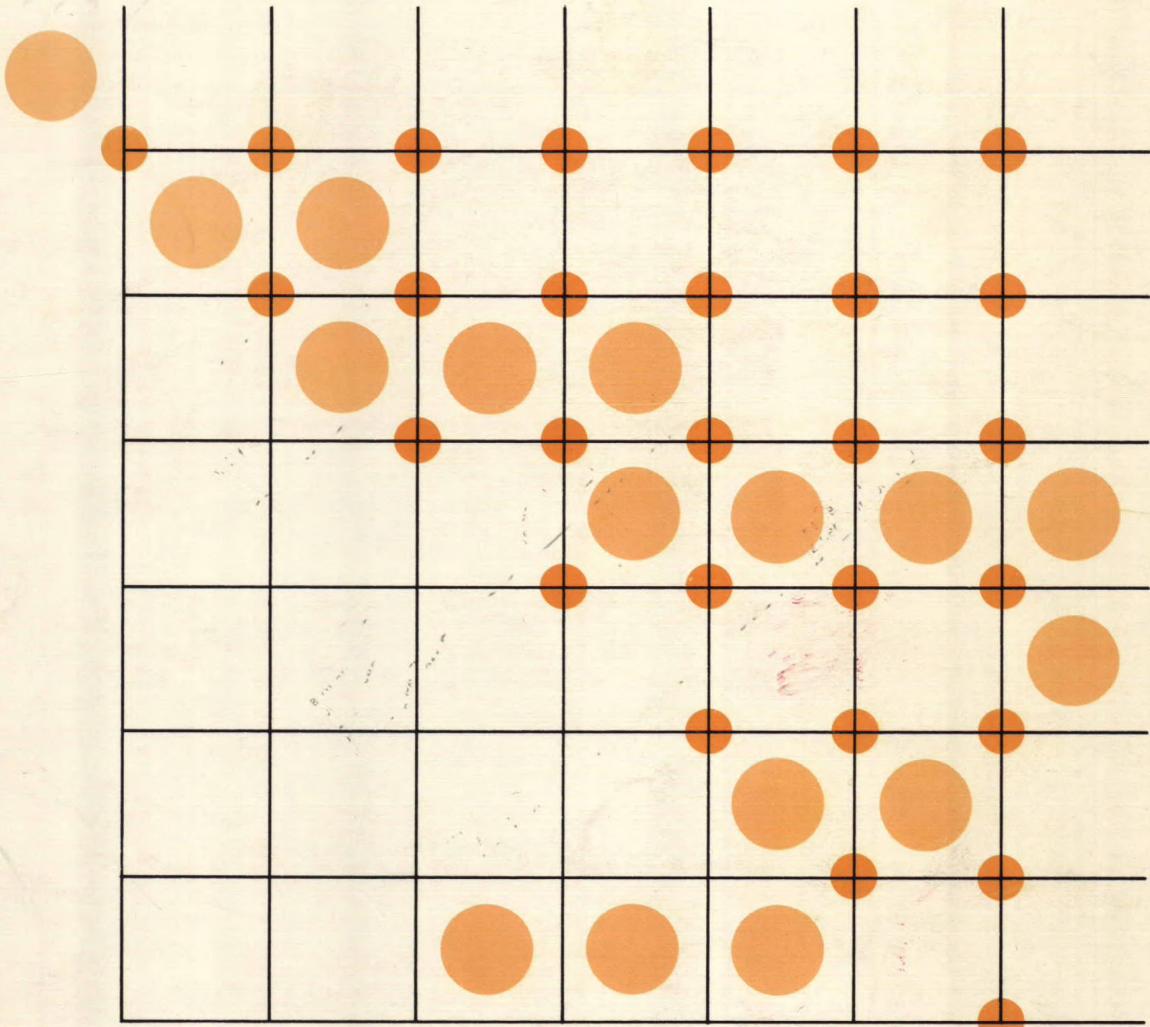


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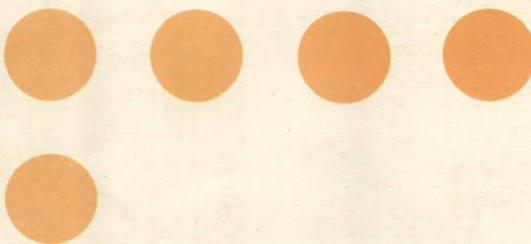
A STUDY OF MUTUAL COUPLING EFFECTS
IN FINITE MICROSTRIP ARRAYS

DSS CONTRACT 06ST.36100-5-0090
FINAL REPORT



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Final Report

Prepared for

Communication Research Centre

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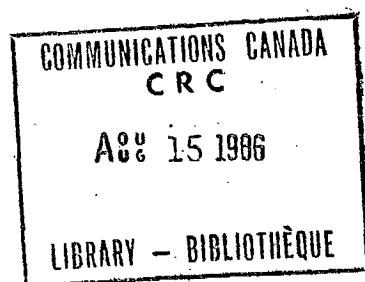
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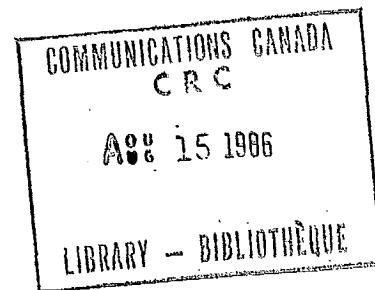
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SUMMARY

In this report the mathematical formulation of mutual coupling between elements of a microstrip array is provided. The generated models are used to investigate both the impedance and the array radiation patterns. The computed array patterns indicate that, generally, the mutual coupling reduces the array gain, but has a small effect on the sidelobe levels. The problem was also studied numerically for coupling between two patch arrays and reported in the program report. It indicated that the mutual coupling increases with beam scan and becomes considerable when the array beam approaches the array plane.

A case of two-dimensional array is also studied. A 19-element array using circular patch elements having a hexagonal grid is considered. It is fabricated using a substrated height of $h = 1/16$ inch and a permittivity of $\epsilon_r = 2.65$. The array beam is scanned to the direction of $\phi_o = 7.5^\circ$, $\theta_o = 55^\circ$, using coaxial lines of appropriate length to feed the array elements. The array beam is then measured experimentally. The array radiation pattern is also computed with and without the mutual coupling. Comparing the results it is found that the array gain is reduced considerably due to the mutual coupling. The computed result with the inclusion of mutual coupling is, otherwise, in good agreement with the experiment.

Recommendation for future work

The work in this phase of study was involved with the development of mathematical models, computation of mutual coupling effect on the

array patterns and a comparison of the results with experiment. Several areas, however, need careful examinations, which should be studied during the second phase of project.

- 1) the effect of mutual coupling on the input impedance has not been examined experimentally. This should be studied with sample arrays and compared with experimental data.
- 2) So far, the effect of coupling is considered for the dominant mode. In certain applications higher order modes may be used and their coupling should also be investigated.
- 3) So far, the experimental comparison is made for the main polarization. The coupling effect on the cross-polarization can provide additional information and because of its small magnitude may become too sensitive to scanning and inter-element coupling.
- 4) The radiation patterns of isolated patches are normally obtained approximately, using the fringing field concept. This method fails to provide accurate radiation patterns near the horizontal plane. Further improvement in the pattern prediction is desirable to enable better co-polar and cross-polar pattern calculations.

1. INTRODUCTION

In any array environment of finite elements in planar configuration, the elements in close proximity to each other behave differently than when they are isolated. The problem is particularly difficult in small arrays since the elements are in asymmetrical positions and the mutual coupling differs for each elements. Qualitatively, the literature survey indicates that the mutual coupling between the array elements can have a substantial effect on the total performance of the antenna as compared with the standard simplified prediction that sums up the relative phase location of every element weighted by its own excitation only. This simplified array factor, when it multiplies the far field of a single isolated patch, generates the array radiation pattern without incorporating the mutual interaction. In order to introduce the coupling effect into the array theory, these weighting factors for all elements are treated as an excitation coefficient column vector that is linearly related to the actual excitation vector by a transformation matrix ($N \times N$) where N is the number of the patches in the array. This transformation matrix can be dimensionless, an admittance in mhos, or an impedance in ohms, as the dependent and independent variables in the array are defined. The diagonal elements in this transformation matrix behave exactly as the simple array theory mentioned above while the off-diagonal elements of the matrix accounts for the mutual interaction between the patches. Since the transformation matrix is deducible from laboratory measurements [11], [23] or by analytical means [20], the current work is devoted to the analytical approach to build a complete and valid model for practical use in an-

tenna analysis or synthesis routine of microstrip planar arrays. It is obvious that any analytical study of the interaction effect must conform to the method of analysis used for the single resonant patch of the array in order to introduce the coupling correctly. The rectangular patch and the circular patch are the most commonly analyzed and measured and accordingly the two cases are selected for the current analysis. The more rigorous treatments of these resonant patches are considered cumbersome and time consuming for regular engineering applications. On the other hand it has been shown [1], for the rectangular patch, that the TEM transmission line model provides a numerically efficient substitute specially as the model has been improved and modified [2] so as to be free from earlier limitations. The rectangular resonant patch in the transmission line model has been viewed in the literature [1], [2], [20], [23] as equivalent to two radiating apertures both linked together internally by the transmission line and its excitation while mutually coupled to themselves (as well as to all other apertures in the array) via the free space radiation interaction mechanism of parallel magnetic current strips. The analysis is simplified further by neglecting the interaction through the surface waves, which limits the validity of the analysis to patches having substrates of small permittivity and thickness. The experimental verification of this validity has been provided [23] for permittivities up to 2.5 and substrate thickness up to $0.15\lambda_0$. The electromagnetic field under the patch can be uniquely identified in terms of the excitation current and its location as well as an impedance (or admittance) boundary condition at the apertures. Note that the common methods of excitation are the coaxial cable from behind the ground plane positioned somewhere under

the patch or microstrip line attached to the perimeter of the patch on the top of the substrate. The impedance (or admittance) boundary condition is identified directly by its stationary integral expression [9] relating the interaction between two separate apertures with magnetic currents of the same spacial distribution. Much of the previous work [20], [21] dealt with two rectangular patches (E-plane coupled, H-plane coupled, or a general translated orientation without relative rotation) where extensive comparison with the experimental observations was made in order to confirm the validity of this model. Hence an attempt is made to generalize the theory and derive analytically the matrix relation between the aperture electric field (or the radiating edge voltage) and the excitation column vector (input currents) for a rectangular patch planar array based on the TEM-Transmission line model.

The other case selected for the analysis is the circular disc patch planar microstrip array. The literature survey [1]-[6] confirms that the cavity model for the circular patch is the simplest and most compact. This model dictates that the electromagnetic field in the substrate is dominantly transverse magnetic with respect to the cylindrical radial direction. Its decomposition into TM modes gives a complete set capable of representing any arbitrary TM-field at the perimeter of the disc where the mutual interaction between different apertures can be introduced. The patch viewed from the top, when separated from the ground plane by the substrate creates a cavity with cylindrical wall of circumferential magnetic current ribbon, formed by dropping perpendiculars from the patch perimeter to the ground plane. Hence the terminology "Cavity Model" is used and the magnetic current ribbon is visualized to radiate from its outer surface while the inside surface

is choked by the cavity. The cavity also serves as the means of transferring the energy needed, by the single radiating aperture, from the current source. As mentioned above, the field inside the cavity can be uniquely identified from the current source, its location, and the impedance boundary condition at the perimeter relating the circumferential magnetic field to the axial electric field of the aperture. In the planar microstrip array situation with mutual interaction included, one invokes a linear relation between the aperture magnetic field of one element and all existing aperture electric fields of elements in the array. This step brings the free space radiation interaction between magnetic current ribbons into effect and the corresponding variational expression for matrix elements are defined without the choking effect of the cavity. Then the Green's function solution of an excited cavity with impedance boundary condition completes the model by relating the diagonal element of the radiation interaction matrix to the cavity characteristics and its excitation. Although this cavity model is also applicable to the rectangular patch case, no attempt is made to pursue it since the trend in the literature has accepted, for the dominant mode of the rectangular patch applications, the TEM-transmission line model as the simplest and the most suited analysis for efficient evaluation of the patch's characteristics.

The analysis to follow starts with the circular patch case as it is the more compact of the two and analytically more tractable because the array has only one aperture per element and the electromagnetic field under the patch behaves according to the dielectric characteristic of the substrate ϵ_r . In the case of the rectangular patch array the analysis deals with two apertures per element and the TEM-mode used

in the transmission line model guides the energy in the air and partially in the dielectric. Hence the analysis must rely on an estimated effective permittivity that is experimentally verifiable. The complete analysis is formulated in this report with the principal electromagnetic field quantities (E, H) and only in the last stage of the model computation a transfer to the more common electric parameter (currents, voltages, impedances, and so on) are executed.

2. THE CIRCULAR PATCH MODEL

A circular disc patch operating in one of its resonant modes is the most prevalent microstrip antenna configurations. It offers performance, in its TM_{11} -mode, similar to that of the rectangular geometry and might have in some applications certain advantages over other shapes. Of the various methods available for the analysis of the disc patch antenna [1]-[6], the Green's function method or its equivalent radial transmission line modelling provides a reasonably good agreement between calculated and measured input impedances and predicts adequately the radiation characteristics.

The two parallel metallic circular discs separated by a small distance apart has been used as a means for guiding the electromagnetic energy flow radially outward in the cylindrical sense as in radial transmission line theory [7]. When the separation between the plates is much less than a half wave length, the transverse magnetic modes with circumferential variation and no variation along the axial direction controls this mechanism of the power transfer. The metallic disc patch over the ground plane separated by a thin substrate of relative permittivity ϵ_r and thickness h provides the same mentioned environ-

ment and the electromagnetic field in the substrate under the patch may be expressed in its circumferential modal expansion as follows:

$$E_z = \sum_n C_n z_n(k\rho) \cos [n(\phi - \phi_0)]$$

$$H_\phi = -\frac{j}{\eta} \sum_n C_n z'_n(k\rho) \cos [n(\phi - \phi_0)]$$

$$H_\rho = -\frac{j}{\eta} \sum_n \frac{C_n}{k\rho} z_n(k\rho) \sin [n(\phi - \phi_0)]$$

where

$$z_n(k\rho) = \begin{cases} J_n(k\rho) + A_n Y_n(k\rho) & , \rho \neq 0 \\ J_n(k\rho) & , (\rho = 0) \text{ eV} \end{cases}$$

The two main field components E_z and H_ϕ in this model can be treated in some normalized sense as voltages and currents on a radial transmission line. For example, $(-E_z h)$ can be taken to be a total voltage between the ground plane and the patch (upper plate as positive) while $(2\pi\rho H_\phi)$ represents the total radial outward current in the upper plate (the patch). Other conventions like $(I = -2\pi\rho H_\phi, V = E_z h)$ is also acceptable and in the following analysis the work concentrates on the proportionality between the two quantities $(-\eta_0 H_\phi)$ and (E_z) which is a dimensionless ratio and leaves the final definition of impedances and admittances until the point of comparison with other published results and the extraction of the equivalent electric circuit parameters when the

array is treated as an N-port network. The analysis proceeds step by step in a self-contained fashion without any outside supplement until the need arises for the empirical extension formula to account for the fringing field effect at the aperture. First we extend the model of the cavity to include a current excitation source where the geometry of a single cell, say the i -th patch, is shown in figure (1) with a current I_i pointed out of the figure (the z-direction) at the location (ρ_i, ϕ_i) and an angular spread (2δ) . Without loss of generality, the effective radial thickness of the source is assumed infinitesimal which makes the cavity react to a surface current source \bar{K}_e^i that is decomposable into a circumferential spectrum at $\rho = \rho_i$.

$$\begin{aligned}\bar{K}_e^i &= \hat{z} \frac{I}{2\pi\rho_i} \delta(\rho-\rho_i) \left\{ U [\phi - (\phi_i - \delta)] - U [\phi - (\phi_i + \delta)] \right\} \\ &= \hat{z} \frac{I}{2\pi\rho_i} \delta(\rho-\rho_i) \left(\sum_0^{\infty} \left(\frac{2}{1+\delta_{no}} \right) \left(\frac{\sin n\delta}{n\delta} \right) \cos [n(\phi-\phi_i)] \right)\end{aligned}\quad (1)$$

But the two dimensional wave equation with no z-variation

$$(\nabla^2 + k^2) E_z^i = jk\eta J_z^i$$

indicates that the electric field (E_z^i), under the i -th patch, have the same ϕ -dependance as its own excitation current especially when outside influences such as coupling and scattering are ignored. Therefore, the electromagnetic field under the patch can be written in the following form for $\rho_i < \rho < b$,

$$E_z^i = \sum_0^\infty E_i^n \left(\frac{J_n(k\rho) + A_n^{-i} Y_n(k\rho)}{J_n(kb) + A_n^{-i} Y_n(kb)} \right) \cos [n(\phi - \phi_i)] \quad (2)$$

$$H_\phi^i = -\frac{j}{\eta} \sum_0^\infty E_i^n \left(\frac{J'_n(k\rho) + A_n^{-i} Y'_n(k\rho)}{J_n(kb) + A_n^{-i} Y_n(kb)} \right) \cos [n(\phi - \phi_i)] \quad (3)$$

and for $0 < \rho < \rho_i$

$$E_z^i = \sum E_i^n \left(\frac{J_n(k\rho_i) + A_n^{-i} Y_n(k\rho_i)}{J_n(kb) + A_n^{-i} Y_n(kb)} \right) \frac{J_n(k\rho)}{J_n(k\rho_i)} \cos [n(\phi - \phi_i)] \quad (4)$$

$$H_\phi^i = -\frac{j}{\eta} \sum E_i^n \left(\frac{J_n(k\rho_i) + A_n^{-i} Y_n(k\rho_i)}{J_n(kb) + A_n^{-i} Y_n(kb)} \right) \frac{J'_n(k\rho)}{J_n(k\rho_i)} \cos [n(\phi - \phi_i)] \quad (5)$$

In the above four equations (2)-(5) the continuity of the electric field from ρ_i^- to ρ_i^+ (the source radial location) is introduced and the constant term E_i^n becomes a representatives of the maximum amplitude of the modal electric field distribution in the aperture. The magnetic field as written is only consistant with the Maxwell's equation (proportional to the derivative of E_z) and its discontinuity at $\rho = \rho_i$ can be invoked to conform to the existance of the source current according to,

$$H_\phi(\rho = \rho_i^+) - H_\phi(\rho = \rho_i^-) = K_e$$

This defines the maximum amplitude E_i^n of the modal electric field at the aperture as a function of the excitation current and the constant A_n^i as follows:

$$E_i^n = jk\eta I \left(\frac{1}{2[1 + \delta_{no}]} \right) \left(\frac{\sin n\delta}{n\delta} \right) J_n(k\rho_i) \left[\frac{J_n(kb)}{A_n^i} + Y_n(kb) \right] \quad (6)$$

Equations (2) and (3) at the aperture defines the parameter A_n^i in terms of the field's ratio $p_i^n = (-\eta_0 H_\phi^i/E_z^i)|_{\rho=b}$. That is,

$$A_n^i = - \begin{pmatrix} p_i^n J_n(kb) - j \frac{\eta_0}{\eta} J'_n(kb) \\ p_i^n Y_n(kb) - j \frac{\eta_0}{\eta} Y'_n(kb) \end{pmatrix} \quad (7)$$

$$E_i^n = - \frac{\eta_0 I_i}{\pi b} \frac{\frac{\sin n\delta}{n\delta}}{1 + \delta_{no}} \begin{pmatrix} \frac{J_n(k\rho_i)}{J_n(kb)} \\ p_i^n - j \frac{\eta_0}{\eta} \frac{J'_n(kb)}{J_n(kb)} \end{pmatrix} \quad (8)$$

Which means, the two dimensional wave equation gives the field in two regions, for $\rho_i < \rho < b$

$$E_z^i = \frac{j\eta k I_i}{2} \sum \frac{\frac{\sin n\delta}{n\delta}}{1 + \delta_{no}} J_n(k\rho_i) \left[Y_n(k\rho) + \frac{1}{A_n^i} J_n(k\rho) \right] \cos [n(\phi - \phi_i)] \quad (2')$$

$$H_\phi^i = \frac{k I_i}{2} \sum \frac{\frac{\sin n\delta}{n\delta}}{1 + \delta_{no}} J_n(k\rho_i) \left[Y'_n(k\rho) + \frac{1}{A_n^i} J'_n(k\rho) \right] \cos [n(\phi - \phi_i)] \quad (3')$$

and for $0 \leq \rho \leq \rho_i$

$$E_z^i = \frac{j\eta k I_i}{2} \sum_n \frac{\sin n\delta}{1+\delta_{no}} J_n(k\rho) \left[Y_n(k\rho_i) + \frac{1}{A_n^i} J_n(k\rho_i) \right] \cos [n(\phi-\phi_i)] \quad (4)$$

$$H_\phi^i = \frac{kI_i}{2} \sum_n \frac{\sin n\delta}{1+\delta_{no}} J'_n(k\rho) \left[Y_n(k\rho_i) + \frac{1}{A_n^i} J_n(k\rho_i) \right] \cos [n(\phi-\phi_i)] \quad (5)$$

consistent with Maxwell's equations, the continuity of the electric field at $\rho = \rho_1$, and the discontinuity of the magnetic field is in accordance with the excitation current as given in Fig. (1). The only unknown left in the equations is the modal constant A_n^i , which is selected in the present work from the boundary condition at $\rho = b$. This is given by equation (7) in terms of a normalized (dimensionless) admittance parameter p_i^n defined above. Therefore, the circumferential spectrum of the main field components at the aperture are represented by the two relations

$$(E_z^i)_{aperture} = -\frac{\eta_o I_i}{\pi b} \sum_n \left(\frac{\sin n\delta}{1+\delta_{no}} \right) \left(\frac{\frac{J_n(k\rho_i)}{J_n(kb)}}{p_i^n - j \frac{\eta_o}{\eta} \frac{J'_n(kb)}{J_n(kb)}} \right) \cos [n(\phi-\phi_i)] \quad (9)$$

$$(H_\phi^i)_{aperture} = \frac{I_i}{\pi b} \sum_n \left(\frac{\sin n\delta}{1+\delta_{no}} \right) \left(\frac{p_i^n \frac{J_n(k\rho_i)}{J_n(kb)}}{p_i^n - j \frac{\eta_o}{\eta} \frac{J'_n(kb)}{J_n(kb)}} \right) \cos [n(\phi-\phi_i)] \quad (10)$$

From equation (2)' the input impedance can be defined according to,

$$Z_{in} = \frac{-j\eta kh}{2} \sum \frac{(\sin n\delta)}{1+\delta_{no}} J_n(k\rho_i) \left[Y_n(k\rho_i) + \frac{1}{A_n} J_n(k\rho_i) \right] \quad (11)$$

and equation (9) indicates that, for the special case of a microstrip line feed connected to the edge, the input impedance becomes

$$Z_e = \frac{\eta_0 h}{\pi b} \sum \frac{\sin n\delta}{1+\delta_{no}} \left(\frac{1}{p_i^n - j \frac{\eta_0}{\eta} \frac{J'_n(kb)}{J_n(kb)}} \right) \quad (12)$$

Up to this step in the analysis, all information on the cavity characteristics and its current excitation have been introduced. The given equations become uniquely defined whenever the parameters p_n^i are established which can be evaluated from the total array configuration, as seen in Fig. (2), and according to the free space radiation interaction mechanism. Before this step is explained, certain simplification can be identified. In order to achieve radiation coherence from the planar array, the individual currents must be injected into the patches in a similar fashion (i.e. ϕ_i is the same for all i 's) and without loss to the generality of the analysis ϕ_i can be taken as zero. This assumption, as Fig. (2) shows, establishes the xz-plane as the E-plane of the array while the yz-plane is taken here as the H-plane of the array. The radiation is emitted from a magnetic current ribbon of radius b and a height h , in place of the aperture, backed by a perfectly conducting metallic pill box in place of the cavity. At any arbitrary aperture in the array, say the i -th, the aperture field when decomposed in its circumferential spectrum reads

$$H_{\phi}^i = \sum H_i^n \cos(n\phi), E_z^i = \sum E_i^n \cos(n\phi)$$

To achieve the mutual coupling and its link to the above cavity formulation, a linear proportionality between the magnetic aperture field of one element and the aperture electric fields of all existing cells in the array must be assumed. This is accomplished by invoking on the above aperture spectrum the following proportionality

$$H_i^n = -\frac{1}{\eta_0} \sum_j^N p_{ij}^n E_j^n \quad (13)$$

That is,

$$p_i^n E_i^n = \sum_j^N p_{ij}^n E_j^n \quad (14)$$

Equation (13) is for an arbitrary element and can be generalized for the total array by considering those maximum amplitudes of the aperture field distribution as column vectors $H^{(n)}$, $E^{(n)}$ connected together by the matrix relation,

$$H^{(n)} = -\frac{1}{\eta_0} p^{(n)} E^{(n)} \quad (15)$$

which is known as the short circuit admittance formulation of the coupling because of the definition of the $p^{(n)}$ matrix,

$$\begin{aligned} p_{ij}^{(n)} &= -\eta_0 \left[\frac{H_i^n}{E_j^n} \right], \quad E_k^n = 0 \text{ when } k \neq j \\ &= -\eta_0 \left[\frac{(H_{\phi}^i)_{w.av.}}{E_j} \right], \quad E_k^n = 0 \text{ when } k \neq j \end{aligned} \quad (16)$$

As seen from the purely mathematical definition of equation (16), the magnetic field H_i^n on the i -th aperture is generated from the existence of an assumed electric field E_j^n on the j -th aperture only, while all others are shorted. That means, the weighted average circumferential magnetic field $(H_\phi^i)_{\text{wav}}$ is generated by a magnetic current ribbon of a known strength ($K_m^j = \hat{\phi} E_j^n \cos n\phi$) under the influence of the existing shorts. The word "weighted average" is clear in this case from the orthogonality of the spectrum,

$$(H_\phi^i)_{\text{w. av.}} = H_i^n = \frac{\int H_\phi^i \cos(n\phi) d\phi}{\int \cos^2(n\phi) d\phi}$$

This definition is extended to the interaction between identical apertures of arbitrary shapes, but similar distribution $w(s)$ on the surface to yield the following general expression

$$(\hat{l}_i \cdot \bar{H}^i)_{\text{w. av.}} = \frac{\iint ds_i W(s_i) (\hat{l}_i \cdot \bar{H}^i)}{\iint ds_i W^2(s_i)}$$

where \hat{l}_i is a tangential unit vector to the circumference ($\hat{z} \times \hat{n}_i = \hat{l}_i$). The unification of this approach to patches (apertures) of different shapes is accomplished by assuming a standardized electric field in the z -direction at the aperture of a real distribution function $w(s)$ that corresponds to the mode in question ($E_z^j = E_j^j W(s_j) \hat{z}$). Let the outward normal \hat{n}_j to the aperture points into the radiation space and is a function of the patch's perimeter. This electric field yields a magnetic surface current ($K_m^j = E_j^j W(s_j) \hat{l}_j$), where the unit vector \hat{l}_j conforms to the right hand rule of the cross product $\hat{l}_j = \hat{z} \times \hat{n}_j$.

Accordingly, the definition of the short circuit admittance parameter $p_{ij}^{(n)}$ becomes,

$$p_{ij}^{(n)} = -\eta_0 \left[\frac{(\hat{l}_i \cdot \bar{H})_{w, av.}}{E_j^n} \right], \quad E_k^n = 0 \text{ for } k \neq j$$

As the magnetic field is generated by the magnetic current \bar{K}_m^j , its mathematical expression becomes

$$\bar{H} = \iint ds_j \mathbf{G}^a \cdot \bar{K}_m^j, \quad \mathbf{G}^a = (\mathbf{I} + \frac{\nabla \nabla}{k_o^2}) \left(\frac{k_o}{j 2\pi \eta_0} \frac{e^{-jk_o R}}{R} \right)$$

Note that the second order tensor \mathbf{G}^a is twice the free space Green's dyadic because of the image needed for the "Equivalence Theorem" [8].

Then,

$$\hat{l}_i \cdot \bar{H} = E_j \iint ds_j W(s_j) [\hat{l}_i \cdot \mathbf{G}^a \cdot \hat{l}_j]$$

$$(\hat{l}_i \cdot \bar{H})_{w, av.} = E_j \frac{\iint ds_i W(s_i) \iint ds_j W(s_j) [\hat{l}_i \cdot \mathbf{G}^a \cdot \hat{l}_j]}{\iint ds_i W^2(s_i)}$$

Therefore,

(17)

$$p_{ij}^{(n)} = - \frac{\iint \iint ds_i ds_j W(s_i) W(s_j) f(\hat{l}_i | \hat{l}_j)}{\iint ds_i W^2(s_i)} \quad (17)$$

$$f(\hat{l}_i | \hat{l}_j) = \frac{e^{-jk_o R}}{j 2\pi k_o R^3} \left[2(1+jk_o R) (\hat{l}_i \cdot \hat{l}_j) + \left(k_o^2 R^2 - 3[1 + jk_o R] \right) (\hat{R} \times \hat{l}_i) \cdot (\hat{R} \times \hat{l}_j) \right] \quad (18)$$

and R represents the distance between the source \wedge observation points.

The above generalized expression for $p_{ij}^{(n)}$, equations (17) and (18), coincides within a normalization factor with the stationary expression for the short circuit admittance derived from the variational principle and is given in the literature [8], [9]. To demonstrate the actual use of this generalized interaction formula, the current case of the circular disc patch array is brought back for evaluation. Since every patch in the array can be defined by its centre coordinates (x_i, y_i) , radius b , thickness h , $\hat{l}_i = \hat{\phi}_i$, and $\hat{n}_i = \hat{\rho}_i$, then the definition of an element of area $ds_i = bdz d\phi_i$ and the distribution function $W(s_i) = \cos(n\phi_i)$ becomes obvious. To investigate the interaction between this patch with any other arbitrary element of the array, say the j -th, the details of the geometry is explained in Fig. (3). The separation between the two patches and their relative orientation becomes,

$$D_m = \sqrt{(x_j - x_i)^2 + (y_j - y_i)^2}, \quad \Phi_m = \tan^{-1} \left[\frac{Y_j - Y_i}{X_j - X_i} \right]$$

Let the two area elements be at two arbitrary locations defined by the variables ϕ_i and ϕ_j respectively, as shown by the horizontal projection plane of Fig. (3). The arbitrary location ϕ_j defines the triangle ($D_m b L$) completely, since the angle between D_m and b becomes $(\pi + \phi_m - \phi_j)$ and yields the expression for L ,

$$L = \sqrt{D_m^2 + b^2 + 2b D_m \cos(\phi_j - \phi_m)}$$

The other arbitrary surface element at ϕ_1 defines the triangle (L₀) and the solution for ρ becomes

$$\rho = \sqrt{L^2 + b^2 - 2b L \cos\theta}, \quad \theta = (\phi_i - \phi_m) - \sin^{-1} \left[\frac{b}{L} \sin(\phi_j - \phi_m) \right]$$

$$\beta = \theta + \sin^{-1} \left[\frac{b}{\rho} \sin\theta \right], \quad R = \sqrt{(z_j - z_i)^2 + \rho^2}$$

$$\hat{l}_i \cdot \hat{l}_j = \cos(\phi_j - \phi_i), \quad \hat{R} = \frac{\rho}{R} \hat{\rho} + \frac{z_j - z_i}{R} \hat{z}$$

$$(\hat{R} \times \hat{l}_i) \cdot (\hat{R} \times \hat{l}_j) = \frac{l^2}{R^2} \cos(\beta) \cos(\beta + \phi_j - \phi_i) + \frac{(z_j - z_i)^2}{R^2} \cos(\phi_j - \phi_i)$$

Therefore,

$$\begin{aligned} p_{ij}^{(n)} &= - \frac{\left(\frac{b}{n}\right)}{(1 + \delta_{no}) \pi} \iiint dz_i d\phi_i dz_j d\phi_j \cos(n\phi_i) \cos(n\phi_j) f_c(\hat{l}_i \mid \hat{l}_j) \\ f_c(\hat{l}_i \mid \hat{l}_j) &= \frac{e^{-jk_o R}}{j 2\pi k_o R^3} \left[\left\{ \left(2 - 3 \frac{(z_j - z_i)^2}{R^2} \right) (1 + jk_o R) + k_o^2 (z_j - z_i)^2 \right\} \cos(\phi_j - \phi_i) \right. \\ &\quad \left. + \left\{ k_o^2 \rho^2 - 3 \frac{\rho^2}{R^2} (1 + jk_o R) \right\} \cos\beta \cos(\beta + \phi_j - \phi_i) \right] \end{aligned} \quad (19)$$

When the minimum separation ($D_m - 2b$) is much greater than h , the approximation of null z -variation in $f_c(\hat{l}_i \mid \hat{l}_j)$ and $R \approx \rho$ is acceptable. This reduces the above double area integration to a more manageable double contour integration,

$$\begin{aligned} p_{ij}^{(n)} &= \frac{j(\frac{h}{b})}{2\pi^2 (1 + \delta_{no})} \iint d\phi_i d\phi_j \cos(n\phi_i) \cos(n\phi_j) \left(\frac{b}{\rho} \right)^2 \frac{e^{-jk_o \rho}}{k_o \rho} \cdot \\ &\quad \left[2(1 + jk_o \rho) \cos(\phi_j - \phi_i) + \left(k_o^2 \rho^2 - 3[1 + jk_o \rho] \right) \cos\beta \cos(\beta + \phi_j - \phi_i) \right] \end{aligned} \quad (20)$$

Equation (20) can be executed numerically for lack of a closed form solution. Although equations (19) and (20) hold also for the self

interaction of the diagonal elements in the matrix, where $s_i = s_j$, the evaluation is sensitive to handle in the neighbourhood of the singularity $R = 0$. If the integral expression is split in real and imaginary parts, the real part is found to be behaving well and render itself to a numerical evaluation.

$$\text{Re} [p_{ii}^{(n)}] = \frac{(\frac{b}{h})}{2\pi} \int \int dz dz \cdot \int d\alpha \frac{\cos(n\alpha)}{k_o R^3} \cdot$$

$$\left[\left(2 - \left(\frac{p}{2b} \right)^2 \right) \left(\sin(k_o R) - k_o R \cos(k_o R) \right) - \left(\frac{p}{2b} \right)^2 k_o^2 R^2 \sin(k_o R) \right. \\ \left. + \left(k_o^2 R^2 \sin(k_o R) - 3[\sin(k_o R) - k_o R \cos(k_o R)] \right) \left(1 - \left(\frac{p}{2b} \right)^2 \right) \frac{(z - z)^2}{R^2} \right]$$

where, $R = \sqrt{p^2 + (z - z)^2}$, $p = 2b \sin(\frac{\alpha}{2})$

The approximation of the above for thin aperture becomes,

$$\text{Re} [p_{ii}^{(n)}] = \frac{(\frac{h}{b})}{2\pi} \int d\alpha \cos(n\alpha) \left(\frac{b}{p} \right)^2 \frac{1}{k_o p} \left[\left(1 + \cos^2 \frac{\alpha}{2} \right) \left(\sin(k_o p) \right. \right. \\ \left. \left. - k_o p \cos(k_o p) \right) - k_o^2 p^2 \sin(k_o p) \sin^2 \left(\frac{\alpha}{2} \right) \right] \quad (21)$$

The radius ($p = 2b \sin(\frac{\alpha}{2})$) upon substitution in Taylor-expansions of the trigonometric function yields an integrand polynomial in $(k_o b)$ within trigonometric coefficients in α that can be integrated exactly and the outcome is as follows:

$$\text{Re} [p_{ii}^{(0)}] = (\frac{h}{b}) \left[\frac{(k_o b)^4}{6} - \frac{(k_o b)^6}{30} + \frac{(k_o b)^8}{336} - \frac{(k_o b)^{10}}{6480} + \dots \right]$$

$$\text{Re} [p_{ii}^{(1)}] = (\frac{h}{b}) \left[\frac{(k_o b)^2}{3} - \frac{2}{15} (k_o b)^4 + \frac{11}{420} (k_o b)^6 - \frac{11}{5436} (k_o b)^8 \right. \\ \left. + \frac{37}{285120} (k_o b)^{10} + \dots \right]$$

$$\text{Re} [p_{ii}^{(2)}] = (\frac{h}{b}) \left[\frac{(k_o b)^4}{20} - \frac{(k_o b)^6}{84} + \frac{29}{22680} (k_o b)^8 - \frac{19}{249480} (k_o b)^{10} + \dots \right]$$

In fact, we found from an intirely different technique that a general formula can be generated according to:

$$\begin{aligned} \text{Re} [p_{ii}^{(n)}] &= \left(\frac{h}{b}\right) \left[\frac{k_o b}{4} \int_0^{2k_o b} \left(J_{2(n-1)}(x) + J_{2(n+1)}(x) \right) dx - \frac{n^2}{2k_o b} \int_0^{2k_o b} J_{2n}(x) dx \right] \\ &= \left(\frac{h}{b}\right) \sum_0^{\infty} \frac{(-1)^m (k_o)^{2(m+n)}}{m! (m+2n)!} \left[n + \frac{m(m-1)}{2(m+n)-1} - \frac{n^2}{2(m+n)+1} \right] \end{aligned} \quad (22)$$

Since it has been shown in reference [2] that a numerical evaluation of equation (21) gives good agreement with Shen's results [10], based on the far field calculation, then the confidence in formulation (22) for accuracy is established. Although the rest of the admittance can still be handled analytically from the imaginary part of equation (19) with a special care, the trend in the literature is to fit the reactive part so as to agree with the measured value of the modal resonance. To account for this step we assume

$$IM [p_{ii}^{(n)}] = \frac{\eta_o}{\eta} \left(\frac{J_n'(kb_e) Y_n'(kb)}{Y_n'(kb_e) J_n(kb)} \right) \quad (23)$$

where b_e is a parameter yet to be determined. Note that the reactive part as assumed is in fact the normalized input susceptance at ($\rho = b$) for a radial transmission line opened at $\rho = b_e$. To see how b_e is to be estimated we go back to equation (9) and write down the relation between the modal electric field amplitude at the aperture and the excitation input current as follows,

$$\left(p_i^n - j \frac{\eta_o}{\eta} \frac{J_n'(kb)}{J_n(kb)} \right) E_i^n = - \frac{\eta_o}{\pi b (1 + \delta_{no})} \left(\frac{\sin n\delta}{n\delta} \right) \frac{J_n(k\rho_i)}{J_n(kb)} I_i ,$$

and substitute relation (14) to obtain

$$\left(p_{ii}^{(n)} - j \frac{\eta_0}{\eta} \frac{J'_n(kb)}{J_n(kb)} \right) E_i^{(n)} + \sum_{j \neq i} p_{ij}^{(n)} E_j^{(n)} =$$

$$- \frac{\eta_0}{\pi b(1 + \delta_{no})} \left(\frac{\sin n\delta}{n\delta} \right) \frac{J_n(k\rho_i)}{J_n(kb)} I_i , \quad (24)$$

Since the above equation is valid for $i = 1, 2, \dots, N$, then the complete general relation of the array in matrix form becomes

$$P^{(n)} E^{(n)} = - \frac{\eta_0}{\pi b(1 + \delta_{no})} \left(\frac{\sin n\delta}{n\delta} \right) S^{(n)} I_z \quad (25)$$

where

$$(S^{(n)})_{ij} = \delta_{ij} \frac{J_n(k\rho_i)}{J_n(kb)} ; \delta_{ij} \text{ is } 1 \text{ for } i = j \text{ and } 0 \text{ otherwise}$$

$$(P^{(n)})_{ij} = (p^{(n)})_{ij} - j \frac{\eta_0}{\eta} \frac{J'_n(kb)}{J_n(kb)} \delta_{ij}$$

Note that the $P^{(n)}$ matrix differs from the radiation admittance matrix $p^{(n)}$ only in the diagonal element. The new diagonal element is the same for all i 's and by using (22) and (23) it can be written as

$$(P^{(n)})_{ii} = \left[\left(\frac{h}{b} \right) \sum_0^{\infty} \frac{(-1)^m (k_e b)^{2(m+n)}}{m! (m+2n)!} \left[n + \frac{m(m-1)}{2(m+n)-1} - \frac{n^2}{2(m+n)+1} \right] \right] \\ - j \frac{\eta_0}{\eta} \frac{2}{\pi kb} \left\{ \frac{\frac{J'_n(kb_e)}{J_n(kb)}}{\frac{Y'_n(kb_e) J_n(kb) - J'_n(kb_e) Y_n(kb)}{}} \right\} \quad (26)$$

The imaginary part of the above expression represents both the assumed fringe field effect and the loading of the cavity, and does indicate resonance at $J'_n(kb_e) = 0$ according to the m -th root $J'_n(K_{mn}) = 0$ of the Bessel's function. That is,

$$f_r = \frac{K_{mn} c}{2\pi b_e \sqrt{\epsilon_r}}, \quad c \text{ is the velocity of light in free space.}$$

It is found [5], [6], [11] that the effective radius b_e given by,

$$b_e = b \sqrt{1 + \frac{1}{\epsilon_r} \left(\frac{2h}{\pi b} \right) \left[\ln\left(\frac{\pi b}{2h}\right) + 1.7726 \right]}$$

predicts the observed resonance very well for $(b/h) \gg 1$.

If we assume no angular spread for the source location or that we are interested only in the earlier resonant modes so that $(\sin(n\delta)/n\delta)$ is approximately one, equation (25) renders the edge characteristics of the array's patch in the standard voltage, current, impedance, admittance form as follows: -

$$\mathbf{V}_e^{(n)} = \mathbf{Z}_e^{(n)} \mathbf{I}_{in} = \mathbf{Y}_n^{-1} \mathbf{S}^{(n)} \mathbf{I}_{in} \quad (27)$$

$$(\mathbf{Y}_n)_{ij} = \frac{j}{2\pi\eta_0} \int \int d\phi d\phi' \cos(n\phi) \cos(n\phi') F_{ij}$$

$$(\mathbf{Y}_n)_{ii} = \frac{1+\delta_{no}}{120} \left[\left(\sum \frac{(-1)^m (k_o)^{2(m+n)}}{m! (m+2n)!} \left[n + \frac{m(m-1)}{2(m+n)-1} - \frac{n^2}{2(m+n)+1} \right] \right) \right.$$

$$\left. - j \frac{2}{\pi(k_o h)} \left(\frac{\frac{J_n'(kb_e)}{J_n(kb)}}{Y_n'(kb_e) J_n(kb) - J_n'(kb_e) Y_n(kb)} \right) \right]$$

$$F_{ij} = \left(\frac{b}{\rho} \right)^2 \frac{e^{-jr}}{r} \left[2(1+jr) \cos(\phi' - \phi) + (r^2 - 3[1+jr]) \cos \beta \cos(\beta + \phi' - \phi) \right]$$

$$r = k_o \rho$$

Equation (2) can now be used at $(\rho = b)$ and $(\rho = \rho_i)$ in order to eliminate A_n^i and relate the input electric field to that at the edge plus the input excitation current which produces the following expression for the modal input network parameters,

$$\mathbf{V}_{in}^{(n)} = \mathbf{Z}_{in}^{(n)} \mathbf{I} = \left(\mathbf{S}^{(n)} \mathbf{Y}_n^{-1} \mathbf{S}^{(n)} + \mathbf{Z}_s^{(n)} \mathbf{S}^{(n)} \right) \mathbf{I} \quad (28)$$

$$\left(\mathbf{Z}_s^{(n)} \right)_{ij} = j \delta_{ij} \frac{\eta_0(k_o h)}{2(1 + \delta_{no})} \left[J_n(k \rho_i) Y_n(kb) - J_n(kb) Y_n(k \rho_i) \right] \quad (29)$$

Since the desired antenna characteristics such as high gain, beam scanning, or steering capability are possible when discrete patches are combined to form an array such as the one discussed or in different arrangement; a general program has to be written for planar array of arbitrary configuration. Array theory has been treated in many texts [8], [12], [13], [14], [15] with emphasis on the array factor ($T_T = T_x T_y$) emanating from relative spacial distribution of the patches ($R_j = R_o^{-\rho_j}$ $\cos \gamma_j$, and $\gamma_j = \cos \theta \cos \theta_j + \sin \theta \sin j \cos(\phi - \phi_j)$) weighted by a coefficient proportional to contribution of the j -th element ability to radiate under the influence of coupling to its neighbours. That is to say,

$$T_t^{(n)} = \sum_{k=1}^N \left(\mathbf{V}_e^{(n)} \right)_k e^{jk_o(X_k \sin \theta \cos \phi + Y_k \sin \theta \sin \phi)} \quad (30)$$

If we accept the theoretical far field radiation pattern for a magnetic current loop of radius b as the value to be modulated, then we can write the total far field of the array as follows: -

$$E_\theta^{(n)} = \left[-(j)^n (bk_o) \frac{e^{-jk_o R_o}}{R_o} \right] J'_n(k_o b \sin \theta) \cos(n\phi) T_t^{(n)} \quad (31)$$

$$E_\phi^{(n)} = \left[(j)^n (bk_o) \frac{e^{-jk_o R_o}}{R_o} \right] \frac{n \cos \theta J_n(kb \sin \theta)}{kb \sin \theta} \sin(n\phi) T_t^{(n)} \quad (32)$$

Relation (31), (32) and (27) completes the tool for studying the total radiation pattern as a function of independent excitation currents. By artificially equating the off diagonal matrix elements $(Y_n)_{ij}$ to zero we can also generate the hypothetical case of the array theory without coupling. Hence the quantitative analytical study of their difference becomes feasible.

3. THE RECTANGULAR PATCH MODEL

The simplest microstrip patch configuration is undoubtedly the rectangular microstrip element as shown in Fig. (4). This simplicity has led a large number of researchers [1], [2], [17], [18], [19] to make numerous attempts to predict and evaluate its radiation and circuit characteristics. The mutual coupling between two rectangular patches has been studied [20], [21] and measured using the TEM-transmission line model as the internal line of communication between the two radiating apertures and the current source. This is the simplest method used by the researchers to develop a step-by-step procedure for the antenna design and to augment other factors as the need arises [2] to improve the model to its final completion. If one wonders why then the choice is made in this report to start with the analysis of the circular disc patch instead. The main rational is that the disc is one aperture patch, where the orthogonality of the circumferential modes are evident and the radial propagation characteristics of the TM-modes under the patch is uniquely defined by the dielectric constant of the substrate and its thickness. In contrast, the case at hand deals with two apertures per each rectangular patch that are linked together and to the excitation source via a TEM-transmission line. Simultaneously

they are also coupled to all other radiating apertures, including each other, from the outside through the radiation interaction mechanism of parallel magnetic current strips. Also an estimate for the propagation characteristics of the transmission line must be established since the TEM-mode propagates the energy in a mix of free space and the dielectric media.

Since in most of the common applications the rectangular patch array excites the transmission line mode of uniform aperture field distribution, the current analysis conforms to this demand and can in future be extended to any other required modal distribution. Because we are dealing with two, left and right, radiating apertures per patch in the planar array seen in Fig. (5), we assume the injected excitation current to split the resonant length L_i to L_i^- and L_i^+ . The corresponding electric field at these apertures are E_i^- and E_i^+ , assumed in the positive z-direction. The quantities $(E_i^- h)$ and $(E_i^+ h)$ are the left and right edge voltages by convention in this case and Fig. (6) indicates the corresponding quantities needed to evaluate the interaction admittance by the generalized variational expression generated in the previous section.

Since only the uniform aperture mode is discussed, the superscript (n) is dropped and the aperture magnetic field is defined according to the radiation convention $\hat{n}_i^- = -\hat{x}$, $\hat{n}_i^+ = +\hat{x}$, $\hat{l}_i^- = \hat{z} \times \hat{n}_i^-$, $\hat{l}_i^+ = \hat{z} \times \hat{n}_i^+$. That is,

$$\bar{H}_{aperture}^- = -\mathbf{H}_i^- \hat{l}_i^- , \quad \mathbf{H}_i^- = -\left(\hat{l}_i^- \cdot \bar{\mathbf{H}}_{aperture}^-\right)_{w. av}$$

$$\bar{H}_{aperture}^+ = -\mathbf{H}_i^+ \hat{l}_i^+ , \quad \mathbf{H}_i^+ = -\left(\hat{l}_i^+ \cdot \bar{\mathbf{H}}_{aperture}^+\right)_{w. av}$$

Note that the weighted average is a simple average in this case since we are dealing with the uniform distribution mode $W(s) = 1$. The electromagnetic field under the patch can be written in a transverse electromagnetic configuration (E_z^i, H_y^i) for the i -th patch

$$E_z^{TL} = A_1^i \cos kx + B_1^i \sin kx \quad \text{for } 0 < x < L_i^-$$

$$H_y^{TL} = \frac{jA_1^i}{\eta} \sin kx + \frac{B_1^i}{j\eta} \cos kx$$

and

$$E_z^{TL} = A_2^i \cos k(L-x) + B_2^i \sin k(L-x) \quad \text{for } L_i^- < x < L_i^+$$

$$H_y^{TL} = \frac{A_2^i}{j\eta_e} \sin k(L-x) + \frac{B_2^i}{\eta_e} \cos k(L-x)$$

The arbitrary constants can be identified from the aperture quantities, and the non-ideal TEM-transmission is compensated for by effective propagation constants k_e and η_e . That is

$$E_z^{TL} = E_i^- \cos (k_e x) + j\eta_e H_i^- \sin (k_e x) \quad \text{for } 0 < x < L_i^-$$

$$H_y^{TL} = j \frac{E_i^-}{\eta_e} \sin (k_e x) + H_i^- \cos (k_e x)$$

and

$$E_z^{TL} = E_i^+ \cos [k_e(L-x)] + j \eta_e H_i^+ \sin [k_e(L-x)] \quad \text{for } L_i^- \leq x \leq L_i^+$$

$$H_y^{TL} = \frac{E_i^+}{j \eta_e} \sin [k_e k(L-x)] - H_i^+ \cos [k_e(L-x)]$$

where, $k_e = k_i \sqrt{\epsilon_e}$, $\eta_e = \eta_0 / \sqrt{\epsilon_e}$

$$\epsilon_e = \frac{1}{2} \left[(\epsilon_r + 1) + (\epsilon_r - 1)(1 + \frac{12h}{W})^{-\frac{1}{2}} \right], \text{ see reference (1).}$$

The above two sets of equations must also satisfy the continuity of the electric field at $x = L_i^-$, and that the discontinuity of the magnetic field $[H_y^{TL}(x = L_i^- + \delta) - H_y^{TL}(x = L_i^- - \delta)]$ has to equal the average excitation current in the z-direction (I_z^i/W) where W represents the width of the patch in the y-direction. That is,

$$\begin{aligned} & \left(H_i^+ \sin(k_e L_i^+) - H_i^- \sin(k_e L_i^-) \right) \\ & - \frac{j}{\eta_e} \left(E_i^+ \cos(k_e L_i^+) - E_i^- \cos(k_e L_i^-) \right) = 0 \end{aligned} \quad (33)$$

$$\begin{aligned} & \left(H_i^+ \cos(k_e L_i^+) + H_i^- \cos(k_e L_i^-) \right) \\ & + \frac{j}{\eta_e} \left(E_i^+ \sin(k_e L_i^+) + E_i^- \sin(k_e L_i^-) \right) = - \frac{I_z^i}{W} \end{aligned} \quad (34)$$

Although equations (33) and (34) are the two physical conditions of the model, and we could introduce the coupling at this stage, they proved to be cumbersome to handle before further simplification. If we multiply (34) by $\sin(k_e L_i^-)$ and (33) by $\cos(k_e L_i^-)$ then add both, a mathematically simpler equation is produced, and if (34) is multiplied by $\sin(k_e L_i^+)$ and (33) by $\cos(k_e L_i^+)$ then by subtraction the other simplified relation is also produced. Hence the physical pair (33), (34) is equivalent to the following mathematical pair,

$$H_i^+ - \frac{j}{\eta_e} \left(\cot(k_e L_i) E_i^+ - \csc(k_o L_i) E_i^- \right) = -\frac{1}{W} \left(\frac{\sin kL_i^+}{\sin kL_i^-} \right) I_z^i \quad (35)$$

$$H_i^- + \frac{j}{\eta_e} \left(\csc(k_e L_i) E_i^+ + \cot(k_o L_i) E_i^- \right) = -\frac{1}{W} \left(\frac{\sin kL_i^+}{\sin kL_i^-} \right) I_z^i \quad (36)$$

These two principal mathematical equations have successfully isolated the dependance of the aperture field on the excitation location to the right hand side, and allowed the loading effect of the rectangular cavity (as in the circular case) to appear explicitly by the two terms $(-j \frac{1}{\eta_e} \cot(k_e L_i))$ and $(j \frac{1}{\eta_e} \csc(k_e L_i))$.

To bring the radiation interaction into the picture we assume a linear proportionality between the aperture magnetic field and all existing aperture electric fields. Since we are dealing with two apertures per patch, an attempt can be made to relate the proportionality to the definition of the normalized radiation admittance used in the circular case, which by the variational expression given earlier by equation (17) provides.

$$H_1^+ = \frac{1}{\eta_o} \sum (p_{ij}^{++} E_j^+ + p_{ij}^{+-} E_j^-)$$

$$H_1^- = \frac{1}{\eta_o} \sum (p_{ij}^{-+} E_j^+ + p_{ij}^{--} E_j^-)$$

Once the array configuration becomes like Fig. (5) and no relative rotation between patches are recognized, then the simple reduction of the above four matrix elements to only two basic matrix elements, gives

$$p_{ij}^{--} = p_{ij}^{++} = p_{ij}$$

$$p_{ij}^+ = p_{ji}^- = q_{ij}$$

$$H_i^+ = \frac{1}{\eta_0} \sum (p_{ij} E_j^+ + q_{ji} E_j^-)$$

$$H_i^- = \frac{1}{\eta_0} \sum (q_{ij} E_j^+ + p_{ij} E_j^-)$$

where p_{ij} is the mutual normalized admittance between any two apertures, of the same polarity, in the i -th and the j -th patches. It obeys formula (17) and takes the explicit expression

$$p_{ij} = \frac{j(\frac{h}{w})}{2\pi} \int_{-\frac{w}{2}}^{\frac{w}{2}} \int_{-\frac{w}{2}}^{\frac{w}{2}} dy dy' \frac{e^{-jk_0\rho}}{k_0 \rho^3} \left[k_0^2 (x_j - x_i)^2 + (1+jk_0\rho) \left(2 - 3 \frac{(x_j - x_i)^2}{\rho^2} \right) \right] \quad (37)$$

$$\rho = \sqrt{(x_j - x_i)^2 + (y_j - y_i + y' - y)^2}$$

where (x_i, y_i) is the coordinate of the i -th patch centre and (y, y') are the dummy variables of integration in the longitudinal aperture direction.

Note that this matrix is symmetric ($p_{ij} = p_{ji}$). On the other hand, the q_{ij} represents the mutual admittance matrix element between the negative aperture of the i -th cell ($x_i - \frac{L}{2}, y_i + y$) and the positive

aperture of the j -th patch ($x_j + \frac{L}{2}, y_j + y'$). The opposite polarity dictates that $\hat{\lambda}_i \cdot \hat{\lambda}_j = -1$.

$$q_{ij} = - \frac{j(\frac{h}{w})}{2\pi} \int_{-w/2}^{w/2} \int_{-\frac{w}{2}}^{\frac{w}{2}} dy dy' \left[\frac{e^{-jk_0 p}}{k_0 p^3} \left[k_0^2 (x_j - x_i + L)^2 + (1+jk_0 p) \left(2 - 3 \frac{(x_j - x_i + L)^2}{p^2} \right) \right] \right] \quad (38)$$

$$p = \sqrt{(x_j - x_i + L)^2 + (y_j - y_i + y' - y)^2}, \quad L_i = L_j = L, \quad q_{ij} \neq q_{ji}$$

When the mutual interaction p_{ij}^{+-} between the apertures $(x_i + \frac{L}{2}, y_i + y)$ and $(x_j - \frac{L}{2}, y_j + y')$ is constructed, the formulation becomes the same as equation (38) with the role of (x_i, x_j) interchanged ($p_{ij}^{+-} = q_{ji}$). Hence, the substitution of the invoked linearity in equation (35), (36) for all i's ($i = 1, 2, \dots, N$) produces the two matrix equations

$$\mathbf{P} \mathbf{E}^+ + \mathbf{Q}_t \mathbf{E}^- = - \frac{\eta_0}{w} \mathbf{S}^- \mathbf{I}_z \quad (39)$$

$$\mathbf{Q} \mathbf{E}^+ + \mathbf{P} \mathbf{E}^- = - \frac{\eta_0}{w} \mathbf{S}^+ \mathbf{I}_z \quad (40)$$

where the P and Q matrices differ only from equations (37), (38) in the diagonal element where the cavity loading takes place,

$$(\mathbf{P})_{ii} = (p)_{ii} - j \frac{\eta_0}{\eta_e} \cot(k_e L_i), \quad (\mathbf{P})_{ij} = (p)_{ij} \quad i \neq j$$

$$(\mathbf{Q})_{ii} = (q)_{ii} + j \frac{\eta_0}{\eta_e} \csc(k_e L_i), \quad (\mathbf{Q})_{ij} = (q)_{ij} \quad i \neq j$$

$$(\mathbf{S}^+)_{ij} = \delta_{ij} [\sin(k_e L_i^+)/\sin(k_e L_i)], \quad (\mathbf{S}^-)_{ij} = \delta_{ij} [\sin(k_e L_i^-)/\sin(k_e L_i)]$$

Equations (39) and (40) are the two coupled equations between the aperture fields and the decoupling yields

$$(\mathbf{P} - \mathbf{Q}_t \mathbf{P}^{-1} \mathbf{Q}) \mathbf{E}^+ = - \frac{\eta_0}{w} (\mathbf{S}^- - \mathbf{Q}_t \mathbf{P}^{-1} \mathbf{S}^+) \mathbf{I}_z \quad (41)$$

$$(\mathbf{P} - \mathbf{Q} \mathbf{P}^{-1} \mathbf{Q}_t) \mathbf{E}^- = - \frac{\eta_0}{w} (\mathbf{S}^+ - \mathbf{Q} \mathbf{P}^{-1} \mathbf{S}^-) \mathbf{I}_z \quad (42)$$

These final matrix relations are the ones to be handled with the aid of a computer program in order to obtain the aperture electric field dependence on the excitation current with the inclusion of the mutual interaction. These equations can be converted to the standard electric circuit parameters (edge voltages, input currents, and admittance/impedance matrices) to look like,

$$\mathbf{V}_e^+ = \mathbf{Z}_e^+ \mathbf{I} = (\mathbf{Y}_s - \mathbf{Y}_{pn} \mathbf{Y}_s^{-1} \mathbf{Y}_{np})^{-1} (\mathbf{S}^- - \mathbf{Y}_{pn} \mathbf{Y}_s^{-1} \mathbf{S}^+) \mathbf{I} \quad (43)$$

$$\mathbf{V}_e^- = \mathbf{Z}_e^- \mathbf{I} = (\mathbf{Y}_s - \mathbf{Y}_{np} \mathbf{Y}_s^{-1} \mathbf{Y}_{pn})^{-1} (\mathbf{S}^+ - \mathbf{Y}_{np} \mathbf{Y}_s^{-1} \mathbf{S}^-) \mathbf{I} \quad (44)$$

$$\mathbf{Y}_s = \frac{(\frac{w}{h})}{\eta_0} \mathbf{P} \quad , \quad \mathbf{Y}_{np} = \frac{(\frac{w}{h})}{\eta_0} \mathbf{Q} \quad , \quad \mathbf{Y}_{pn} = \mathbf{Y}_{np}'$$

$$\begin{aligned} (\mathbf{P})_{ii} = \eta_0 \left(\frac{h}{w} \right) & \left[\left(\frac{(\frac{w}{\lambda_0})^2}{15} \sum \frac{(-1)^m (k_0 w)^{2m}}{(2m+3)! (2m+1)} \right) \right. \\ & \left. + j \left(\frac{1}{60} \left(\frac{w}{\lambda_0} \right) \left(\frac{\Delta l}{h} \right) \epsilon_e - \frac{w}{\eta_e h} \cot(k_e L_i) \right) \right] \end{aligned} \quad (45)$$

$$\epsilon_e = \frac{1}{2} [(\epsilon_r + 1) + (\epsilon_r - 1) (1 + 12 \frac{h}{w})^{-\frac{1}{2}}]$$

$$\frac{\Delta l}{h} = .412 \left[\frac{(\epsilon_e + 3) (\frac{w}{h} + .264)}{(\epsilon_e - .258) (\frac{w}{h} + .8)} \right]$$

Note that the real part has been accomplished exactly while the imaginary part of the self-admittance of the aperture has been estimated with the extention formula supplied by reference (1). That is to say, the self-admittance of a radiating aperture becomes:

$$Y_a = G_a + j B_a = \left(\frac{\left(\frac{w}{\lambda_o}\right)^2}{15} \sum \frac{(-1)^m (k_o w)^{2m}}{(2m+3)! (2m+1)} \right) + j \left(\frac{1}{60} \left(\frac{w}{\lambda_o}\right) \left(\frac{\Delta l}{h}\right) \epsilon_e \right) \quad (46)$$

where the imaginary part is the same as reference [1], equation (2.30b) page 46, and real part gives the actual value of the radiation conductance that in the limit of $w \ll \lambda_o$ coincides with equation (2.50) of the reference mentioned. As a check on the mathematical model, we considered the special case of one patch without mutual coupling between the negative and positive edges $[Y_{np}]_{11} = j \frac{W}{\eta_e h} \csc(k_e L)$ and injected the input current at the left edge $[S^+ = 1, S^- = 0]$. The corresponding expression, using equation (44), for the input admittance becomes

$$Y_{in} = Y_a + \frac{w}{\eta_e h} \frac{y_a + j \frac{w}{\eta_e h} \tan(k_e L)}{\frac{w}{\eta_e h} + j Y_a \tan(k_e L)}$$

which is identical to reference [1], equation 2.55, once the characteristic admittance is recognized by $y_o = \frac{W}{\eta_e h}$. As a final check the case of two different patches, arbitrarily translated from each other, are injected both from the negative edge so that $S^+ = 1$, and $S^- = 0$. The (2×2) matrices can be handled analytically and produces the final input admittance parameters,

$$Y_{11} = Y_{ST1} - \left(Y_{MT1}^2 Y_{ST2} - 2 Y_s^{12} Y_{np}^{12} Y_{MT1} + [Y_{np}^{12}]^2 Y_{ST1} \right) / N$$

$$Y_{11} = Y_{ST2} - \left\{ Y_{MT2}^2 Y_{ST1} - 2 Y_s^{12} Y_{np}^{21} Y_{MT2} + [Y_{np}^{21}]^2 Y_{ST2} \right\} / N$$

$$Y_{12} = Y_{21} = Y_s^{12} - \left\{ Y_{MT1} Y_{ST2} + Y_{MT2} Y_{ST1} - Y_{MT1} Y_{MT2} Y_s^{12} - Y_s^{12} Y_{np}^{12} Y_{np}^{21} \right\} / N$$

$$Y_{ST1} = Y_a^{-1} - j Y_o^{-1} \cot(k_e L_1) , Y_{MT1} = Y_{np}^{-1} + j Y_o^{-1} \csc(k_e L_1)$$

$$Y_{ST2} = Y_a^{-2} - j Y_o^{-2} \cot(k_e L_2) , Y_{MT2} = Y_{np}^{-2} + j Y_o^{-2} \csc(k_e L_2)$$

$$N = Y_{ST1} Y_{ST2} - (Y_s^{12})^2$$

These results are identical to those of Van Lil and Van de Cappelle [20], and we conclude that the model as it stands is correct and ready for numerical programming of the general case.

It is desirable from the feed point of view to know the input impedance at any arbitrary location and not the special case of $L_i^- = 0$ discussed above. This step can be explained by going back to the transmission line equation, where the input electric field E_{in}^i is the same as E_i^{TL} at $x = L_i^-$. That is,

$$E_{in} = E_i^+ \cos(k_e L_i^+) + j \eta_e H_i^+ \sin(k_e L_i^+) \quad (47)$$

or from left of the source

$$E_{in} = E_i^- \cos(k_e L_i^-) + j \eta_e H_i^- \sin(k_e L_i^-) \quad (48)$$

use any of the two equations and substitute the corresponding aperture magnetic field (H_i^+ / H_i^-) from equations (35/36) to obtain the following relation

$$E_{in}^i = \left(\frac{\sin(k_e L_i^-)}{\sin(k_e L_i)} \right) E_i^+ + \left(\frac{\sin(k_e L_i^+)}{\sin(k_e L_i)} \right) E_i^- \\ - j \frac{\eta_e}{w} \sin k_e L_i \left(\frac{\sin(k_e L_i^-)}{\sin(k_e L_i)} \right) \left(\frac{\sin(k_e L_i^+)}{\sin(k_e L_i)} \right) I_z^i$$

which must hold for all $i = 1, 2, \dots, N$ in the array.

Therefore,

$$E_{in} = S^- E^+ + S^+ E^- - j \frac{\eta_e}{w} \sin(k_e L) S^+ S^- I_z \quad (49)$$

Hence, the conversion to the circuit parameters gives

$$V_{in} = Z_{in} I_{in} \quad (50)$$

$$Z_{in} = S^- Z_e^+ + S^+ Z_e^- + j \frac{\eta_e h}{w} \sin(k_e L) S^- S^+ \quad (51)$$

where the edge impedance are defined by equations (43) and (44).

4. NUMERICAL CHECKS OF THE MODELS

Both models clearly exhibit complexity far beyond the capability of just a disc calculator even for the limiting case of only two patch array. The main obstacles are the numerical integration required for the variational form of the mutual interaction and the need to invert complex matrices in order to study the array performance. Therefore, an efficient computer program has been written to manipulate these given formulations toward the different characteristics of interest to the designer. Two prototype programs are written in FORTRAN for both the circular and rectangular planar arrays. While they are still in infancy, some of their results are shown here for comparison with the available data in the literature and to validate the acceptability of the model formulation as given in this report. The case of a single rectangular patch, resonant at 2.2 GHz, with the dimension ($L = 4.14$ cm, $W = 6.858$ cm, $h = 0.159$ cm) and a relative permittivity $\epsilon_r = 2.62$, is injected from a microstrip line at $L = 1.016$ cm is put to the program for the calculation of the input impedance. Since this is a duplicate of a case studied [2] in the literature, to see the effect of including the mutual interaction between the left and right apertures in the transmission line theory, we have the chance to validate our calcualtion. Figure (7) shows the modal input impedance of the TEM_1 -mode in the neighbourhood of its resonance for the case above as compared with a circular disc patch (TM_{11} -mode) that resonates at the same frequency and is fed by a hypothetical strip line at $\phi = 0$. The results seems to be reasonable compared to the available data. Since we evaluate the self-conductance of a radiating aperture analyti-

cally, a test for the variation of the radiation admittance, for different operating modes of the circular disc patch, with frequency are shown in figures [8]-[12] and the result of the TM_{11} mode agrees with the available data [2], [10]. Figures [13]-[18] investigate the variation of the mutual coupling with the separation, for various orientations (E-coupling, H-coupling, and the middle plane) of two circular patches in different modes of operation. Finally figures [19]-[22] show the effect of coupling on the array radiation pattern of a five element array with both E-plane and H-plane coupled configurations. The circular array operates in the TM_{11} -mode and the rectangular case is in the TEM_1 -mode. The result is acceptable and the testing of the programs are completed. A complete listing of the two programs are documented at the end of the report with a short user manual. Most of the parameters used in the program itself are the one used in this report to facilitate any future augmentation and modification by the users.

EXPERIMENTAL VERIFICATION

In the previous section mutual coupling effects in linear arrays was investigated. No experimental verification was provided. To examine the validity of the solutions it is, of course, desirable to compare the computed results with experimental data. In this section, we provide a sample case and compare the computed radiation patterns with measured ones.

For the array, a planar geometry is considered, which is shown in Fig. 23. It consists of 19-circular patches designed with a hexagonal grid. Each element is excited by a coaxial probe, at an appropriate point to match the line impedance. Since, the effect of mutual coupling increases with the beam scan, the array elements were phased to generate a main beam along the 35° elevation. The element phases were achieved by using cables of appropriate lengths. The radiation pattern without the effect of mutual coupling gives a main beam with a peak gain of about 12.5 dB, which is considerably higher than the measured gain of about 10 dB. Including the effects of mutual coupling reduces the gain to that of the measured one. Their agreement is quite satisfactory, except near the second sidelobe, where the measured results give a higher level.

These results indicate that, for the selected array the inclusion of the mutual coupling effect improves the prediction of the array pattern, essentially in all ranges of the θ -angle, except near the zero elevation angle. This is of course expected, since the analytic model is based on the assumption of infinite ground plane and the measurement is for an array with a finite ground plane. This single comparison also validates the accuracy of the array modelling. Further comparison

should be carried out to investigate the accuracy of the cross-polar components and also the effects of higher order modes.

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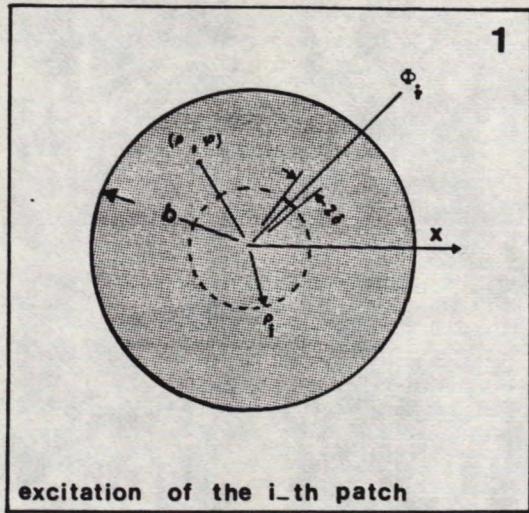


Fig. 1

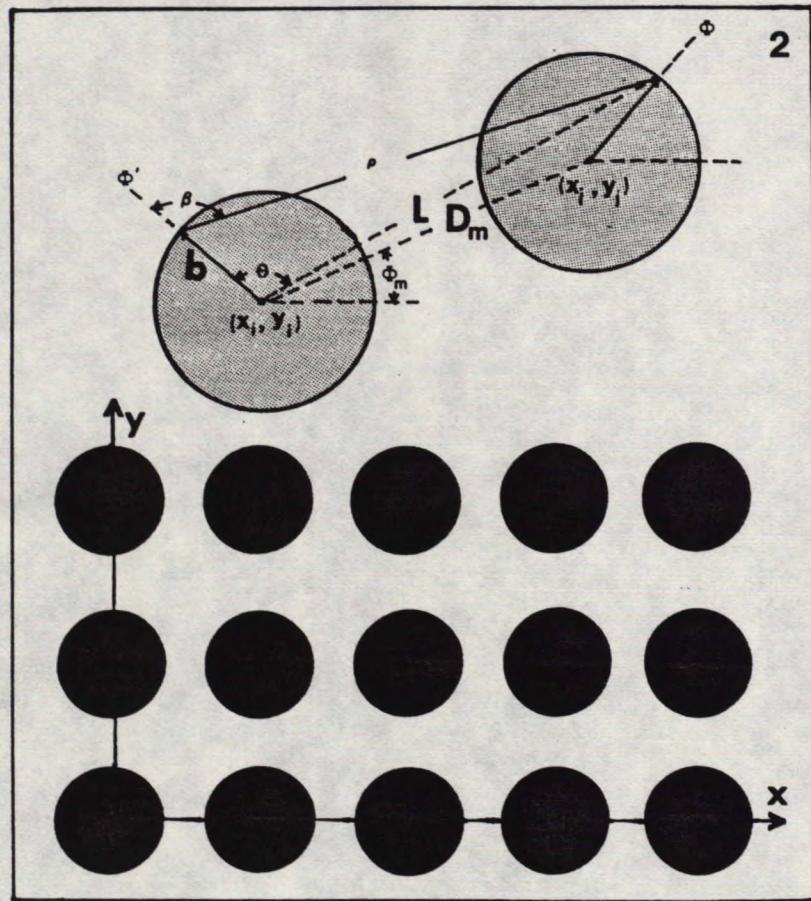


Fig. 2

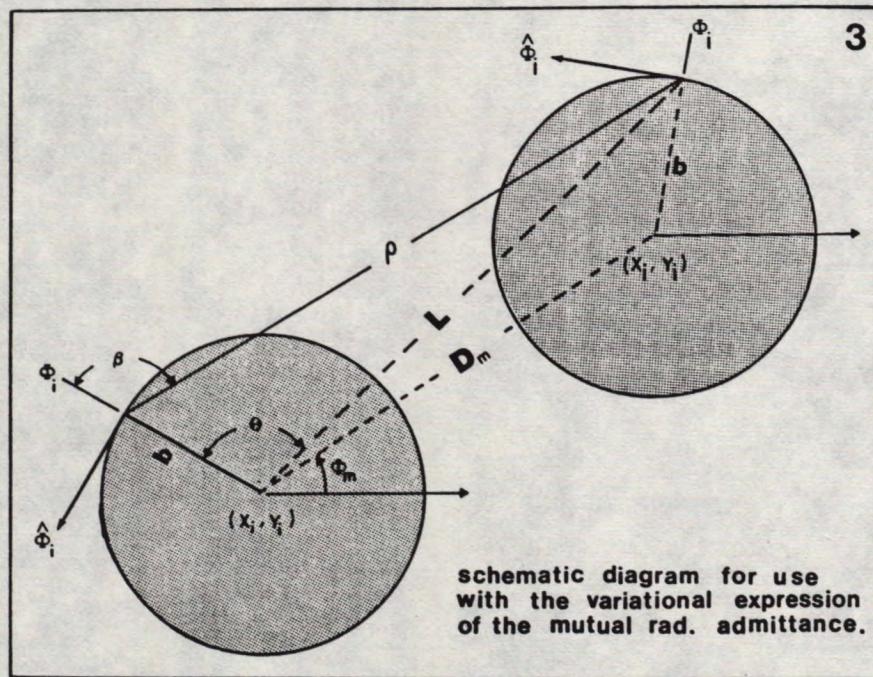


Fig. 3

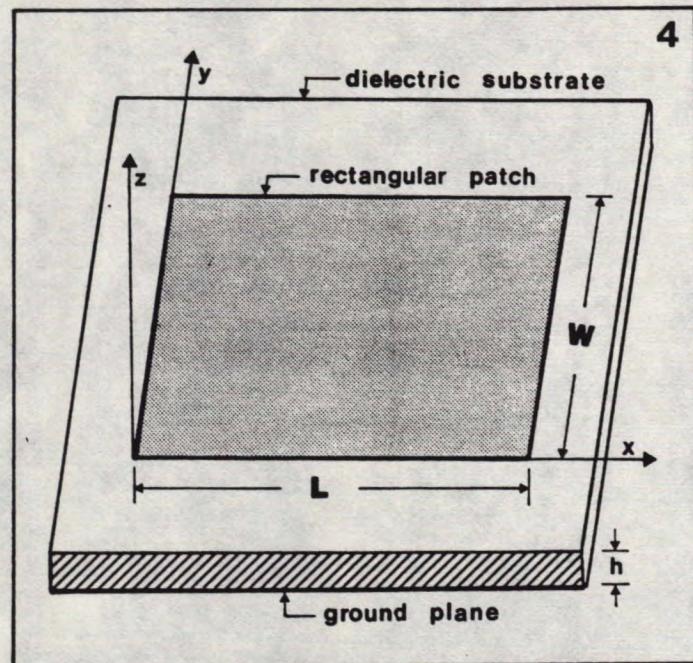


Fig. 4

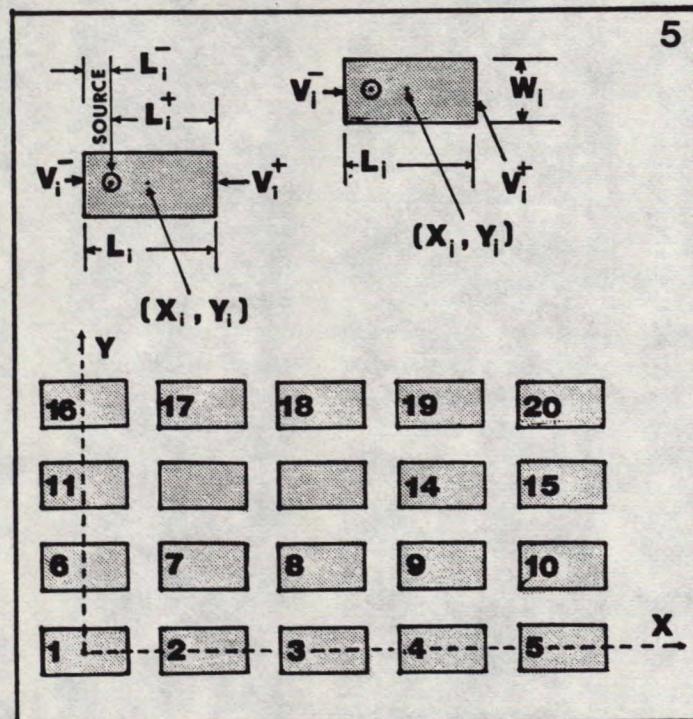
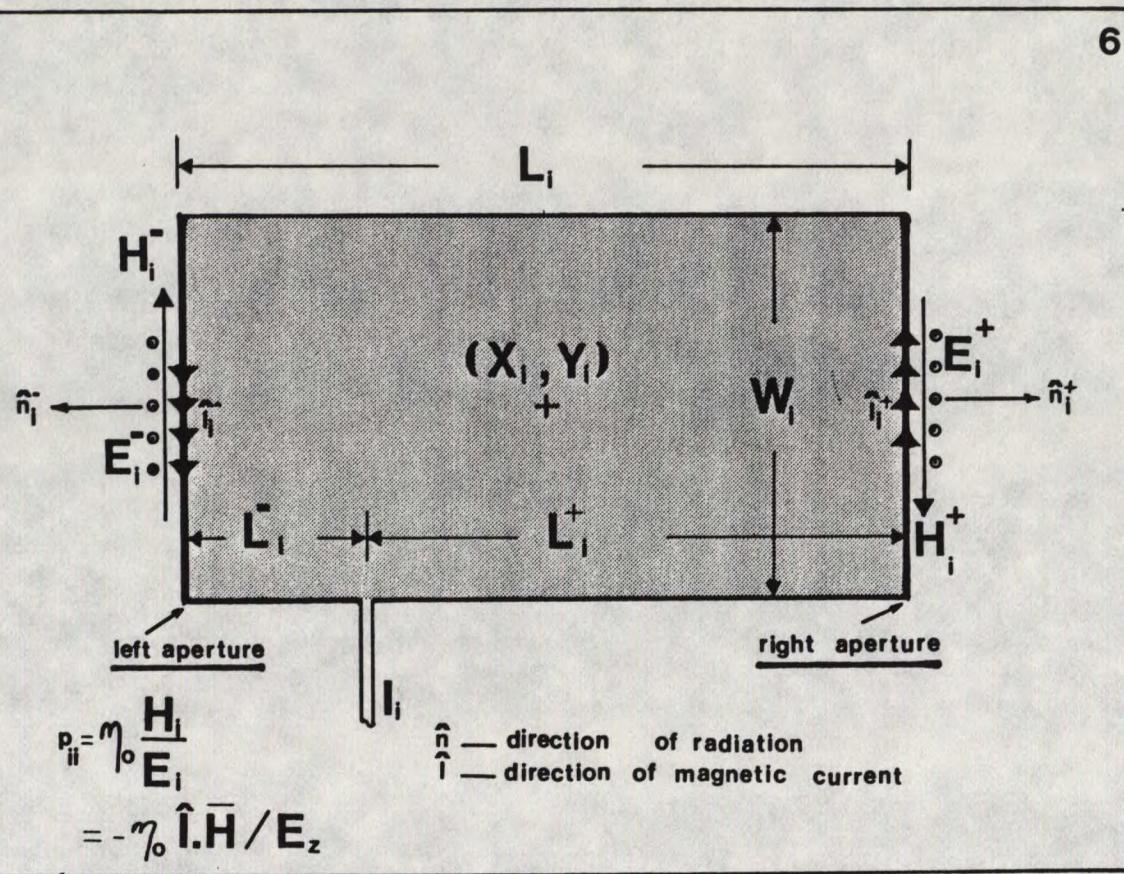


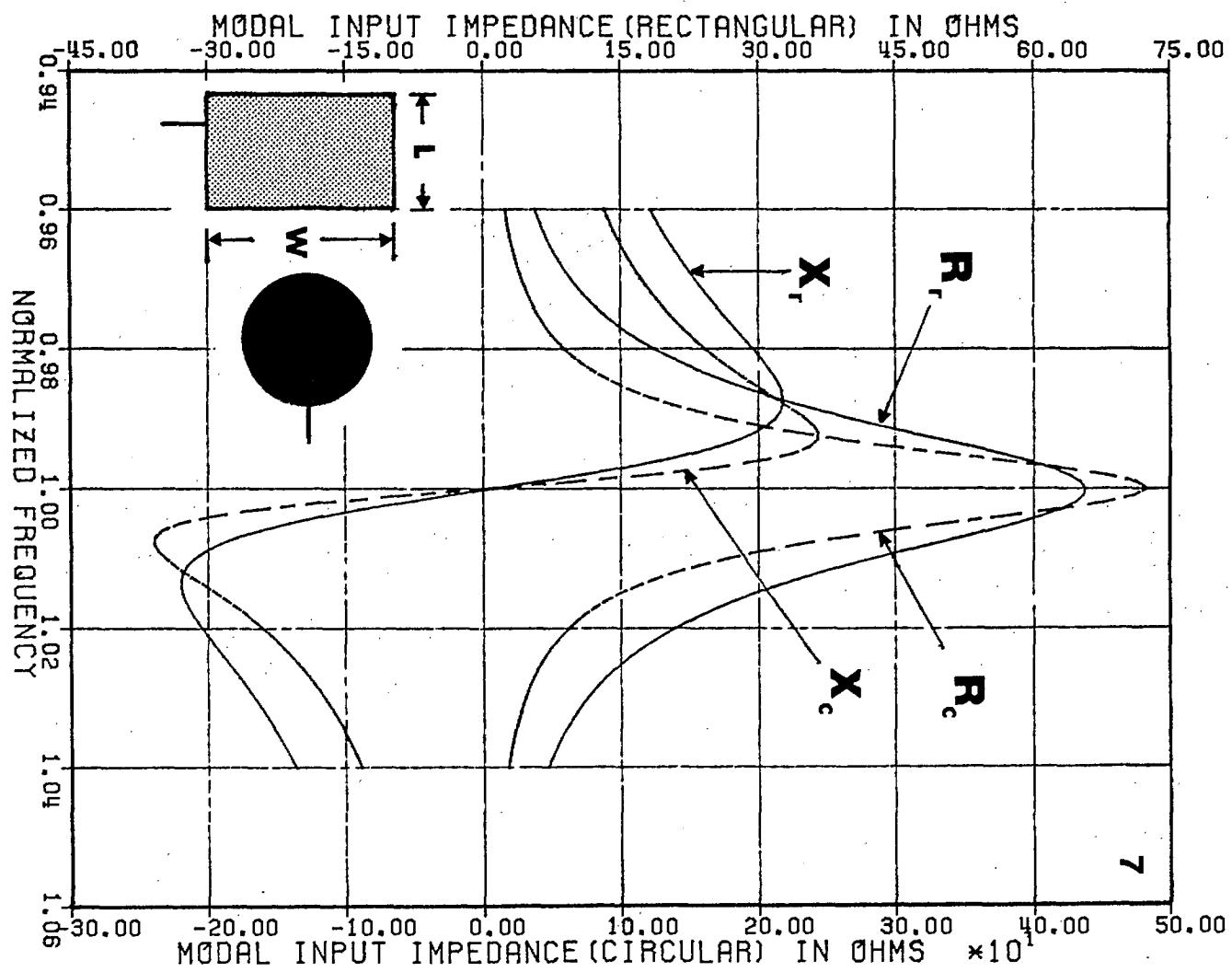
Fig. 5

6

Fig. 6



CELL'S PARAMETERS FOR RECTANGULAR MICROSTRIP PATCH PLANAR ARRAY



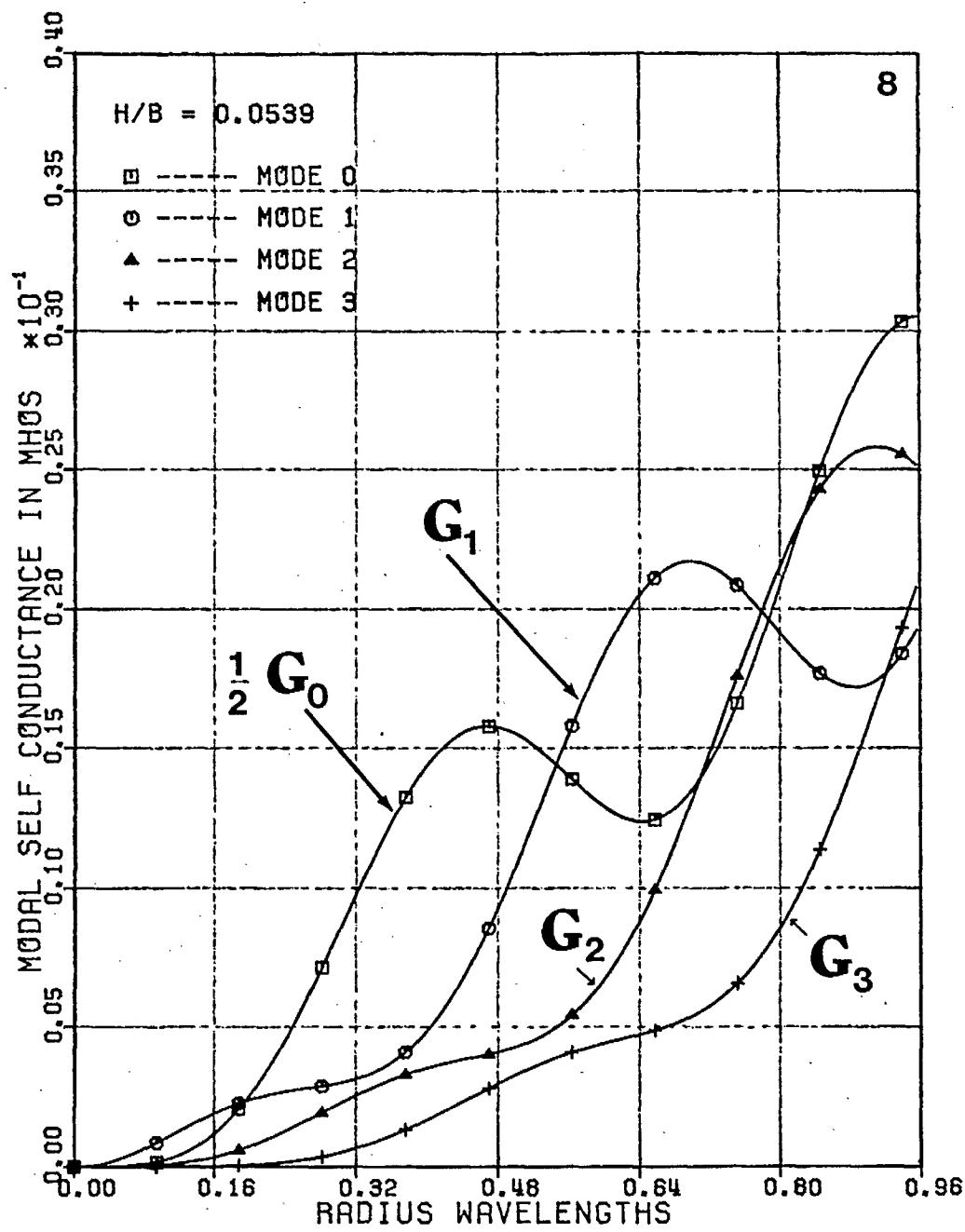


Fig. 8

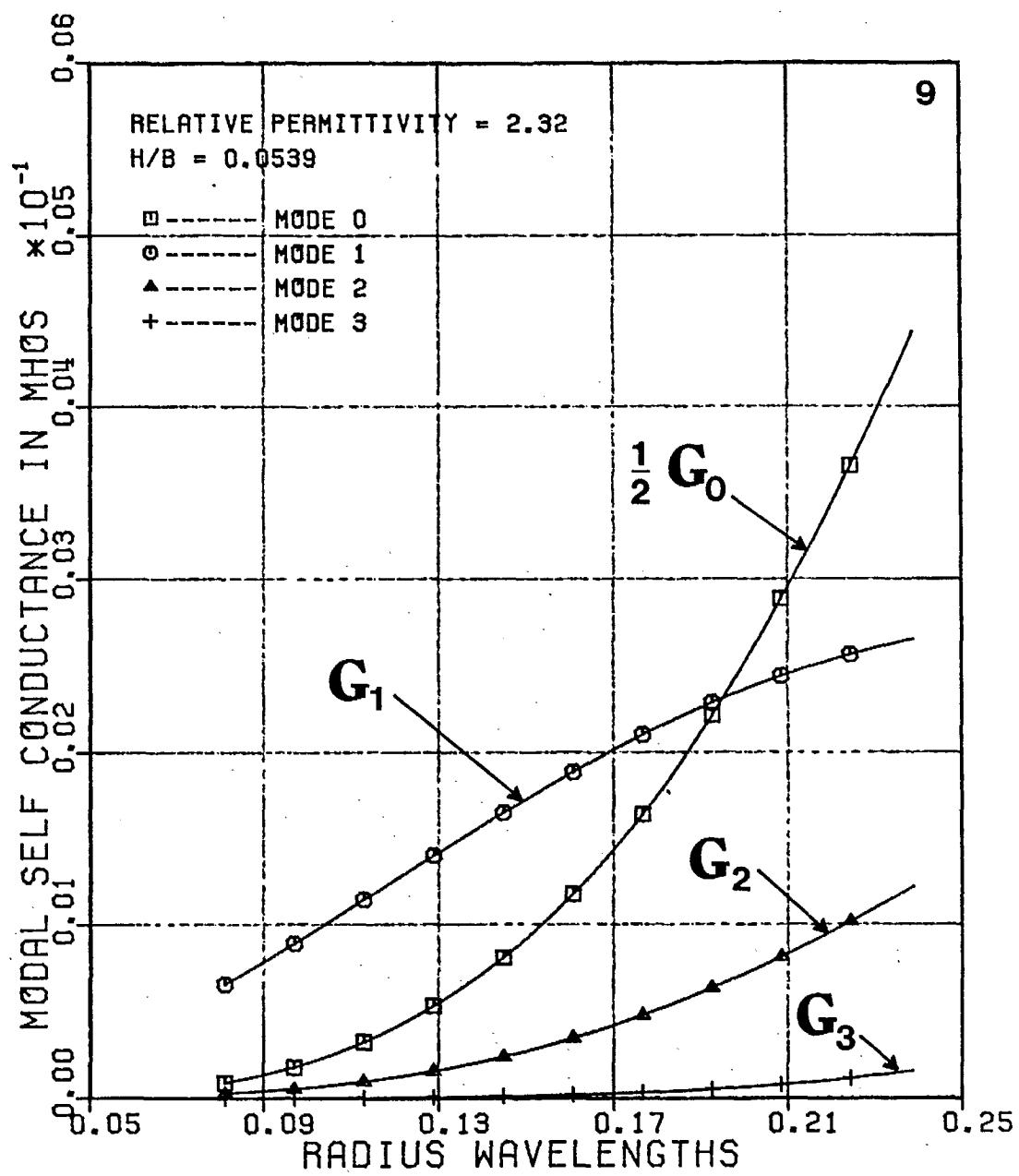


Fig. 9

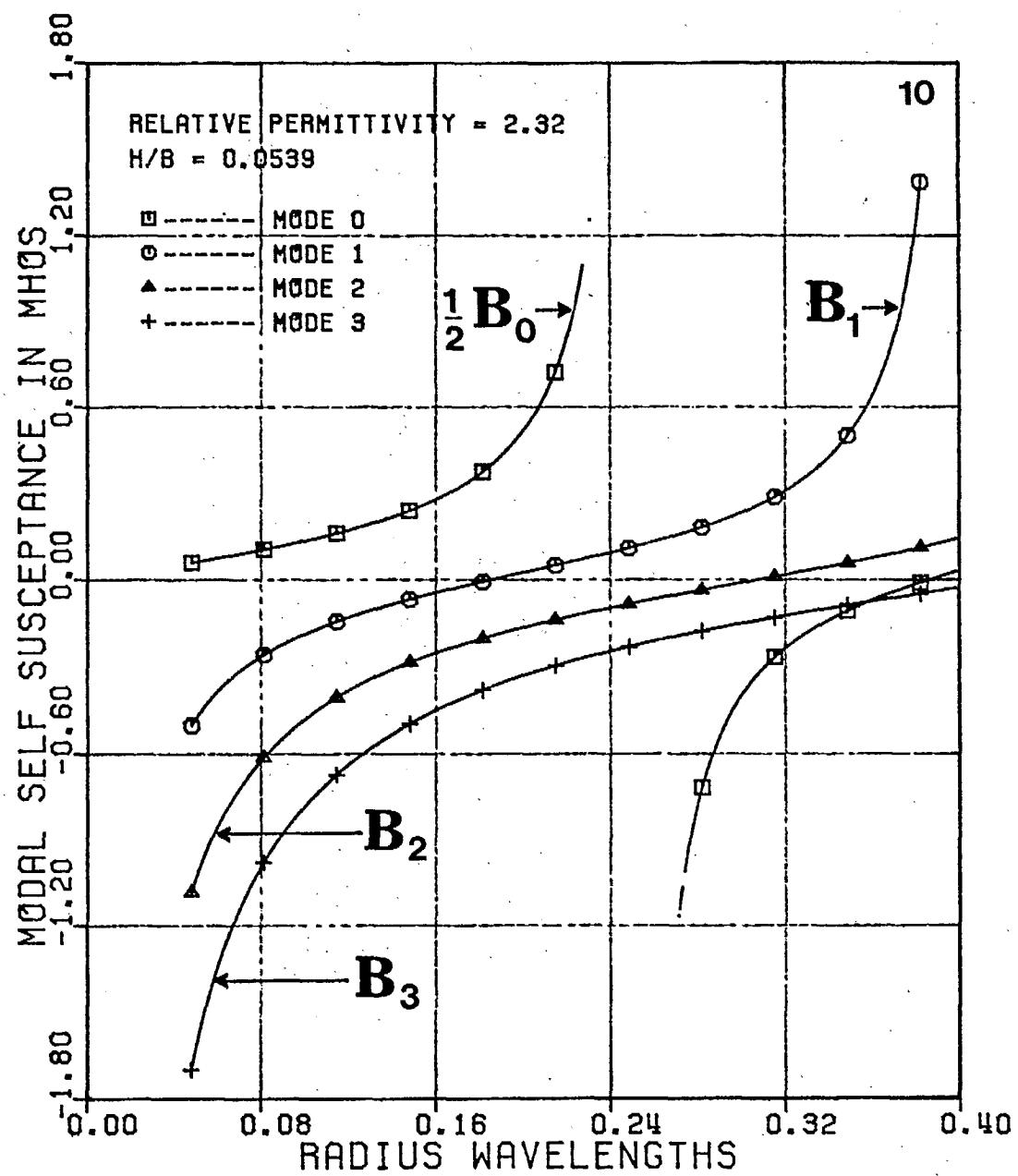


Fig. 10

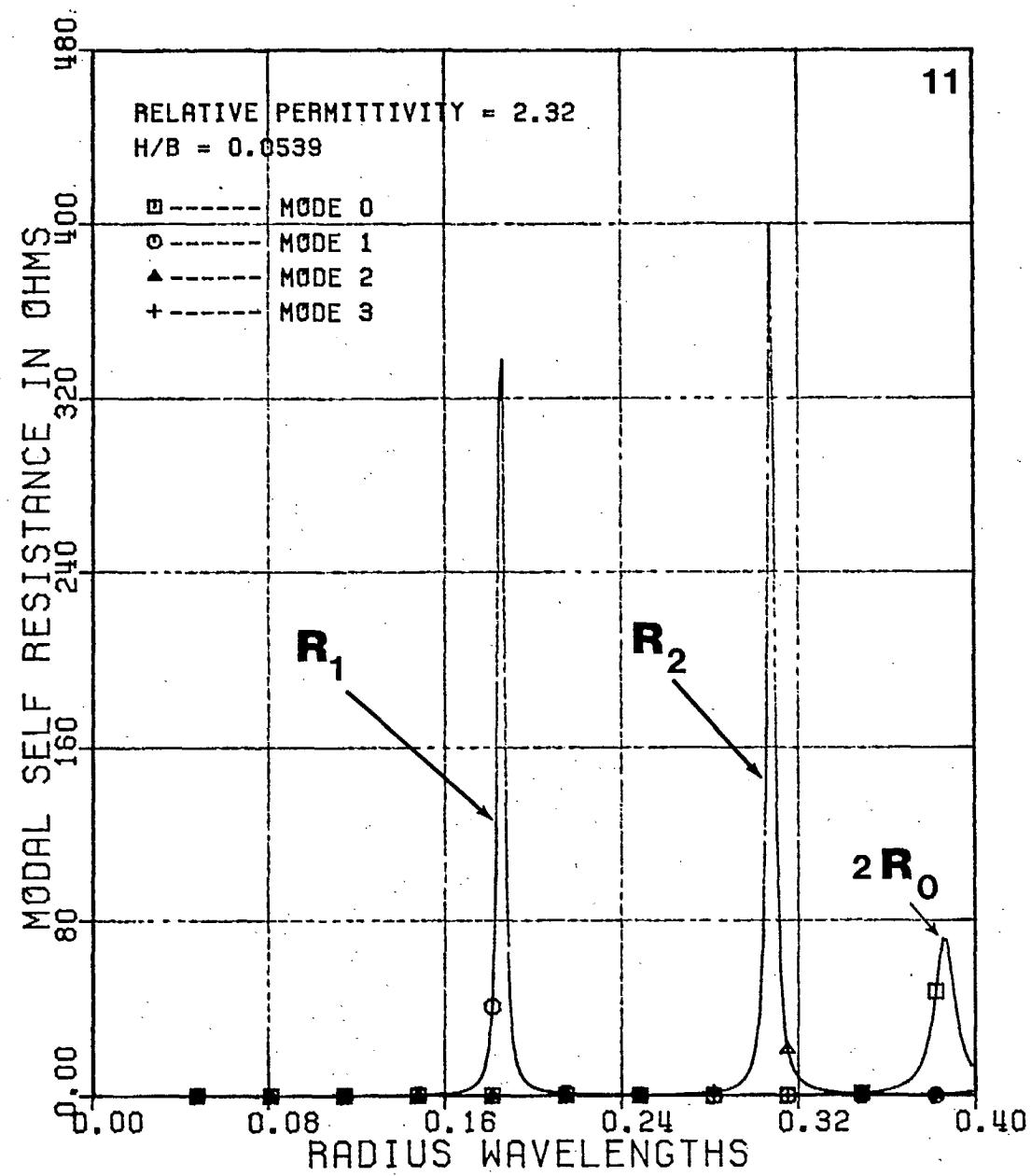


Fig.II

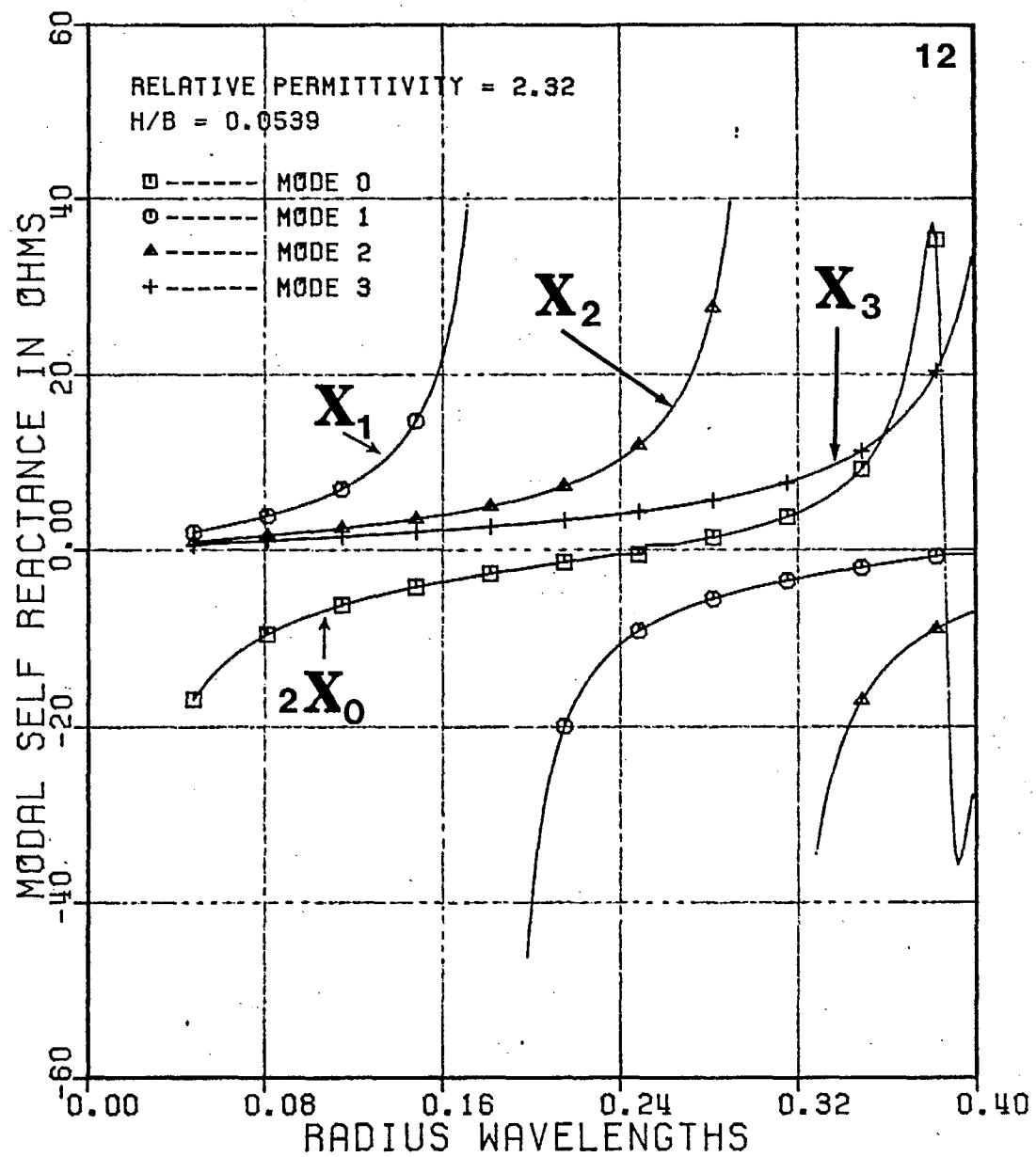


Fig. 12

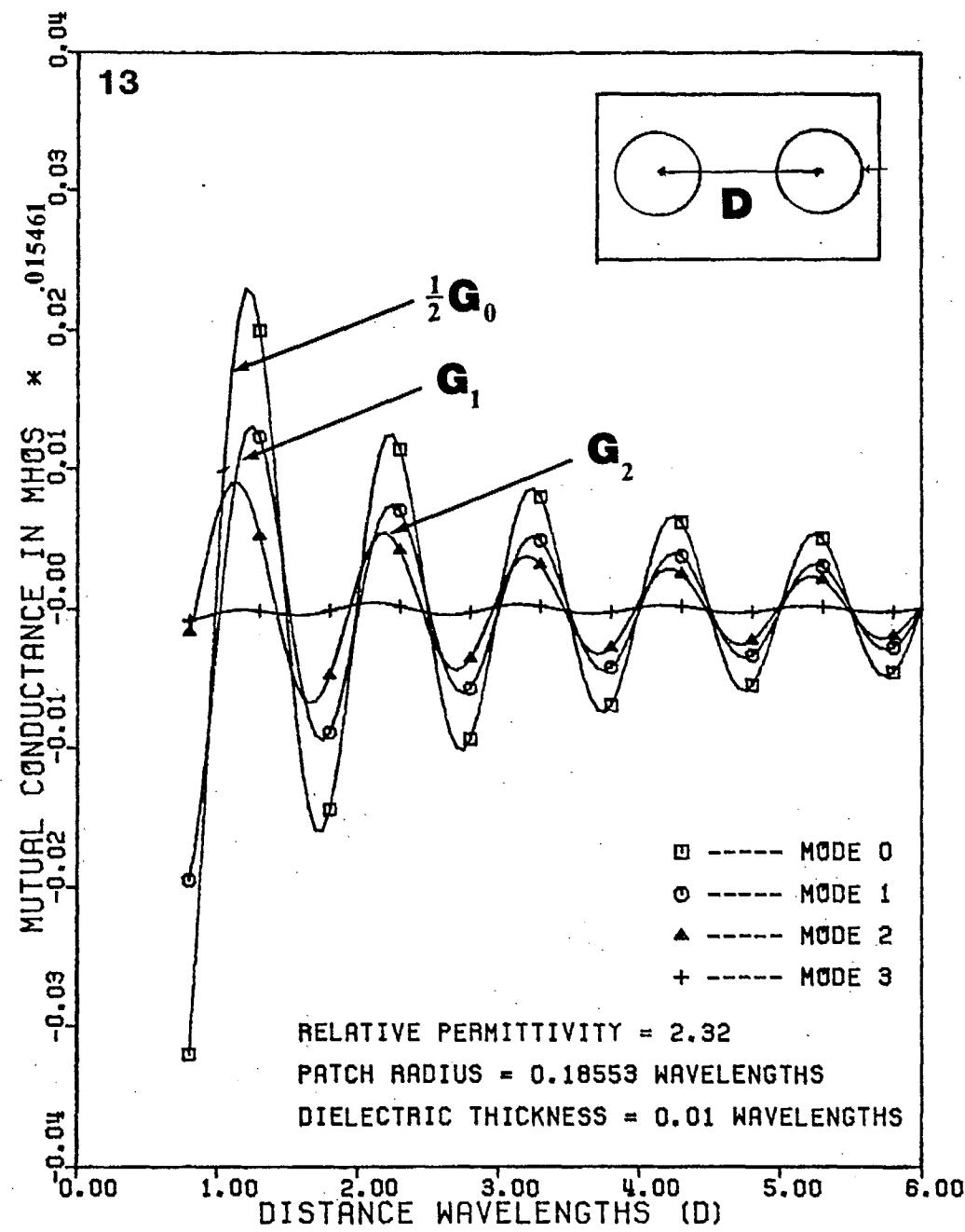


Fig. 13

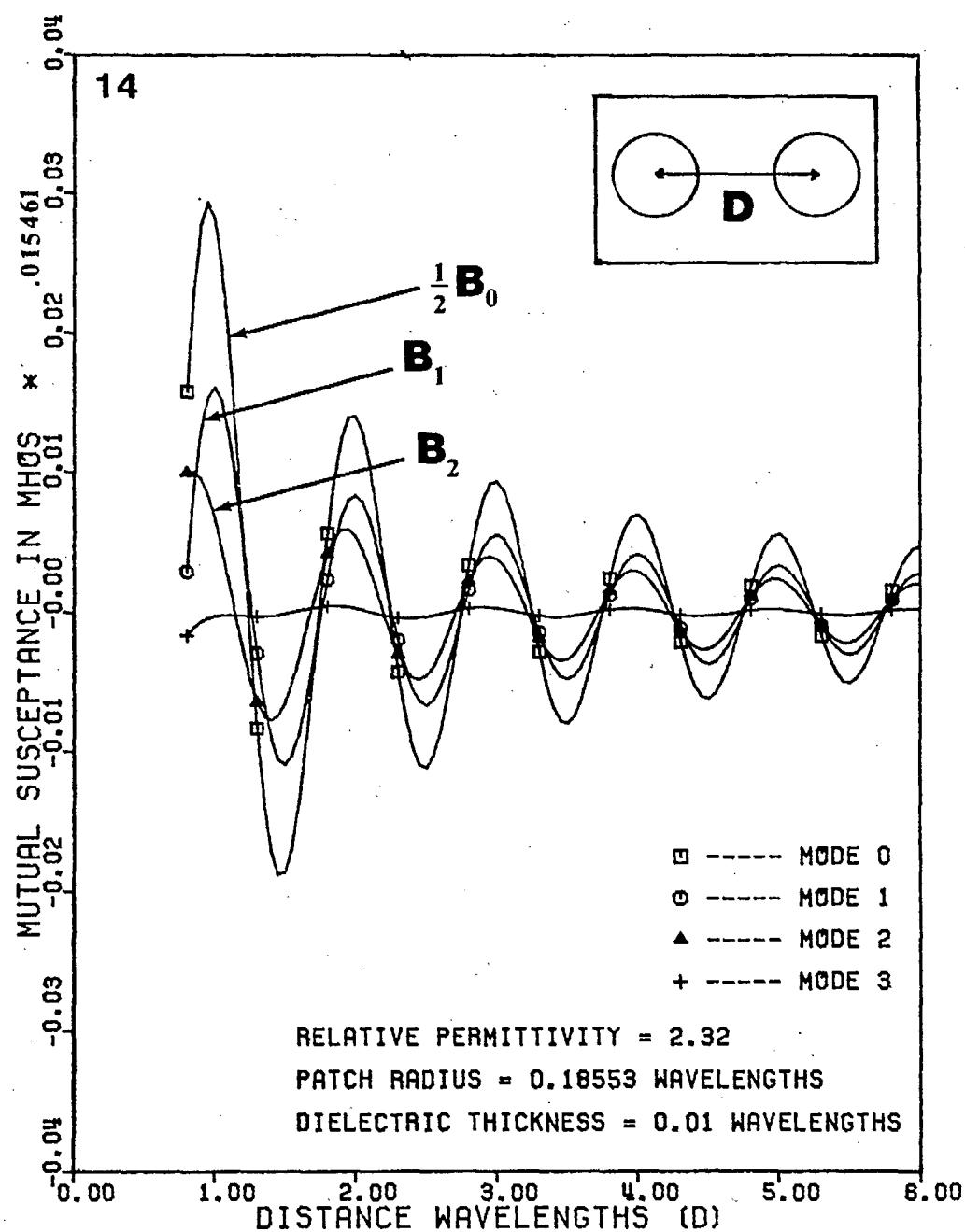


Fig. 14

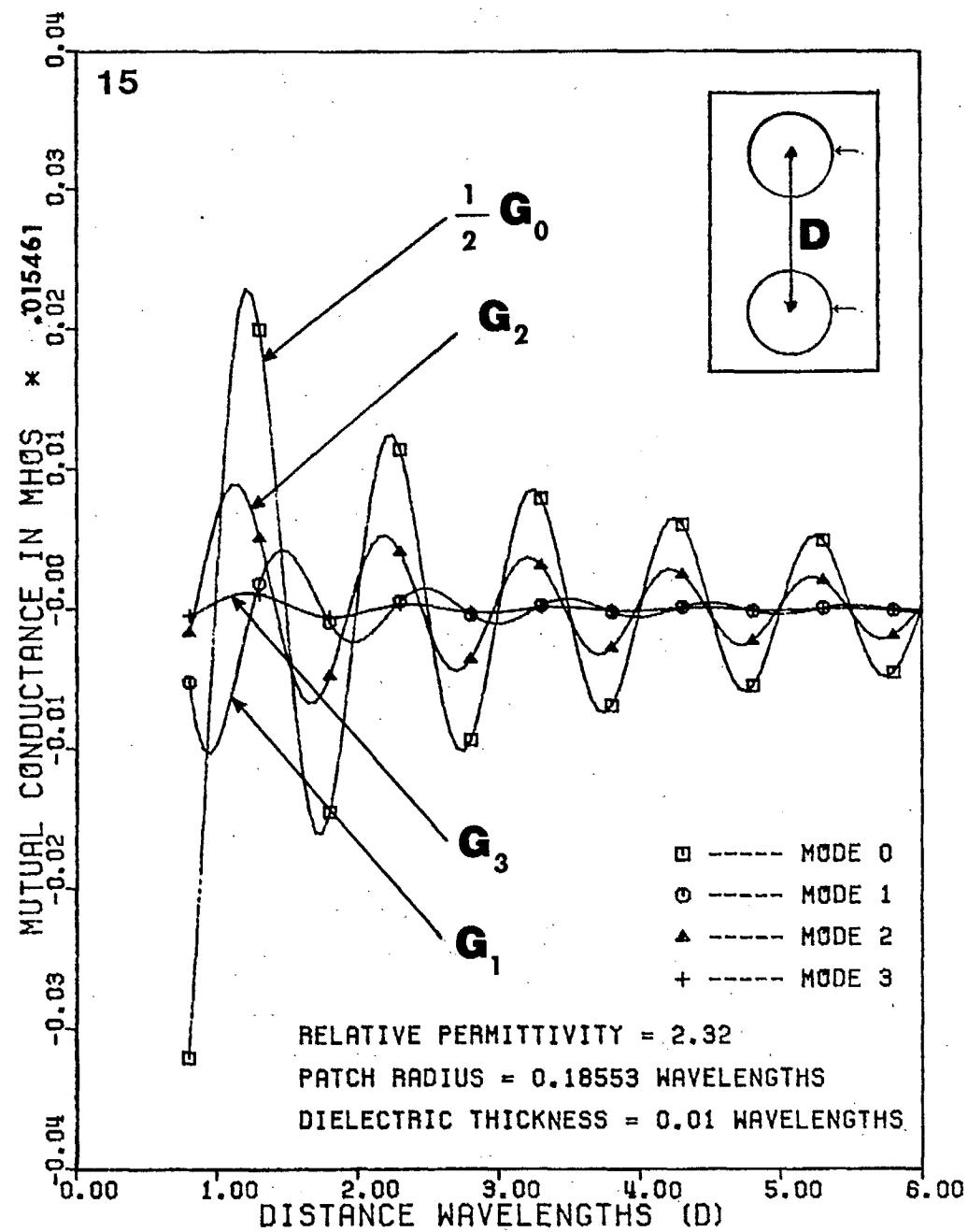


Fig. 15

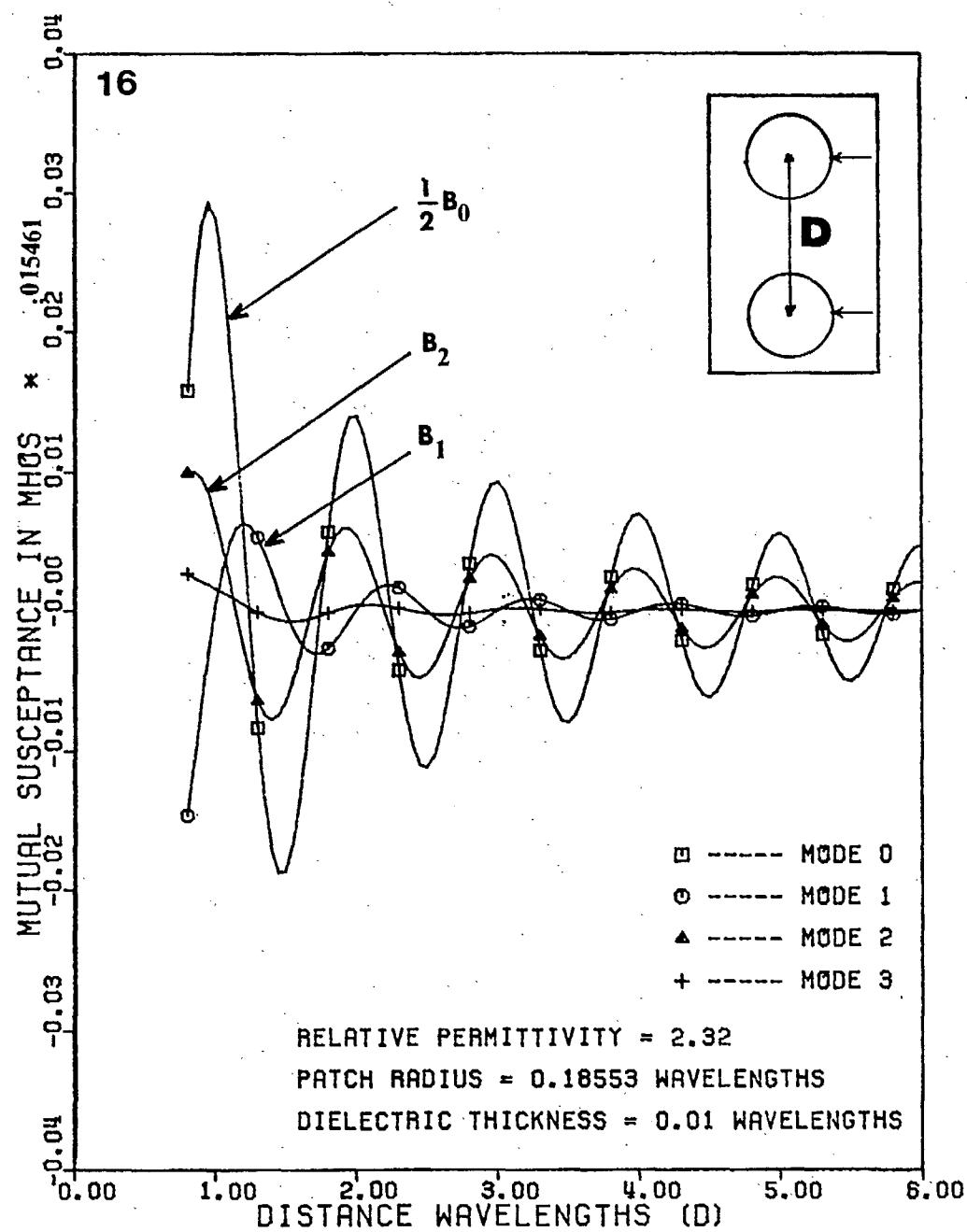


Fig. 16

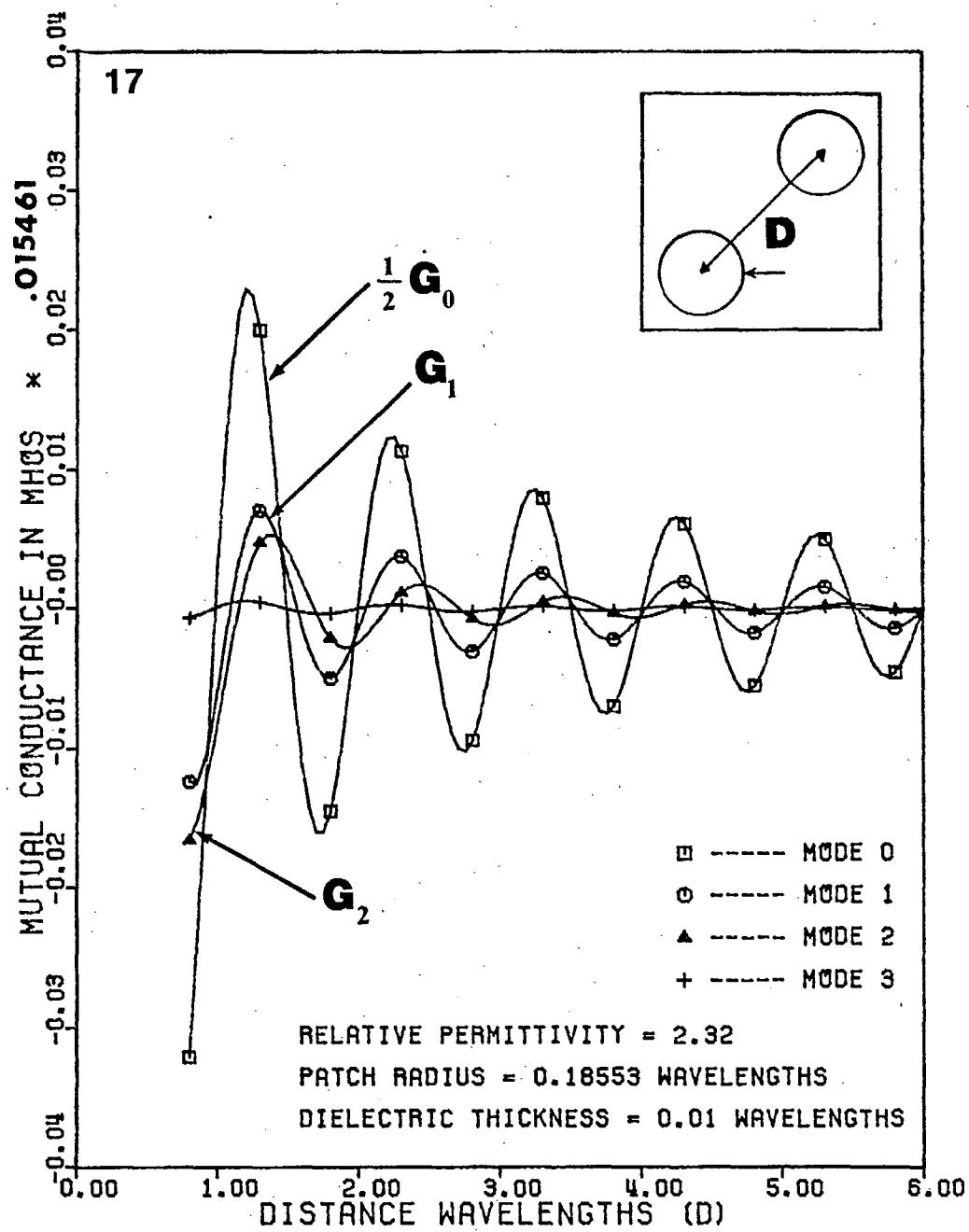


Fig. 17

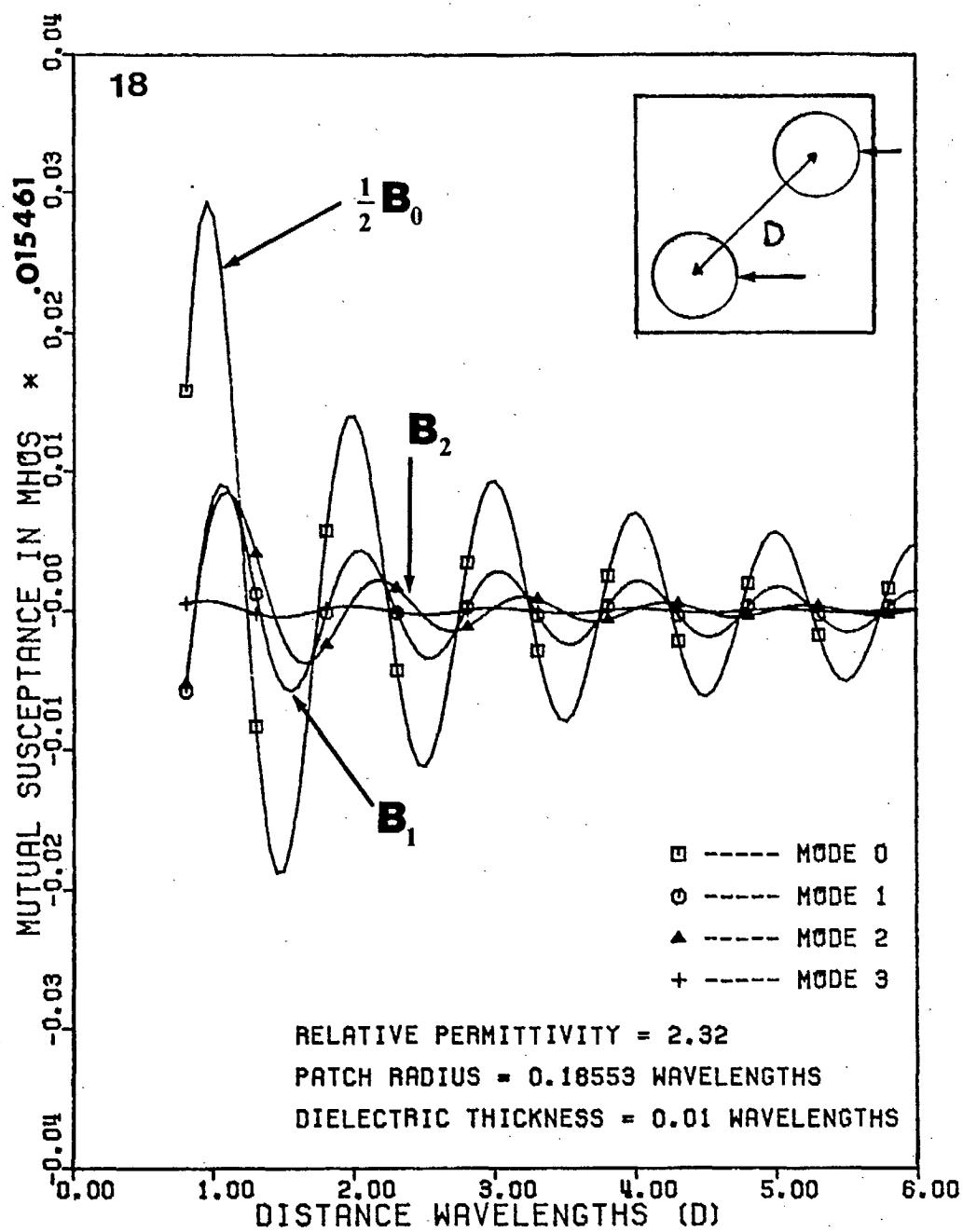
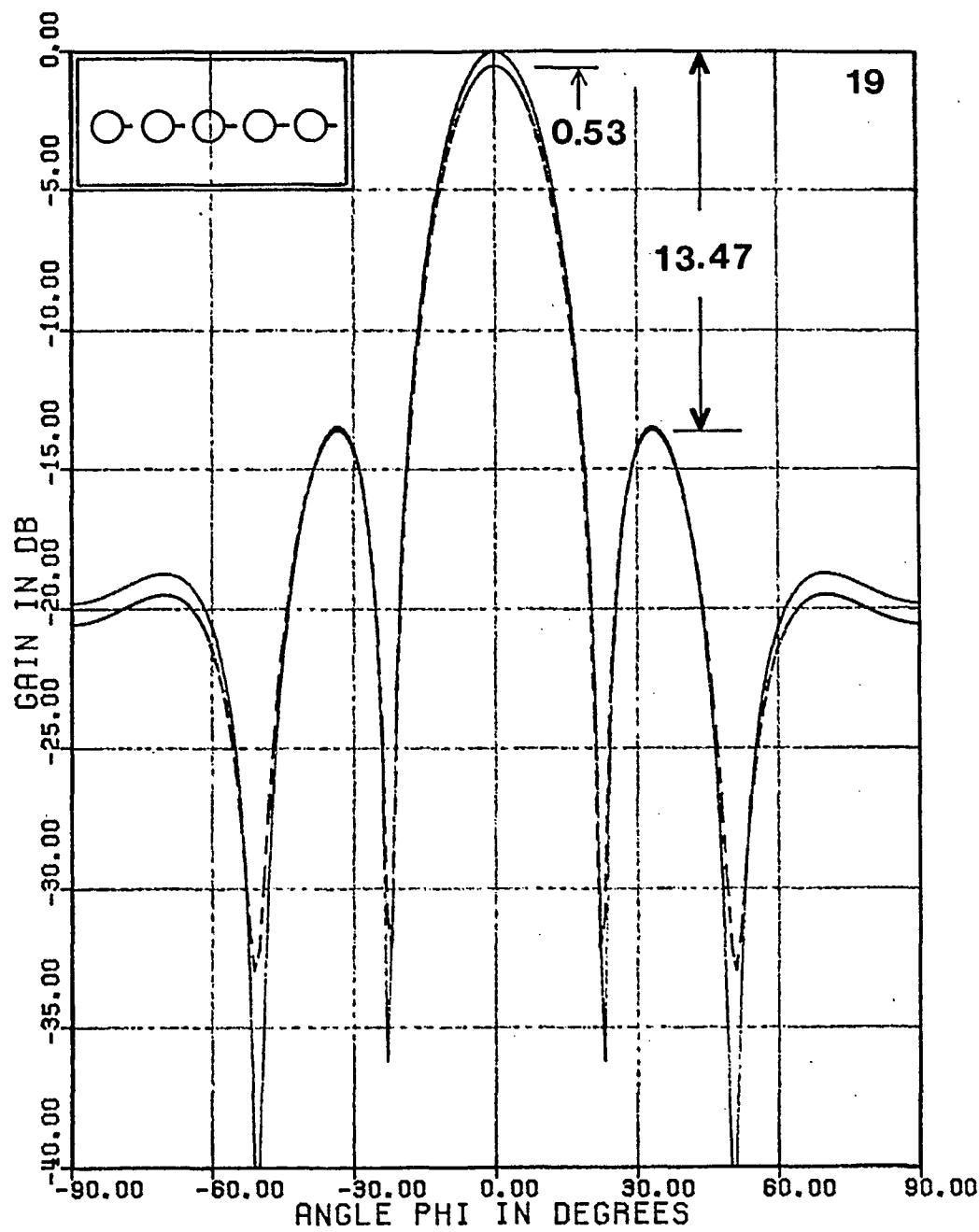
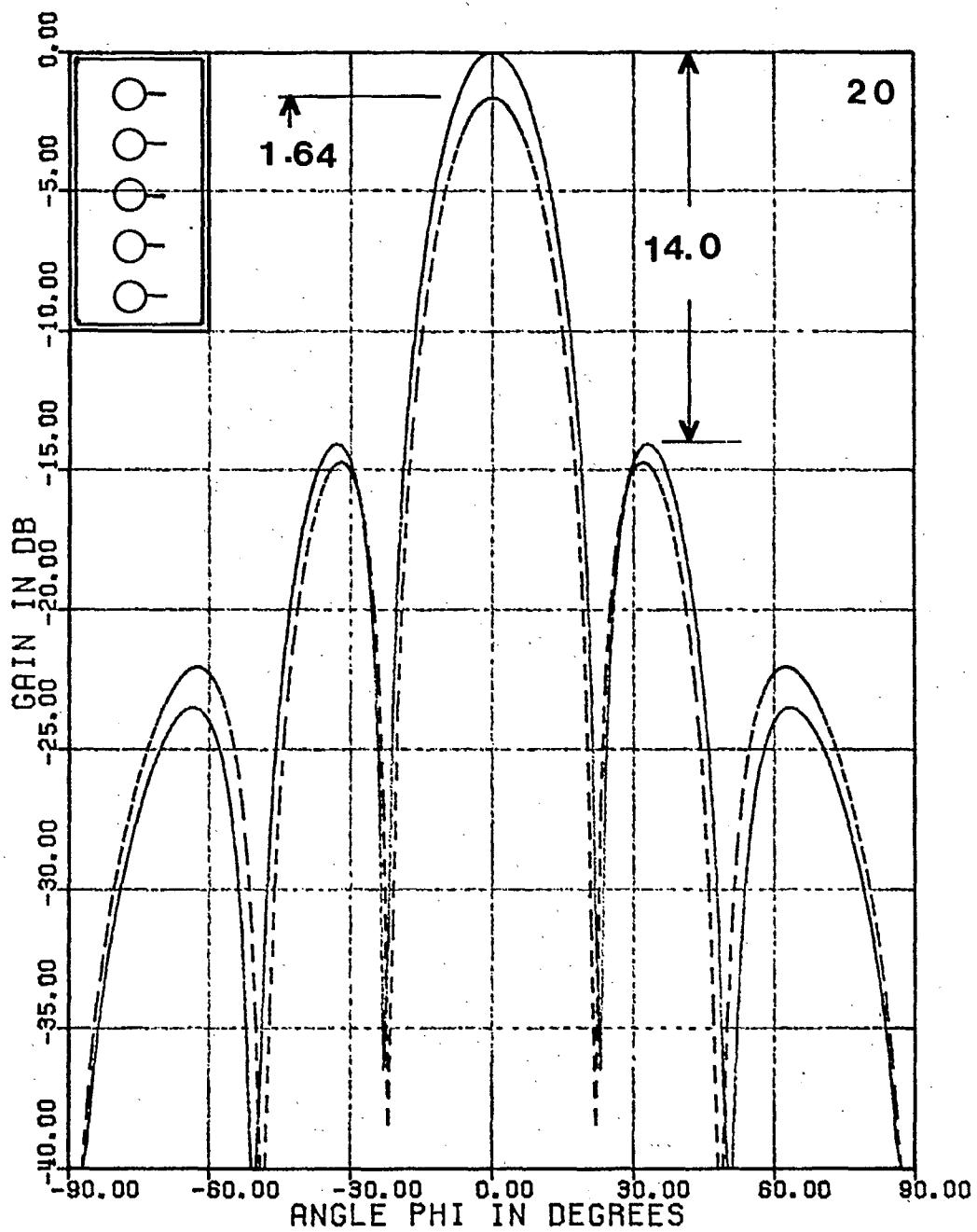


Fig. 18



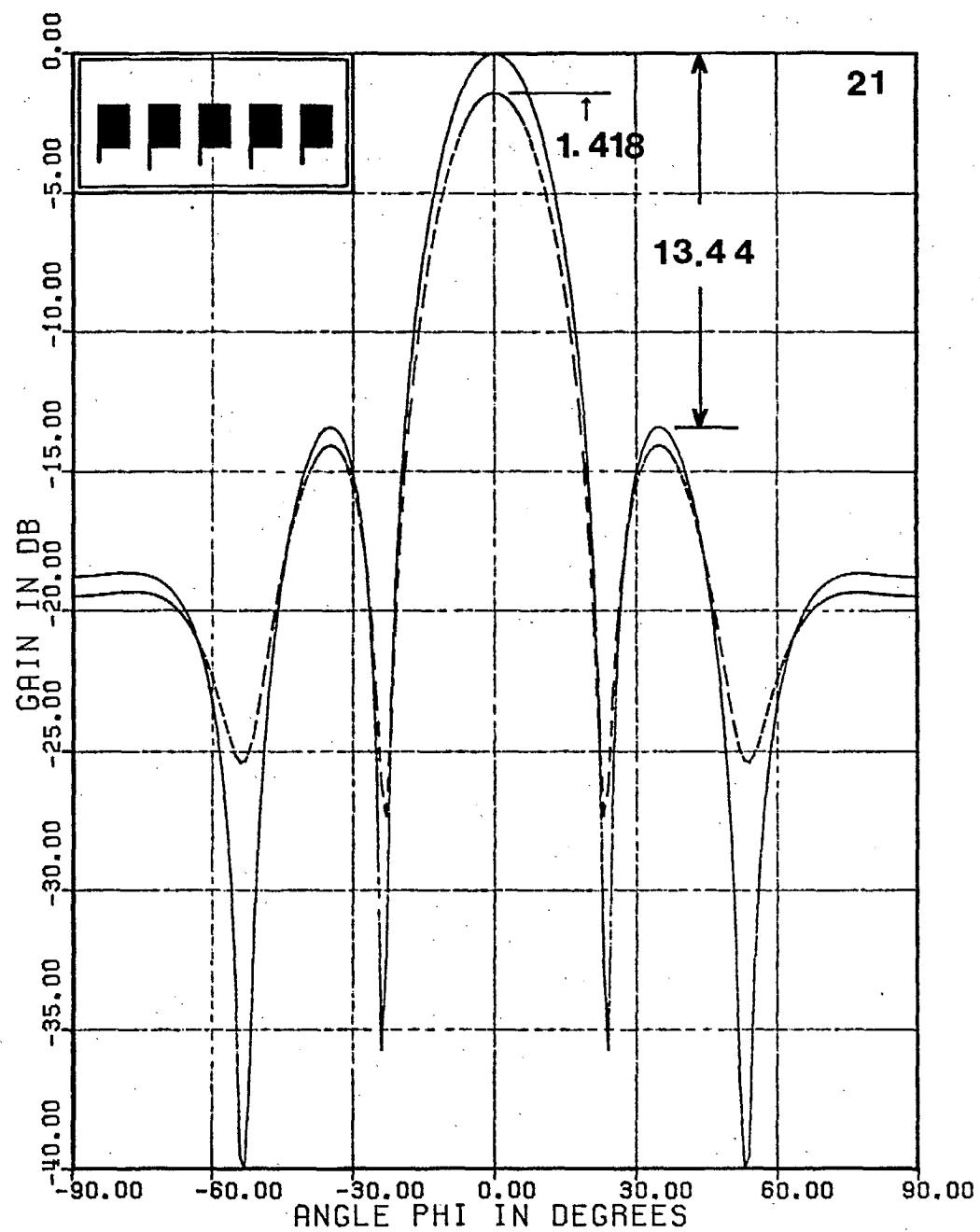
A 5 CIRCULAR PATCH... ARRAY
PATCHES ARE 0.52λ APART ON X-AXIS
E-PLANE RADIATION PATTERN

Fig. I9



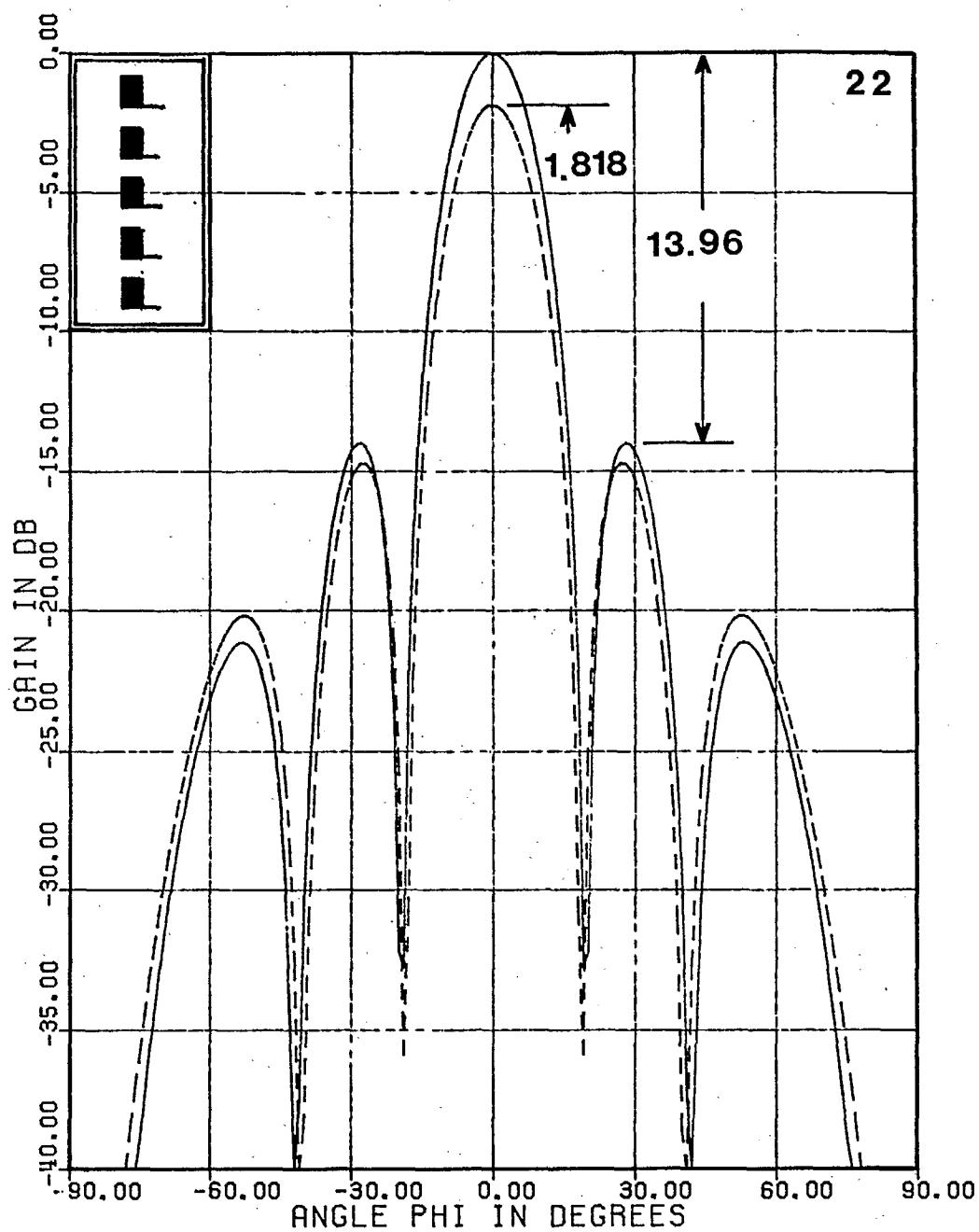
A 5 CIRCULAR PATCH ARRAY
PATCHES ARE 0.52λ APART ON Y-AXIS
H-PLANE RADIATION PATTERN

Fig. 20



A 5 RECTANGULAR PATCH ARRAY
PATCHES ARE 0.5λ APART ON X-AXIS
E-PLANE RADIATION PATTERN

Fig. 21



A 5 RECTANGULAR PATCH ARRAY
PATCHES ARE 0.6λ APART ON Y-AXIS
H-PLANE RADIATION PATTERN

Fig. 22

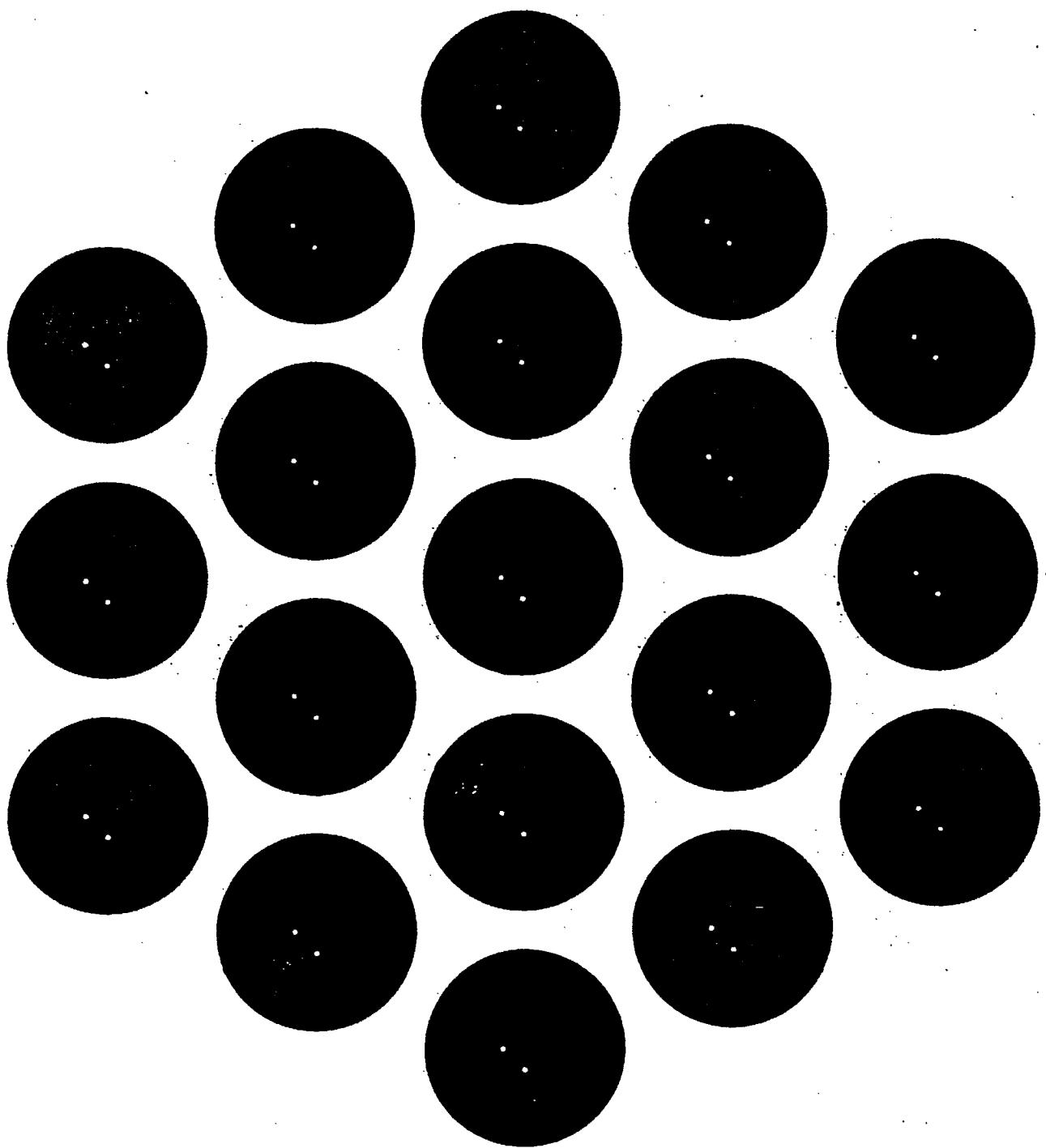


Fig. 23. Geometry of a 19-element circular patch array

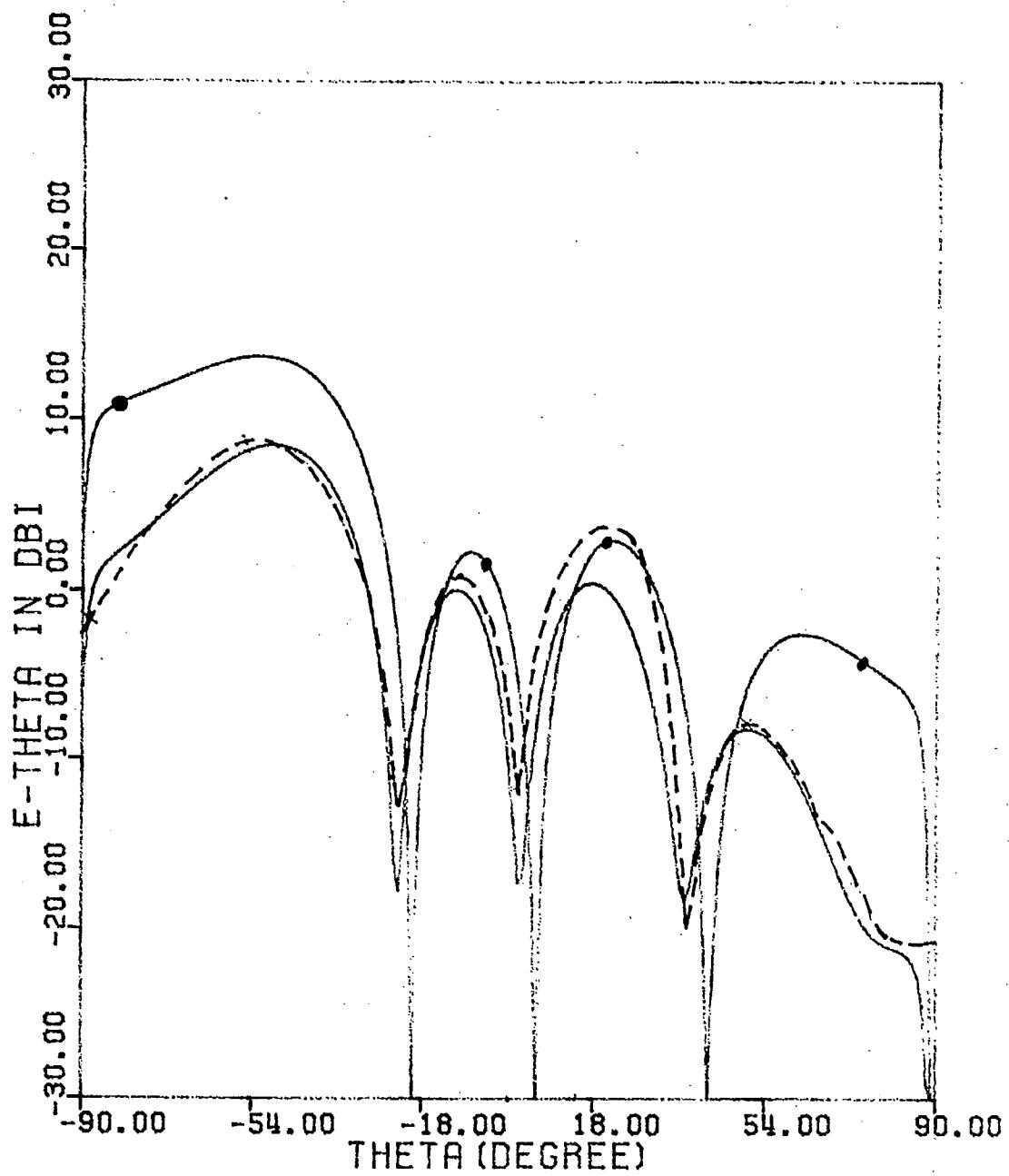


Fig. 24. Measured and computed patterns of the array in Fig. 23.

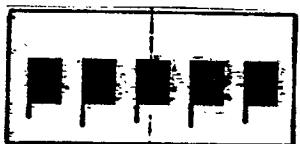
— measured, —●— computed no coupling

----- computed with coupling

APPENDIX A
RADIATION PATTERN OUTPUTS

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF RECTANGULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
-90.000	-19.486	-18.791	-40.000	-40.000
-89.000	-19.484	-18.789	-40.000	-38.887
-88.000	-19.479	-18.784	-34.282	-32.864
-87.000	-19.470	-18.775	-30.756	-29.338
-86.000	-19.459	-18.763	-28.251	-26.833
-85.000	-19.444	-18.749	-26.306	-24.888
-84.000	-19.428	-18.734	-24.714	-23.295
-83.000	-19.411	-18.717	-23.364	-21.946
-82.000	-19.393	-18.701	-22.193	-20.775
-81.000	-19.375	-18.685	-21.156	-19.738
-80.000	-19.360	-18.673	-20.226	-18.808
-79.000	-19.347	-18.664	-19.382	-17.964
-78.000	-19.338	-18.661	-18.609	-17.191
-77.000	-19.335	-18.665	-17.895	-16.477
-76.000	-19.339	-18.678	-17.231	-15.813
-75.000	-19.353	-18.702	-16.610	-15.192
-74.000	-19.377	-18.741	-16.027	-14.609
-73.000	-19.415	-18.796	-15.477	-14.059
-72.000	-19.468	-18.870	-14.956	-13.538
-71.000	-19.539	-18.967	-14.461	-13.042
-70.000	-19.630	-19.091	-13.988	-12.570
-69.000	-19.745	-19.245	-13.537	-12.119
-68.000	-19.887	-19.436	-13.105	-11.687
-67.000	-20.058	-19.667	-12.690	-11.271
-66.000	-20.263	-19.947	-12.290	-10.872
-65.000	-20.506	-20.284	-11.906	-10.487
-64.000	-20.790	-20.687	-11.534	-10.116
-63.000	-21.121	-21.170	-11.175	-9.757
-62.000	-21.500	-21.750	-10.828	-9.410
-61.000	-21.933	-22.449	-10.491	-9.073
-60.000	-22.419	-23.299	-10.165	-8.747
-59.000	-22.956	-24.344	-9.848	-8.430
-58.000	-23.532	-25.657	-9.541	-8.123
-57.000	-24.123	-27.356	-9.242	-7.824
-56.000	-24.682	-29.668	-8.951	-7.533
-55.000	-25.135	-33.119	-8.668	-7.250
-54.000	-25.382	-39.511	-8.392	-6.974
-53.000	-25.329	-40.000	-8.124	-6.705
-52.000	-24.936	-36.780	-7.862	-6.444
-51.000	-24.243	-31.074	-7.607	-6.189
-50.000	-23.347	-27.550	-7.358	-5.940



TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF RECTANGULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
-49.000	-22.348	-24.983	-7.116	-5.698
-48.000	-21.326	-22.965	-6.879	-5.461
-47.000	-20.330	-21.308	-6.649	-5.231
-46.000	-19.389	-19.912	-6.424	-5.006
-45.000	-18.516	-18.718	-6.205	-4.787
-44.000	-17.718	-17.687	-5.992	-4.573
-43.000	-16.997	-16.794	-5.783	-4.365
-42.000	-16.354	-16.022	-5.580	-4.162
-41.000	-15.789	-15.359	-5.383	-3.965
-40.000	-15.302	-14.798	-5.190	-3.772
-39.000	-14.893	-14.333	-5.003	-3.585
-38.000	-14.564	-13.964	-4.820	-3.402
-37.000	-14.317	-13.690	-4.643	-3.225
-36.000	-14.156	-13.512	-4.470	-3.052
-35.000	-14.086	-13.437	-4.303	-2.885
-34.000	-14.114	-13.471	-4.140	-2.722
-33.000	-14.251	-13.626	-3.982	-2.564
-32.000	-14.510	-13.918	-3.829	-2.411
-31.000	-14.910	-14.370	-3.681	-2.262
-30.000	-15.477	-15.017	-3.537	-2.119
-29.000	-16.250	-15.911	-3.398	-1.980
-28.000	-17.284	-17.139	-3.264	-1.846
-27.000	-18.667	-18.856	-3.135	-1.717
-26.000	-20.536	-21.381	-3.010	-1.592
-25.000	-23.079	-25.560	-2.890	-1.472
-24.000	-26.213	-35.699	-2.775	-1.357
-23.000	-27.332	-32.575	-2.664	-1.246
-22.000	-24.074	-23.499	-2.559	-1.140
-21.000	-20.358	-18.912	-2.457	-1.039
-20.000	-17.370	-15.770	-2.361	-0.943
-19.000	-14.978	-13.365	-2.269	-0.851
-18.000	-13.008	-11.415	-2.182	-0.764
-17.000	-11.346	-9.778	-2.100	-0.681
-16.000	-9.918	-8.374	-2.022	-0.604
-15.000	-8.675	-7.152	-1.949	-0.531
-14.000	-7.584	-6.080	-1.880	-0.462
-13.000	-6.622	-5.133	-1.817	-0.399
-12.000	-5.770	-4.294	-1.758	-0.340
-11.000	-5.016	-3.551	-1.704	-0.286
-10.000	-4.350	-2.895	-1.654	-0.236
-9.000	-3.764	-2.317	-1.609	-0.191

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF RECTANGULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
-8.000	-3.252	-1.812	-1.569	-0.151
-7.000	-2.810	-1.375	-1.534	-0.116
-6.000	-2.433	-1.003	-1.503	-0.085
-5.000	-2.118	-0.692	-1.477	-0.059
-4.000	-1.864	-0.441	-1.456	-0.038
-3.000	-1.668	-0.247	-1.439	-0.021
-2.000	-1.529	-0.109	-1.428	-0.009
-1.000	-1.446	-0.027	-1.420	-0.002
0.000	-1.418	0.000	-1.418	0.000
1.000	-1.446	-0.027	-1.420	-0.002
2.000	-1.529	-0.110	-1.428	-0.009
3.000	-1.668	-0.247	-1.439	-0.021
4.000	-1.864	-0.441	-1.456	-0.038
5.000	-2.119	-0.693	-1.477	-0.059
6.000	-2.433	-1.004	-1.503	-0.085
7.000	-2.810	-1.376	-1.534	-0.116
8.000	-3.253	-1.813	-1.569	-0.151
9.000	-3.765	-2.318	-1.609	-0.191
10.000	-4.351	-2.896	-1.654	-0.236
11.000	-5.017	-3.553	-1.704	-0.286
12.000	-5.771	-4.296	-1.758	-0.340
13.000	-6.622	-5.135	-1.817	-0.399
14.000	-7.585	-6.082	-1.880	-0.462
15.000	-8.676	-7.154	-1.949	-0.531
16.000	-9.919	-8.376	-2.022	-0.604
17.000	-11.347	-9.780	-2.100	-0.681
18.000	-13.009	-11.417	-2.182	-0.764
19.000	-14.978	-13.368	-2.269	-0.851
20.000	-17.370	-15.773	-2.361	-0.943
21.000	-20.357	-18.915	-2.457	-1.039
22.000	-24.072	-23.502	-2.559	-1.140
23.000	-27.331	-32.579	-2.664	-1.246
24.000	-26.215	-35.702	-2.775	-1.357
25.000	-23.083	-25.564	-2.890	-1.472
26.000	-20.539	-21.384	-3.010	-1.592
27.000	-18.670	-18.860	-3.135	-1.717
28.000	-17.287	-17.143	-3.264	-1.846
29.000	-16.253	-15.915	-3.398	-1.980
30.000	-15.480	-15.021	-3.537	-2.119
31.000	-14.913	-14.374	-3.681	-2.262
32.000	-14.513	-13.923	-3.829	-2.411

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF RECTANGULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
33.000	-14.254	-13.631	-3.982	-2.564
34.000	-14.117	-13.476	-4.140	-2.722
35.000	-14.089	-13.442	-4.303	-2.885
36.000	-14.158	-13.518	-4.470	-3.052
37.000	-14.319	-13.695	-4.643	-3.225
38.000	-14.566	-13.970	-4.820	-3.402
39.000	-14.895	-14.339	-5.003	-3.585
40.000	-15.304	-14.803	-5.190	-3.772
41.000	-15.791	-15.365	-5.383	-3.965
42.000	-16.356	-16.028	-5.580	-4.162
43.000	-16.998	-16.800	-5.783	-4.365
44.000	-17.718	-17.693	-5.992	-4.573
45.000	-18.516	-18.724	-6.205	-4.787
46.000	-19.388	-19.919	-6.424	-5.006
47.000	-20.328	-21.315	-6.649	-5.231
48.000	-21.323	-22.972	-6.879	-5.461
49.000	-22.344	-24.991	-7.116	-5.698
50.000	-23.342	-27.557	-7.358	-5.940
51.000	-24.238	-31.081	-7.607	-6.189
52.000	-24.932	-36.788	-7.862	-6.444
53.000	-25.326	-40.000	-8.124	-6.705
54.000	-25.382	-39.519	-8.392	-6.974
55.000	-25.138	-33.127	-8.668	-7.250
56.000	-24.687	-29.676	-8.951	-7.533
57.000	-24.130	-27.364	-9.242	-7.824
58.000	-23.541	-25.665	-9.541	-8.123
59.000	-22.966	-24.353	-9.848	-8.430
60.000	-22.430	-23.308	-10.165	-8.747
61.000	-21.944	-22.458	-10.491	-9.073
62.000	-21.512	-21.759	-10.828	-9.410
63.000	-21.132	-21.180	-11.175	-9.757
64.000	-20.802	-20.696	-11.534	-10.116
65.000	-20.518	-20.293	-11.906	-10.487
66.000	-20.276	-19.957	-12.290	-10.872
67.000	-20.071	-19.677	-12.690	-11.271
68.000	-19.899	-19.446	-13.105	-11.687
69.000	-19.758	-19.256	-13.537	-12.119
70.000	-19.643	-19.101	-13.988	-12.570
71.000	-19.552	-18.978	-14.461	-13.042
72.000	-19.481	-18.881	-14.956	-13.538
73.000	-19.428	-18.807	-15.477	-14.059

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF RECTANGULAR PATCHES

PHI ***	E-PLANE		H-PLANE	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
74.000	-19.390	-18.752	-16.027	-14.609
75.000	-19.366	-18.713	-16.610	-15.192
76.000	-19.352	-18.689	-17.231	-15.813
77.000	-19.348	-18.676	-17.895	-16.477
78.000	-19.351	-18.672	-18.609	-17.191
79.000	-19.360	-18.675	-19.382	-17.964
80.000	-19.372	-18.684	-20.226	-18.808
81.000	-19.388	-18.697	-21.156	-19.738
82.000	-19.406	-18.712	-22.193	-20.775
83.000	-19.424	-18.729	-23.364	-21.946
84.000	-19.441	-18.745	-24.714	-23.295
85.000	-19.457	-18.761	-26.306	-24.888
86.000	-19.472	-18.775	-28.251	-26.833
87.000	-19.483	-18.786	-30.756	-29.338
88.000	-19.492	-18.795	-34.282	-32.864
89.000	-19.497	-18.800	-40.000	-38.887
90.000	-19.499	-18.802	-40.000	-40.000

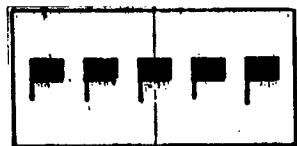
**** TABULATION COMPLETED ****

TABULATION OF THE RADIATION PATTERN
 A PLANAR ARRAY OF RECTANGULAR PATCHES

 PHI

 E-PLANE

 H-PLANE



IN DEGREES	FIELD IN DB		FIELD IN DB	
	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
-90.000	-6.628	-4.811	-40.000	-40.000
-89.000	-6.626	-4.809	-40.000	-40.000
-88.000	-6.621	-4.804	-40.000	-40.000
-87.000	-6.612	-4.795	-40.000	-40.000
-86.000	-6.599	-4.782	-40.000	-40.000
-85.000	-6.583	-4.766	-40.000	-40.000
-84.000	-6.563	-4.747	-40.000	-40.000
-83.000	-6.540	-4.724	-40.000	-40.000
-82.000	-6.514	-4.697	-40.000	-40.000
-81.000	-6.484	-4.667	-40.000	-40.000
-80.000	-6.450	-4.634	-40.000	-40.000
-79.000	-6.414	-4.597	-40.000	-40.000
-78.000	-6.374	-4.558	-39.959	-40.000
-77.000	-6.332	-4.515	-38.478	-40.000
-76.000	-6.286	-4.469	-37.065	-40.000
-75.000	-6.237	-4.421	-35.717	-38.913
-74.000	-6.186	-4.369	-34.431	-37.279
-73.000	-6.132	-4.315	-33.206	-35.756
-72.000	-6.075	-4.258	-32.039	-34.332
-71.000	-6.016	-4.199	-30.929	-33.001
-70.000	-5.954	-4.138	-29.873	-31.754
-69.000	-5.891	-4.074	-28.871	-30.586
-68.000	-5.825	-4.008	-27.922	-29.492
-67.000	-5.757	-3.940	-27.024	-28.469
-66.000	-5.687	-3.870	-26.178	-27.514
-65.000	-5.615	-3.798	-25.383	-26.625
-64.000	-5.541	-3.724	-24.638	-25.799
-63.000	-5.466	-3.649	-23.946	-25.037
-62.000	-5.390	-3.573	-23.305	-24.338
-61.000	-5.312	-3.495	-22.717	-23.701
-60.000	-5.233	-3.416	-22.183	-23.129
-59.000	-5.153	-3.336	-21.706	-22.622
-58.000	-5.072	-3.255	-21.286	-22.183
-57.000	-4.991	-3.174	-20.928	-21.814
-56.000	-4.908	-3.091	-20.634	-21.519
-55.000	-4.825	-3.008	-20.408	-21.304
-54.000	-4.741	-2.924	-20.257	-21.174
-53.000	-4.658	-2.840	-20.185	-21.139
-52.000	-4.573	-2.756	-20.203	-21.209
-51.000	-4.489	-2.672	-20.319	-21.399
-50.000	-4.405	-2.588	-20.548	-21.728

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF RECTANGULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
-49.000	-4.320	-2.503	-20.909	-22.223
-48.000	-4.236	-2.419	-21.425	-22.924
-47.000	-4.153	-2.335	-22.133	-23.892
-46.000	-4.069	-2.252	-23.087	-25.226
-45.000	-3.986	-2.169	-24.370	-27.108
-44.000	-3.904	-2.086	-26.131	-29.928
-43.000	-3.822	-2.005	-28.663	-34.819
-42.000	-3.741	-1.924	-32.671	-40.000
-41.000	-3.661	-1.843	-40.000	-37.466
-40.000	-3.581	-1.764	-38.521	-30.206
-39.000	-3.503	-1.686	-31.113	-26.155
-38.000	-3.426	-1.608	-26.839	-23.328
-37.000	-3.350	-1.532	-23.907	-21.173
-36.000	-3.275	-1.457	-21.703	-19.456
-35.000	-3.201	-1.384	-19.968	-18.059
-34.000	-3.129	-1.312	-18.572	-16.914
-33.000	-3.058	-1.241	-17.440	-15.981
-32.000	-2.989	-1.171	-16.529	-15.235
-31.000	-2.921	-1.103	-15.811	-14.662
-30.000	-2.855	-1.037	-15.272	-14.257
-29.000	-2.790	-0.973	-14.906	-14.022
-28.000	-2.727	-0.910	-14.714	-13.963
-27.000	-2.666	-0.849	-14.704	-14.099
-26.000	-2.607	-0.790	-14.896	-14.458
-25.000	-2.550	-0.732	-15.318	-15.087
-24.000	-2.494	-0.677	-16.023	-16.065
-23.000	-2.441	-0.623	-17.096	-17.534
-22.000	-2.389	-0.571	-18.695	-19.785
-21.000	-2.340	-0.522	-21.151	-23.565
-20.000	-2.292	-0.474	-25.358	-32.246
-19.000	-2.247	-0.429	-35.873	-32.756
-18.000	-2.204	-0.386	-30.798	-22.425
-17.000	-2.163	-0.345	-22.388	-17.525
-16.000	-2.124	-0.306	-17.919	-14.226
-15.000	-2.087	-0.269	-14.830	-11.725
-14.000	-2.053	-0.235	-12.466	-9.715
-13.000	-2.021	-0.203	-10.561	-8.046
-12.000	-1.991	-0.173	-8.980	-6.633
-11.000	-1.963	-0.145	-7.644	-5.422
-10.000	-1.938	-0.120	-6.505	-4.378
-9.000	-1.915	-0.098	-5.530	-3.477

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF RECTANGULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
-8.000	-1.895	-0.077	-4.696	-2.703
-7.000	-1.877	-0.059	-3.987	-2.041
-6.000	-1.861	-0.043	-3.391	-1.482
-5.000	-1.848	-0.030	-2.899	-1.020
-4.000	-1.837	-0.019	-2.504	-0.648
-3.000	-1.829	-0.011	-2.201	-0.362
-2.000	-1.823	-0.005	-1.988	-0.160
-1.000	-1.819	-0.001	-1.860	-0.040
0.000	-1.818	0.000	-1.818	0.000
1.000	-1.819	-0.001	-1.860	-0.040
2.000	-1.823	-0.005	-1.988	-0.160
3.000	-1.829	-0.011	-2.201	-0.362
4.000	-1.838	-0.020	-2.504	-0.648
5.000	-1.849	-0.031	-2.899	-1.020
6.000	-1.862	-0.044	-3.391	-1.482
7.000	-1.878	-0.060	-3.987	-2.041
8.000	-1.896	-0.078	-4.696	-2.703
9.000	-1.917	-0.099	-5.530	-3.477
10.000	-1.940	-0.122	-6.505	-4.378
11.000	-1.965	-0.147	-7.644	-5.422
12.000	-1.993	-0.175	-8.980	-6.633
13.000	-2.023	-0.205	-10.561	-8.046
14.000	-2.055	-0.237	-12.466	-9.715
15.000	-2.090	-0.271	-14.830	-11.725
16.000	-2.127	-0.308	-17.919	-14.226
17.000	-2.166	-0.347	-22.388	-17.525
18.000	-2.207	-0.388	-30.798	-22.425
19.000	-2.250	-0.432	-35.873	-32.756
20.000	-2.296	-0.477	-25.358	-32.246
21.000	-2.343	-0.525	-21.151	-23.565
22.000	-2.393	-0.575	-18.695	-19.785
23.000	-2.445	-0.626	-17.096	-17.534
24.000	-2.498	-0.680	-16.023	-16.065
25.000	-2.554	-0.736	-15.318	-15.087
26.000	-2.612	-0.793	-14.896	-14.458
27.000	-2.671	-0.853	-14.704	-14.099
28.000	-2.732	-0.914	-14.714	-13.963
29.000	-2.795	-0.977	-14.906	-14.022
30.000	-2.860	-1.042	-15.272	-14.257
31.000	-2.926	-1.108	-15.811	-14.662
32.000	-2.994	-1.176	-16.529	-15.235

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF RECTANGULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
33.000	-3.064	-1.245	-17.440	-15.981
34.000	-3.135	-1.316	-18.572	-16.914
35.000	-3.207	-1.389	-19.968	-18.059
36.000	-3.281	-1.463	-21.703	-19.456
37.000	-3.356	-1.538	-23.907	-21.173
38.000	-3.433	-1.614	-26.839	-23.328
39.000	-3.510	-1.691	-31.113	-26.155
40.000	-3.589	-1.770	-38.521	-30.206
41.000	-3.668	-1.849	-40.000	-37.466
42.000	-3.748	-1.930	-32.671	-40.000
43.000	-3.830	-2.011	-28.663	-34.819
44.000	-3.912	-2.093	-26.131	-29.928
45.000	-3.994	-2.175	-24.370	-27.108
46.000	-4.077	-2.259	-23.087	-25.226
47.000	-4.161	-2.342	-22.133	-23.892
48.000	-4.245	-2.426	-21.425	-22.924
49.000	-4.329	-2.510	-20.909	-22.223
50.000	-4.414	-2.595	-20.548	-21.728
51.000	-4.498	-2.679	-20.319	-21.399
52.000	-4.583	-2.764	-20.203	-21.209
53.000	-4.667	-2.848	-20.185	-21.139
54.000	-4.751	-2.932	-20.257	-21.174
55.000	-4.835	-3.016	-20.408	-21.304
56.000	-4.918	-3.099	-20.634	-21.519
57.000	-5.001	-3.182	-20.928	-21.814
58.000	-5.083	-3.264	-21.286	-22.183
59.000	-5.164	-3.345	-21.706	-22.622
60.000	-5.244	-3.425	-22.183	-23.129
61.000	-5.323	-3.504	-22.717	-23.701
62.000	-5.401	-3.582	-23.305	-24.338
63.000	-5.478	-3.659	-23.946	-25.037
64.000	-5.553	-3.734	-24.638	-25.799
65.000	-5.627	-3.808	-25.383	-26.625
66.000	-5.699	-3.879	-26.178	-27.514
67.000	-5.769	-3.950	-27.024	-28.469
68.000	-5.837	-4.018	-27.922	-29.492
69.000	-5.903	-4.084	-28.871	-30.586
70.000	-5.967	-4.148	-29.873	-31.754
71.000	-6.029	-4.210	-30.929	-33.001
72.000	-6.088	-4.269	-32.039	-34.332
73.000	-6.145	-4.326	-33.206	-35.756

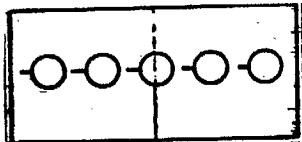
TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF RECTANGULAR PATCHES

PHI ***	E-PLANE		H-PLANE	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
74.000	-6.199	-4.380	-34.431	-37.279
75.000	-6.251	-4.432	-35.717	-38.913
76.000	-6.300	-4.480	-37.065	-40.000
77.000	-6.345	-4.526	-38.478	-40.000
78.000	-6.388	-4.569	-39.959	-40.000
79.000	-6.428	-4.609	-40.000	-40.000
80.000	-6.464	-4.645	-40.000	-40.000
81.000	-6.498	-4.678	-40.000	-40.000
82.000	-6.528	-4.708	-40.000	-40.000
83.000	-6.554	-4.735	-40.000	-40.000
84.000	-6.578	-4.758	-40.000	-40.000
85.000	-6.597	-4.778	-40.000	-40.000
86.000	-6.613	-4.794	-40.000	-40.000
87.000	-6.626	-4.807	-40.000	-40.000
88.000	-6.635	-4.816	-40.000	-40.000
89.000	-6.641	-4.821	-40.000	-40.000
90.000	-6.642	-4.823	-40.000	-40.000

**** TABULATION COMPLETED ****

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF CIRCULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
-90.000	-20.558	-19.805	-40.000	-40.000
-89.000	-20.552	-19.799	-37.213	-36.682
-88.000	-20.536	-19.783	-31.193	-30.661
-87.000	-20.509	-19.756	-27.671	-27.139
-86.000	-20.472	-19.719	-25.172	-24.640
-85.000	-20.425	-19.673	-23.233	-22.702
-84.000	-20.370	-19.618	-21.649	-21.118
-83.000	-20.307	-19.556	-20.310	-19.779
-82.000	-20.237	-19.487	-19.149	-18.618
-81.000	-20.162	-19.413	-18.126	-17.595
-80.000	-20.084	-19.335	-17.210	-16.679
-79.000	-20.003	-19.254	-16.382	-15.851
-78.000	-19.921	-19.173	-15.626	-15.095
-77.000	-19.841	-19.093	-14.930	-14.399
-76.000	-19.764	-19.017	-14.286	-13.755
-75.000	-19.692	-18.945	-13.686	-13.155
-74.000	-19.628	-18.882	-13.125	-12.594
-73.000	-19.574	-18.828	-12.599	-12.067
-72.000	-19.533	-18.786	-12.102	-11.571
-71.000	-19.507	-18.760	-11.633	-11.101
-70.000	-19.500	-18.752	-11.187	-10.656
-69.000	-19.515	-18.765	-10.764	-10.233
-68.000	-19.555	-18.803	-10.360	-9.829
-67.000	-19.625	-18.871	-9.975	-9.444
-66.000	-19.729	-18.972	-9.607	-9.076
-65.000	-19.873	-19.113	-9.254	-8.722
-64.000	-20.062	-19.298	-8.915	-8.384
-63.000	-20.303	-19.536	-8.589	-8.058
-62.000	-20.606	-19.834	-8.276	-7.745
-61.000	-20.981	-20.204	-7.974	-7.443
-60.000	-21.440	-20.660	-7.683	-7.152
-59.000	-22.000	-21.218	-7.402	-6.871
-58.000	-22.682	-21.902	-7.131	-6.599
-57.000	-23.515	-22.747	-6.868	-6.337
-56.000	-24.536	-23.799	-6.615	-6.083
-55.000	-25.796	-25.137	-6.369	-5.838
-54.000	-27.360	-26.890	-6.131	-5.600
-53.000	-29.278	-29.313	-5.901	-5.369
-52.000	-31.433	-33.021	-5.677	-5.146
-51.000	-32.944	-40.000	-5.461	-4.929
-50.000	-32.189	-40.000	-5.250	-4.719



TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF CIRCULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
-49.000	-29.734	-34.805	-5.047	-4.516
-48.000	-27.135	-29.609	-4.849	-4.318
-47.000	-24.870	-26.284	-4.657	-4.126
-46.000	-22.956	-23.831	-4.471	-3.940
-45.000	-21.332	-21.893	-4.290	-3.759
-44.000	-19.941	-20.302	-4.114	-3.583
-43.000	-18.742	-18.966	-3.944	-3.413
-42.000	-17.704	-17.831	-3.779	-3.248
-41.000	-16.805	-16.861	-3.619	-3.087
-40.000	-16.030	-16.032	-3.463	-2.932
-39.000	-15.369	-15.331	-3.312	-2.781
-38.000	-14.815	-14.746	-3.166	-2.635
-37.000	-14.364	-14.271	-3.024	-2.493
-36.000	-14.015	-13.905	-2.887	-2.356
-35.000	-13.769	-13.648	-2.754	-2.223
-34.000	-13.629	-13.501	-2.625	-2.094
-33.000	-13.603	-13.473	-2.501	-1.970
-32.000	-13.699	-13.573	-2.380	-1.849
-31.000	-13.933	-13.817	-2.264	-1.733
-30.000	-14.326	-14.229	-2.152	-1.621
-29.000	-14.910	-14.841	-2.044	-1.513
-28.000	-15.732	-15.709	-1.940	-1.409
-27.000	-16.871	-16.920	-1.839	-1.308
-26.000	-18.460	-18.632	-1.743	-1.212
-25.000	-20.759	-21.176	-1.650	-1.119
-24.000	-24.380	-25.441	-1.562	-1.031
-23.000	-31.085	-36.207	-1.477	-0.946
-22.000	-32.077	-31.560	-1.396	-0.865
-21.000	-23.946	-22.850	-1.318	-0.787
-20.000	-19.257	-18.341	-1.245	-0.713
-19.000	-16.035	-15.232	-1.175	-0.643
-18.000	-13.576	-12.846	-1.108	-0.577
-17.000	-11.591	-10.909	-1.046	-0.515
-16.000	-9.931	-9.284	-0.987	-0.456
-15.000	-8.514	-7.892	-0.931	-0.400
-14.000	-7.287	-6.684	-0.880	-0.348
-13.000	-6.214	-5.626	-0.831	-0.300
-12.000	-5.272	-4.696	-0.787	-0.256
-11.000	-4.442	-3.876	-0.746	-0.215
-10.000	-3.713	-3.154	-0.709	-0.177
-9.000	-3.074	-2.521	-0.675	-0.144

TABULATION OF THE RADIATION PATTERN
 A PLANAR ARRAY OF CIRCULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
-8.000	-2.517	-1.970	-0.645	-0.114
-7.000	-2.036	-1.494	-0.618	-0.087
-6.000	-1.628	-1.088	-0.595	-0.064
-5.000	-1.288	-0.751	-0.575	-0.044
-4.000	-1.013	-0.478	-0.560	-0.028
-3.000	-0.801	-0.268	-0.547	-0.016
-2.000	-0.651	-0.119	-0.538	-0.007
-1.000	-0.561	-0.030	-0.533	-0.002
0.000	-0.531	0.000	-0.531	0.000
1.000	-0.561	-0.030	-0.533	-0.002
2.000	-0.651	-0.119	-0.538	-0.007
3.000	-0.801	-0.268	-0.547	-0.016
4.000	-1.013	-0.478	-0.560	-0.028
5.000	-1.288	-0.751	-0.575	-0.044
6.000	-1.628	-1.088	-0.595	-0.064
7.000	-2.036	-1.494	-0.618	-0.087
8.000	-2.517	-1.970	-0.645	-0.114
9.000	-3.074	-2.521	-0.675	-0.144
10.000	-3.713	-3.154	-0.709	-0.177
11.000	-4.442	-3.876	-0.746	-0.215
12.000	-5.272	-4.696	-0.787	-0.256
13.000	-6.214	-5.626	-0.831	-0.300
14.000	-7.287	-6.684	-0.880	-0.348
15.000	-8.514	-7.892	-0.931	-0.400
16.000	-9.931	-9.284	-0.987	-0.456
17.000	-11.591	-10.909	-1.046	-0.515
18.000	-13.576	-12.846	-1.108	-0.577
19.000	-16.035	-15.232	-1.175	-0.643
20.000	-19.257	-18.341	-1.245	-0.713
21.000	-23.946	-22.850	-1.318	-0.787
22.000	-32.077	-31.560	-1.396	-0.865
23.000	-31.085	-36.207	-1.477	-0.946
24.000	-24.380	-25.441	-1.562	-1.031
25.000	-20.759	-21.176	-1.650	-1.119
26.000	-18.460	-18.632	-1.743	-1.212
27.000	-16.871	-16.920	-1.839	-1.308
28.000	-15.732	-15.709	-1.940	-1.409
29.000	-14.910	-14.841	-2.044	-1.513
30.000	-14.326	-14.229	-2.152	-1.621
31.000	-13.933	-13.817	-2.264	-1.733
32.000	-13.699	-13.573	-2.380	-1.849

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF CIRCULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
33.000	-13.603	-13.473	-2.501	-1.970
34.000	-13.629	-13.501	-2.625	-2.094
35.000	-13.769	-13.648	-2.754	-2.223
36.000	-14.015	-13.905	-2.887	-2.356
37.000	-14.364	-14.271	-3.024	-2.493
38.000	-14.815	-14.746	-3.166	-2.635
39.000	-15.369	-15.331	-3.312	-2.781
40.000	-16.030	-16.032	-3.463	-2.932
41.000	-16.805	-16.861	-3.619	-3.087
42.000	-17.704	-17.831	-3.779	-3.248
43.000	-18.742	-18.966	-3.944	-3.413
44.000	-19.941	-20.302	-4.114	-3.583
45.000	-21.332	-21.893	-4.290	-3.759
46.000	-22.956	-23.831	-4.471	-3.940
47.000	-24.870	-26.284	-4.657	-4.126
48.000	-27.135	-29.609	-4.849	-4.318
49.000	-29.734	-34.805	-5.047	-4.516
50.000	-32.189	-40.000	-5.250	-4.719
51.000	-32.944	-40.000	-5.461	-4.929
52.000	-31.433	-33.021	-5.677	-5.146
53.000	-29.278	-29.313	-5.901	-5.369
54.000	-27.360	-26.890	-6.131	-5.600
55.000	-25.796	-25.137	-6.369	-5.838
56.000	-24.536	-23.799	-6.615	-6.083
57.000	-23.515	-22.747	-6.868	-6.337
58.000	-22.682	-21.902	-7.131	-6.599
59.000	-22.000	-21.218	-7.402	-6.871
60.000	-21.440	-20.660	-7.683	-7.152
61.000	-20.981	-20.204	-7.974	-7.443
62.000	-20.606	-19.834	-8.276	-7.745
63.000	-20.303	-19.536	-8.589	-8.058
64.000	-20.062	-19.298	-8.915	-8.384
65.000	-19.873	-19.113	-9.254	-8.722
66.000	-19.729	-18.972	-9.607	-9.076
67.000	-19.625	-18.871	-9.975	-9.444
68.000	-19.555	-18.803	-10.360	-9.829
69.000	-19.515	-18.765	-10.764	-10.233
70.000	-19.500	-18.752	-11.187	-10.656
71.000	-19.507	-18.760	-11.633	-11.101
72.000	-19.533	-18.786	-12.102	-11.571
73.000	-19.574	-18.828	-12.599	-12.067

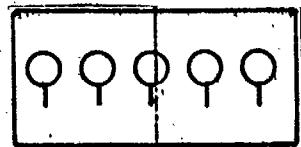
TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF CIRCULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
74.000	-19.628	-18.882	-13.125	-12.594
75.000	-19.692	-18.945	-13.686	-13.155
76.000	-19.764	-19.017	-14.286	-13.755
77.000	-19.841	-19.093	-14.930	-14.399
78.000	-19.921	-19.173	-15.626	-15.095
79.000	-20.003	-19.254	-16.382	-15.851
80.000	-20.084	-19.335	-17.210	-16.679
81.000	-20.162	-19.413	-18.126	-17.595
82.000	-20.237	-19.487	-19.149	-18.618
83.000	-20.307	-19.556	-20.310	-19.779
84.000	-20.370	-19.618	-21.649	-21.118
85.000	-20.425	-19.673	-23.233	-22.702
86.000	-20.472	-19.719	-25.172	-24.640
87.000	-20.509	-19.756	-27.671	-27.139
88.000	-20.536	-19.783	-31.193	-30.661
89.000	-20.552	-19.799	-37.213	-36.682
90.000	-20.558	-19.805	-40.000	-40.000

**** TABULATION COMPLETED ****

TABULATION OF THE RADIATION PATTERN
 A PLANAR ARRAY OF CIRCULAR PATCHES

PHI ***	E-PLANE		H-PLANE	
	*****		*****	
IN	FIELD IN DB		FIELD IN DB	
--	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
DEGREES	-----	-----	-----	-----
-90.000	-7.048	-5.406	-40.000	-40.000
-89.000	-7.046	-5.404	-40.000	-40.000
-88.000	-7.040	-5.398	-40.000	-40.000
-87.000	-7.029	-5.388	-40.000	-40.000
-86.000	-7.015	-5.373	-37.740	-38.986
-85.000	-6.997	-5.355	-35.779	-37.020
-84.000	-6.974	-5.333	-34.168	-35.404
-83.000	-6.948	-5.306	-32.798	-34.028
-82.000	-6.917	-5.276	-31.606	-32.829
-81.000	-6.883	-5.242	-30.549	-31.766
-80.000	-6.845	-5.204	-29.600	-30.810
-79.000	-6.804	-5.162	-28.739	-29.943
-78.000	-6.758	-5.117	-27.953	-29.151
-77.000	-6.710	-5.068	-27.231	-28.424
-76.000	-6.658	-5.016	-26.567	-27.755
-75.000	-6.602	-4.961	-25.954	-27.140
-74.000	-6.544	-4.902	-25.388	-26.574
-73.000	-6.482	-4.841	-24.868	-26.054
-72.000	-6.417	-4.776	-24.391	-25.581
-71.000	-6.350	-4.709	-23.955	-25.152
-70.000	-6.280	-4.639	-23.562	-24.769
-69.000	-6.208	-4.566	-23.210	-24.432
-68.000	-6.133	-4.491	-22.901	-24.142
-67.000	-6.055	-4.414	-22.636	-23.901
-66.000	-5.976	-4.334	-22.417	-23.714
-65.000	-5.895	-4.253	-22.247	-23.582
-64.000	-5.811	-4.170	-22.129	-23.512
-63.000	-5.726	-4.085	-22.066	-23.509
-62.000	-5.640	-3.999	-22.065	-23.580
-61.000	-5.552	-3.911	-22.131	-23.737
-60.000	-5.463	-3.821	-22.273	-23.990
-59.000	-5.372	-3.731	-22.500	-24.358
-58.000	-5.281	-3.640	-22.825	-24.862
-57.000	-5.189	-3.547	-23.267	-25.537
-56.000	-5.096	-3.454	-23.849	-26.428
-55.000	-5.002	-3.361	-24.604	-27.614
-54.000	-4.908	-3.267	-25.582	-29.223
-53.000	-4.814	-3.172	-26.862	-31.510
-52.000	-4.719	-3.077	-28.575	-35.089
-51.000	-4.624	-2.983	-30.975	-40.000
-50.000	-4.529	-2.888	-34.622	-40.000



TABULATION OF THE RADIATION PATTERN
 A PLANAR ARRAY OF CIRCULAR PATCHES

PHI ***	E-PLANE		H-PLANE	
	*****		*****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
-49.000	-4.435	-2.793	-40.000	-36.527
-48.000	-4.340	-2.699	-40.000	-31.227
-47.000	-4.246	-2.605	-34.319	-27.805
-46.000	-4.153	-2.511	-29.889	-25.259
-45.000	-4.060	-2.418	-26.814	-23.234
-44.000	-3.967	-2.326	-24.474	-21.559
-43.000	-3.876	-2.234	-22.600	-20.145
-42.000	-3.785	-2.144	-21.052	-18.935
-41.000	-3.695	-2.054	-19.750	-17.894
-40.000	-3.607	-1.965	-18.645	-16.999
-39.000	-3.519	-1.878	-17.706	-16.234
-38.000	-3.433	-1.791	-16.911	-15.589
-37.000	-3.348	-1.706	-16.248	-15.058
-36.000	-3.264	-1.623	-15.708	-14.638
-35.000	-3.182	-1.540	-15.285	-14.330
-34.000	-3.101	-1.460	-14.980	-14.136
-33.000	-3.022	-1.381	-14.795	-14.062
-32.000	-2.945	-1.303	-14.734	-14.119
-31.000	-2.869	-1.228	-14.809	-14.323
-30.000	-2.796	-1.154	-15.034	-14.695
-29.000	-2.724	-1.082	-15.433	-15.272
-28.000	-2.654	-1.012	-16.042	-16.105
-27.000	-2.586	-0.944	-16.915	-17.284
-26.000	-2.520	-0.878	-18.143	-18.966
-25.000	-2.456	-0.814	-19.894	-21.481
-24.000	-2.394	-0.753	-22.516	-25.719
-23.000	-2.334	-0.693	-26.972	-36.460
-22.000	-2.277	-0.636	-38.386	-31.789
-21.000	-2.222	-0.581	-31.292	-23.057
-20.000	-2.169	-0.528	-23.423	-18.526
-19.000	-2.119	-0.477	-19.109	-15.398
-18.000	-2.071	-0.429	-16.092	-12.994
-17.000	-2.025	-0.384	-13.764	-11.040
-16.000	-1.982	-0.340	-11.871	-9.399
-15.000	-1.941	-0.300	-10.283	-7.992
-14.000	-1.903	-0.262	-8.925	-6.771
-13.000	-1.867	-0.226	-7.749	-5.701
-12.000	-1.834	-0.193	-6.724	-4.759
-11.000	-1.804	-0.162	-5.826	-3.929
-10.000	-1.776	-0.134	-5.040	-3.198
-9.000	-1.750	-0.109	-4.354	-2.556

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF CIRCULAR PATCHES

PHI ***	E-PLANE *****		H-PLANE *****	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
-8.000	-1.727	-0.086	-3.757	-1.997
-7.000	-1.707	-0.066	-3.244	-1.515
-6.000	-1.690	-0.048	-2.808	-1.104
-5.000	-1.675	-0.034	-2.446	-0.761
-4.000	-1.663	-0.022	-2.153	-0.485
-3.000	-1.654	-0.012	-1.928	-0.272
-2.000	-1.647	-0.005	-1.768	-0.120
-1.000	-1.643	-0.001	-1.673	-0.030
0.000	-1.641	0.000	-1.641	0.000
1.000	-1.643	-0.001	-1.673	-0.030
2.000	-1.647	-0.005	-1.768	-0.120
3.000	-1.654	-0.012	-1.928	-0.272
4.000	-1.663	-0.022	-2.153	-0.485
5.000	-1.675	-0.034	-2.446	-0.761
6.000	-1.690	-0.048	-2.808	-1.104
7.000	-1.707	-0.066	-3.244	-1.515
8.000	-1.727	-0.086	-3.757	-1.997
9.000	-1.750	-0.109	-4.354	-2.556
10.000	-1.776	-0.134	-5.040	-3.198
11.000	-1.804	-0.162	-5.826	-3.929
12.000	-1.834	-0.193	-6.724	-4.759
13.000	-1.867	-0.226	-7.749	-5.701
14.000	-1.903	-0.262	-8.925	-6.771
15.000	-1.941	-0.300	-10.283	-7.992
16.000	-1.982	-0.340	-11.871	-9.399
17.000	-2.025	-0.384	-13.764	-11.040
18.000	-2.071	-0.429	-16.092	-12.994
19.000	-2.119	-0.477	-19.109	-15.398
20.000	-2.169	-0.528	-23.423	-18.526
21.000	-2.222	-0.581	-31.292	-23.057
22.000	-2.277	-0.636	-38.386	-31.789
23.000	-2.334	-0.693	-26.972	-36.460
24.000	-2.394	-0.753	-22.516	-25.719
25.000	-2.456	-0.814	-19.894	-21.481
26.000	-2.520	-0.878	-18.143	-18.966
27.000	-2.586	-0.944	-16.915	-17.284
28.000	-2.654	-1.012	-16.042	-16.105
29.000	-2.724	-1.082	-15.433	-15.272
30.000	-2.796	-1.154	-15.034	-14.695
31.000	-2.869	-1.228	-14.809	-14.323
32.000	-2.945	-1.303	-14.734	-14.119

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF CIRCULAR PATCHES

PHI ***	E-PLANE		H-PLANE	
IN --	FIELD IN DB		FIELD IN DB	
DEGREES	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
33.000	-3.022	-1.381	-14.795	-14.062
34.000	-3.101	-1.460	-14.980	-14.136
35.000	-3.182	-1.540	-15.285	-14.330
36.000	-3.264	-1.623	-15.708	-14.638
37.000	-3.348	-1.706	-16.248	-15.058
38.000	-3.433	-1.791	-16.911	-15.589
39.000	-3.519	-1.878	-17.706	-16.234
40.000	-3.607	-1.965	-18.645	-16.999
41.000	-3.695	-2.054	-19.750	-17.894
42.000	-3.785	-2.144	-21.052	-18.935
43.000	-3.876	-2.234	-22.600	-20.145
44.000	-3.967	-2.326	-24.474	-21.559
45.000	-4.060	-2.418	-26.814	-23.234
46.000	-4.153	-2.511	-29.889	-25.259
47.000	-4.246	-2.605	-34.319	-27.805
48.000	-4.340	-2.699	-40.000	-31.227
49.000	-4.435	-2.793	-40.000	-36.527
50.000	-4.529	-2.888	-34.622	-40.000
51.000	-4.624	-2.983	-30.975	-40.000
52.000	-4.719	-3.077	-28.575	-35.089
53.000	-4.814	-3.172	-26.862	-31.510
54.000	-4.908	-3.267	-25.582	-29.223
55.000	-5.002	-3.361	-24.604	-27.614
56.000	-5.096	-3.454	-23.849	-26.428
57.000	-5.189	-3.547	-23.267	-25.537
58.000	-5.281	-3.640	-22.825	-24.862
59.000	-5.372	-3.731	-22.500	-24.358
60.000	-5.463	-3.821	-22.273	-23.990
61.000	-5.552	-3.911	-22.131	-23.737
62.000	-5.640	-3.999	-22.065	-23.580
63.000	-5.726	-4.085	-22.066	-23.509
64.000	-5.811	-4.170	-22.129	-23.512
65.000	-5.895	-4.253	-22.247	-23.582
66.000	-5.976	-4.334	-22.417	-23.714
67.000	-6.055	-4.414	-22.636	-23.901
68.000	-6.133	-4.491	-22.901	-24.142
69.000	-6.208	-4.566	-23.210	-24.432
70.000	-6.280	-4.639	-23.562	-24.769
71.000	-6.350	-4.709	-23.955	-25.152
72.000	-6.417	-4.776	-24.391	-25.581
73.000	-6.482	-4.841	-24.868	-26.054

TABULATION OF THE RADIATION PATTERN
A PLANAR ARRAY OF CIRCULAR PATCHES

PHI	E-PLANE		H-PLANE	
***	*****		*****	
IN	FIELD IN DB		FIELD IN DB	
--	COUPLED	UNCOUPLED	COUPLED	UNCOUPLED
DEGREES	-----	-----	-----	-----
74.000	-6.544	-4.902	-25.388	-26.574
75.000	-6.602	-4.961	-25.954	-27.140
76.000	-6.658	-5.016	-26.567	-27.755
77.000	-6.710	-5.068	-27.231	-28.424
78.000	-6.758	-5.117	-27.953	-29.151
79.000	-6.804	-5.162	-28.739	-29.943
80.000	-6.845	-5.204	-29.600	-30.810
81.000	-6.883	-5.242	-30.549	-31.766
82.000	-6.917	-5.276	-31.606	-32.829
83.000	-6.948	-5.306	-32.798	-34.028
84.000	-6.974	-5.333	-34.168	-35.404
85.000	-6.997	-5.355	-35.779	-37.020
86.000	-7.015	-5.373	-37.740	-38.986
87.000	-7.029	-5.388	-40.000	-40.000
88.000	-7.040	-5.398	-40.000	-40.000
89.000	-7.046	-5.404	-40.000	-40.000
90.000	-7.048	-5.406	-40.000	-40.000

**** TABULATION COMPLETED ****

APPENDIX B
COMPUTER PROGRAM LISTING

```

100. //RZIN JOB ',,T=10,I=5','LAI'
101. // EXEC WATFIV,P=D
102. /*ROUTE PRINT SELF
103. //FT08F001 DD DSN=LLAI.RZIN,DISP=OLD
104. //SYSIN DD *
105. $JOB WATFIV LAI,NOEXT
106. C ****
107. C * INPUT IMPEDANCE OF AN ISOLATED RECTANGULAR(SQUARE) PATCH. *
108. C ****
109. IMPLICIT REAL*8(A-B,D-H,O-Z),COMPLEX*16(C),INTEGER(I-N)
110. REAL*8 L,LN,LPOS,LNEG
111. REAL Z0,Z1,Z2
112. PI=3.14159265D0
113. L=0.3036D0*1.001
114. LPOS=0.22909D0*1.001
115. LNEG=0.07451D0*1.001
116. W=0.50292D0
117. H=0.01166D0
118. CJ=(0.D0,1.D0)
119. FACT=376.73D0*H/W
120. EPSR=2.62
121. EPSE=(EPSR-1)/DSQRT(1.+12.*H/W)
122. EPSE=EPSR+1.+EPSE
123. EPSE=0.5*EPSE
124. FN=0.96D0
125. WHILE (FN.LE.1.04D0) DO
126. HN=H*FN
127. WN=W*FN
128. LN=L*FN
129. ARG=DSQRT(EPSE)*2.*PI*FN
130. SNPT=DSIN(ARG*LPOS)
131. SNMT=DSIN(ARG*LNEG)
132. BOTT=DSIN(ARG*L)
133. SNP=SNPT/BOTT
134. SNM=SNMT/BOTT
135. CALL PSLFR(LN,WN,HN,EPSR,GII,BII)
136. CP=DCMPLX(GII,BII)
137. CALL QSLFR(LN,WN,HN,EPSR,QSR,QSI)
138. CQ=DCMPLX(QSR,QSI)
139. CZ1=(SNP**2+SNM**2)*CP
140. CZ2=2.*SNP*SNM*CQ
141. CZ=(CZ1-CZ2)/(CP**2-CQ**2)+SNP*SNM*BOTT*CJ/DSQRT(EPSE)
142. CZ=FACT*CZ
143. Z0=FN
144. Z1=DREAL(CZ)
145. Z2=DAIMAG(CZ)
146. WRITE(8,*)Z0,Z1,Z2
147. FN=FN+0.001D0
148. END WHILE
149. STOP
150. END
151. C ****
152. C * SUBROUTINE OF SELF APERTURE NORMALIZED ADMITTANCE. *
153. C ****
154. SUBROUTINE PSLFR(LN,WN,HN,EPSR,GII,BII)
155. IMPLICIT REAL*8(A-B,D-H,O-Z),COMPLEX*16(C),INTEGER(I-N)
156. REAL*8 LN,LPOS,LNEG
157. TOPI=2.*3.14159265D0
158. WK0=TOPI*WN
159. X2=WK0**2
160. GII=0.D0
161. K=1
162. DO 1 I=1,26
163. K=-K

```

```
164.      J=26-I
165.      M=2*I+1
166.      MM=(M+1)*(M+2)
167. 1     GII=(GII+K/M)*X2/MM
168.      GII=(4/TOPI)*(HN/WN)*GII
169.      EPSE=(EPSR-1)/DSQRT(1.+12.*HN/WN)
170.      EPSE=EPSE+EPSR+1.
171.      EPSE=0.5*EPSE
172.      EP=DSQRT(EPSE)
173.      F1=(EPSE+.3)*(.264+WN/HN)
174.      F2=(EPSE-.258)*(.8+WN/HN)
175.      F3=.412*F1/F2
176.      BII=TOPI*HN*F3*EPSE
177.      BII=BII-EP*DCOTAN(TOPI*LN*EP)
178.      RETURN
179.      END
180. C ****
181. C * SUBROUTINE OF THE MUTUAL ADMITTANCE BETWEEN TWO APERTURES *
182. C * OF THE SAME PATCH. *
183. C ****
184.      SUBROUTINE QSLFR(LN,WN,HN,EPSR,QSR,QSI)
185.      IMPLICIT REAL*8(A-B,D-H,O-Z),COMPLEX*16(C),INTEGER(I-N)
186.      REAL*8 LN,LPOS,LNEG,LN2
187.      N=30
188.      TOPI=2.*3.14159265D0
189.      R11=LN
190.      LN2=LN*LN
191.      R12=DSQRT(LN2+WN**2)
192.      X11=TOPI*R11
193.      X12=TOPI*R12
194.      QSR=2.* (DSIN(X12)/X12-DSIN(X11)/X11)
195.      QSI=2.* (DCOS(X12)/X12-DCOS(X11)/X11)
196.      SUM1=0.D0
197.      SUM2=0.D0
198.      DO 1 I=1,N
199.      DO 1 J=1,N
200.      RIJ=DSQRT(LN2+((J-I)*WN/N)**2)
201.      XIJ=TOPI*RIJ
202.      SUM1=SUM1+DSIN(XIJ)/XIJ
203. 1     SUM2=SUM2+DCOS(XIJ)/XIJ
204.      DEL2=(TOPI*WN/N)**2
205.      SUM1=SUM1*DEL2
206.      SUM2=SUM2*DEL2
207.      QSR=-(HN/WN)*(QSR+SUM1)/TOPI
208.      QSI=-(HN/WN)*(QSI+SUM2)/TOPI
209.      EPSE=(EPSR-1.)/DSQRT(1.+12.*HN/WN)
210.      EPSE=EPSE+EPSR+1.
211.      EPSE=0.5*EPSE
212.      EP=DSQRT(EPSE)
213.      QSI=QSI+EP/DSIN(TOPI*LN*EP)
214.      RETURN
215.      END
216. $ENTRY
```

User Guide to the Numerical Programming

Program function: Radiation pattern of a microstrip planar array.
Patch shape: Rectangular(square) elements of identical dimensions.

- (1) User must fill in the array dimensions(line 117) with the exact number of radiating patches because of the library subroutine LEQT1C.
- (2) Output data are written into a dataset(allocated by the user, line 103 and 334) from the radiation pattern program. The data are then read by a plot-program from the dataset for plotting.
- (3) Due to different array shape, three cases are introduced to the user for time saving purpose.
Case A: A planar array with arbitrary patch locations.
Case B: A rectangular array shape with a fixed horizontal and vertical separations between adjacent patches.
Case C: An array of Case B above with missing patch(patches).
- (4) Starting from the bottom row, the patches are numbered from left to right.
- (5) Input data in the dimension of length are normalized to the wavelength of the corresponding operating frequency.
- (6) Input data are in double precision and are supplied in free format as follows:

Case A: N,0,0,0
 LN,WN,HN,EPSR
 XCR(I),YCR(I),CI(I),SNP(I),SNM(I)

 XCR(N),YCR(N),CI(N),SNP(N),SNM(N)
 NPDG

Case B: 0,NX,NY,0
 LN,WN,HN,EPSR
 HDN,VDN
 CI(1),SNP(1),SNM(1)
 . . .
 . . .
 CI(K),SNP(K),SNM(K)
 NPDG

Case C: 0,NX,NY,M
 LN,WN,HN,EPSR
 HDN,VDN
 CI(1),SNP(1),SNM(1)
 . . .
 . . .
 CI(K),SNP(K),SNM(K)
 MCEL(1),...,MCEL(M)
 NPDG

Variable Dictionary

N	- integer, total number of radiating elements not in rectangular array form.
NX,NY	- integer, number of radiating elements on the x-, and y-axis which are in a rectangular array form.
LN,WN	- real, normalized length and width of the patch.
HN	- real, normalized depth of the substrate.
EPSR	- real, relative permittivity of the substrate.
XCR(I),YCR(I)	- real, normalized x and y coordinates of the Ith radiating element.
CI(I)	- complex, normalized feeding current of the Ith cell.
SNP(I),SNM(I)	- real, the mathematical expression $\sin(KL^+)/\sin(KL)$ and $\sin(KL^-)/\sin(KL)$ respectively.
K	- integer, product of NX and NY.
HDN,VDN	- real, normalized horizontal and vertical distance between any two adjacent patches.
MCEL(I)	- integer, position number of the Ith missing patch in ascending order.
NPDG	- integer, number of points per degree for the radiation pattern plot.

Remark: Subroutine LEQT1C of the program is the IMSL (International Mathematical and Statistical Library) routine for complex matrix equation solution.

```

100. //RARRAY JOB ',,T=2,I=5,F=BDK1','LAI'
101. // EXEC WATFIV,SIZE=1024K,P=D
102. /*ROUTE PRINT LOCAL
103. //FT08F001 DD DSN=LLAI.DATA,DISP=OLD
104. //SYSIN DD *
105. $JOB WATFIV LAI,NOEXT
106. C ****
107. C * A PROGRAM OF FINDING THE RADIATION PATTERN OF A SQUARE OR *
108. C * RECTANGULAR MICROSTRIP PLANAR ARRAY. CASES A,B,AND C ARE *
109. C * INTRODUCED DUE TO DIFFERENT ARRAY SHAPE IN ORDER TO *
110. C * REDUCE THE COMPUTING TIME. *
111. C ****
112. IMPLICIT REAL*8(A-B,D-H,O-Z),COMPLEX*16(C),INTEGER(I-N)
113. C ****
114. C * USER MUST USE A DECLARED DUMMY ARRAYS WITH EXACT NUMBER *
115. C * OF RADIATING PATCHES BECAUSE OF THE SUBROUTINE LEQT1C. *
116. C ****
117. COMPLEX*16 CRPL( , ),CRM1( , ),CEP( ),CEM( )
118. COMPLEX*16 CP(100,100),CQ(100,100),CQT(100,100),CI(100)
119. COMPLEX*16 CE(100),DCMPLX
120. REAL*8 XCR(100),YCR(100),SNP(100),SNM(100),LN,WA(100)
121. REAL*8 ANG(500),EC(500),E(500),HC(500),H(500),CDABS
122. REAL ALOG10,Z0,Z1,Z2,Z3,Z4
123. INTEGER MCEL(20)
124. READ,NR,NX,NY,NMIS
125. READ,LN,WN,HN,EPSR
126. IF (NR.EQ.0) GOTO 5
127. C ****
128. C * CASE A *
129. C * A GENERAL PLANAR ARRAY WITH ARBITRARY PATCH LOCATIONS. *
130. C ****
131. DO 1 I=1,NR
132. 1 READ,XCR(I),YCR(I),CI(I),SNP(I),SNM(I)
133. CALL PSLFR(LN,WN,HN,EPSR,GII,BII)
134. CALL QSLFR(LN,WN,HN,EPSR,QSR,QSI)
135. DO 2 I=1,NR
136. CE(I)=CI(I)
137. CP(I,I)=DCMPLX(GII,BII)
138. 2 CQ(I,I)=DCMPLX(QSR,QSI)
139. L=NR-1
140. IF (L.GT.0) THEN
141. DO 3 I=1,L
142. XI=XCR(I)
143. YI=YCR(I)
144. K=I+1
145. DO 3 J=K,NR
146. XJ=XCR(J)
147. YJ=YCR(J)
148. CALL PMUTR(LN,WN,HN,XI,YI,XJ,YJ,PMR,PMI)
149. CP(I,J)=DCMPLX(PMR,PMI)
150. CALL QMUTR(LN,WN,HN,XI,YI,XJ,YJ,QMR,QMI)
151. 3 CQ(I,J)=DCMPLX(QMR,QMI)
152. DO 4 I=2,NR
153. XI=XCR(I)
154. YI=YCR(I)
155. K=I-1
156. DO 4 J=1,K
157. XJ=XCR(J)
158. YJ=YCR(J)
159. CP(I,J)=CP(J,I)
160. CALL QMUTR(LN,WN,HN,XI,YI,XJ,YJ,QMR,QMI)
161. 4 CQ(I,J)=DCMPLX(QMR,QMI)
162. END IF
163. NEND=NR

```

164. GOTO 19
165. C ****
166. C * CASE B
167. C * PLANAR ARRAY IN A RECTANGULAR SHAPE WITH A FIXED VERTICAL *
168. C * AND HORIZONTAL SEPARATIONS BETWEEN ADJACENT PATCHES. *
169. C ****
170. 5 NCAL=NX*NY
171. READ,HDN,VDN
172. DO 6 J=1,NY
173. DO 6 I=1,NX
174. XCR(NX*(J-1)+I)=(I-1)*HDN
175. 6 YCR(NX*(J-1)+I)=(J-1)*VDN
176. CALL PSLFR(LN,WN,HN,EPSR,GII,BII)
177. CALL QSLFR(LN,WN,HN,EPSR,QSR,QSI)
178. DO 7 I=1,NCAL
179. CP(I,I)=DCMPLX(GII,BII)
180. 7 CQ(I,I)=DCMPLX(QSR,QSI)
181. YI=YCR(1)
182. DO 8 J=2,NCAL
183. XI=XCR(1)
184. XJ=XCR(J)
185. YJ=YCR(J)
186. CALL PMUTR(LN,WN,HN,XI,YI,XJ,YJ,PMR,PMI)
187. CP(1,J)=DCMPLX(PMR,PMI)
188. CALL QMUTR(LN,WN,HN,XI,YI,XJ,YJ,QMR,QMI)
189. CQ(1,J)=DCMPLX(QMR,QMI)
190. XI=XCR(J)
191. XJ=XCR(1)
192. CALL QMUTR(LN,WN,HN,XI,YI,XJ,YJ,QMR,QMI)
193. 8 CQ(J,1)=DCMPLX(QMR,QMI)
194. IF (NX.EQ.1) GOTO 11
195. DO 10 I=2,NX
196. DO 10 K=1,NY
197. MM=(K-1)*NX+I
198. NN=K*NX
199. IM1=I-1
200. DO 9 L=1,IM1
201. 9 CQ(I,MM-L)=CQ(MM+L+1-I,1)
202. CP(I,MM-L)=CP(I-L,MM)
203. DO 10 J=MM,NN
204. CQ(I,J)=CQ(I-1,J-1)
205. 10 CP(I,J)=CP(I-1,J-1)
206. 11 CONTINUE
207. IF (NY.EQ.1) GOTO 14
208. DO 12 K=2,NY
209. I=(K-1)*NX
210. DO 12 L=1,NX
211. DO 12 J=1,NX
212. CQ(I+L,J)=CQ(L,J+I)
213. 12 CP(I+L,J)=CP(J,I+L)
214. NXP1=NX+1
215. DO 13 I=NXP1,NCAL
216. DO 13 J=NXP1,NCAL
217. CQ(I,J)=CQ(I-NX,J-NX)
218. 13 CP(I,J)=CP(I-NX,J-NX)
219. 14 DO 15 J=1,NCAL
220. 15 READ,CI(J),SNP(J),SNM(J)
221. NEND=NCAL
222. IF (NMIS.EQ.0) GOTO 19
223. C ****
224. C * CASE C
225. C * A PLANAR ARRAY OF CASE B WITH MISSING PATCHES. *
226. C ****
227. READ,(MCEL(I),I=1,NMIS)

228. K=1
229. WHILE (K.LE.NMIS) DO
230. NREV=NMIS+1-K
231. NSTRT=MCEL(NREV)
232. NEND=NCAL-K
233. IF (NSTRT.GT.NEND) GOTO 18
234. NNI=NEND+1
235. DO 16 I=1,NNI
236. DO 16 J=NSTRT,NEND
237. CQ(I,J)=CQ(I,J+1)
238. 16 CP(I,J)=CP(I,J+1)
239. DO 17 I=NSTRT,NEND
240. CI(I)=CI(I+1)
241. XCR(I)=XCR(I+1)
242. YCR(I)=YCR(I+1)
243. DO 17 J=1,NEND
244. CQ(I,J)=CQ(I+1,J)
245. 17 CP(I,J)=CP(I+1,J)
246. 18 K=K+1
247. END WHILE
248. 19 N=NEND
249. DO 20 I=1,N
250. DO 20 J=1,N
251. CQT(I,J)=CQ(J,I)
252. CRPL(I,J)=CP(I,J)
253. 20 CRM1(I,J)=(0.D0,0.D0)
254. DO 21 I=1,N
255. 21 CRM1(I,I)=(1.0D0,0.0D0)
256. CALL LEQT1C(CRPL,N,N,CRM1,N,N,0,WA,IER)
257. CCP=1./(CP(1,1)+CQ(1,1))
258. CCN=1./(CP(1,1)-CQ(1,1))
259. DO 24 I=1,N
260. CXX1=CYY1=(0.D0,0.D0)
261. DO 23 J=1,N
262. CXIJ=CXX=CYY=(0.D0,0.D0)
263. DO 22 K=1,N
264. CXX=CXX+CQT(I,K)*CRM1(K,J)
265. CYY=CYY+CQ(I,K)*CRM1(K,J)
266. DO 22 L=1,N
267. 22 CXIJ=CXIJ+CQT(I,K)*CRM1(K,L)*CQ(L,J)
268. CRPL(I,J)=CP(I,J)-CXIJ
269. CP(I,J)=CRPL(I,J)
270. CXX1=CXX1+SNP(J)*CI(J)*CXX
271. 23 CYY1=CYY1+SNM(J)*CI(J)*CYY
272. CEP(I)=SNM(I)*CI(I)-CXX1
273. 24 CEM(I)=SNP(I)*CI(I)-CYY1
274. CALL LEQT1C(CRPL,N,N,CEP,1,N,0,WA,IER)
275. DO 25 I=1,N
276. DO 25 J=1,N
277. 25 CRPL(I,J)=CP(I,J)
278. CALL LEQT1C(CRPL,N,N,CEM,1,N,0,WA,IER)
279. PI=3.14159265D0
280. READ,NPDG
281. DEL=1./NPDG
282. CJ=(0.0D0,1.0D0)
283. THETA=-90.D0
284. K=1
285. EMAX=HMAX=0.D0
286. WHILE (THETA.LE.90.D0) DO
287. RAD=PI*THETA/180.D0
288. SKL=PI*LN*DSIN(RAD)
289. SKW=PI*WN*DSIN(RAD)
290. CKL=SKL*CJ
291. IF (DABS(SKW).LE.1.D-6) FACT=1.D0

292. IF (DABS(SKW).GT.1.D-6) FACT=DCOS(RAD)*DSIN(SKW)/SKW
293. CEIX=CIIIX=CEIY=CIIY=(0.D0,0.D0)
294. DO 26 I=1,N
295. SNA=SNP(I)+SNM(I)
296. SND=SNP(I)-SNM(I)
297. CPOS=CI(I)*CCP
298. CNEG=CI(I)*CCN
299. CKXI=2.*PI*XCR(I)*DSIN(RAD)*CJ
300. CKYI=2.*PI*YCR(I)*DSIN(RAD)*CJ
301. CEIX=CEIX+(CEP(I)*CDEXP(CKL)-CEM(I)*CDEXP(-CKL))
302. * *CDEXP(CKXI)
303. CIIIX=CIIIX+(-SND*CNEG*DCOS(SKL)+SNA*CPOS*DSIN(SKL)*CJ)
304. * *CDEXP(CKXI)
305. CEIY=CEIY+(CEP(I)-CEM(I))*CDEXP(CKYI)
306. 26 CIIY=CIIY+(-SND*CNEG*CDEXP(CKYI))
307. ABEC=CDABS(CEIX)
308. ABE=CDABS(CIIIX)
309. ABHC=CDABS(FACT*CEIY)
310. ABH=CDABS(FACT*CIIY)
311. IF (ABE.GT.EMAX) EMAX=ABE
312. IF (ABH.GT.HMAX) HMAX=ABH
313. ANG(K)=THETA
314. EC(K)=ABEC
315. E(K)=ABE
316. HC(K)=ABHC
317. H(K)=ABH
318. THETA=THETA+DEL
319. K=K+1
320. END WHILE
321. J=K-1
322. DO 27 K=1,J
323. Z0=ANG(K)
324. Z1=EC(K)/EMAX
325. Z2=E(K)/EMAX
326. Z3=HC(K)/HMAX
327. Z4=H(K)/HMAX
328. IF (Z1.LE.0.01) Z1=0.01
329. IF (Z2.LE.0.01) Z2=0.01
330. IF (Z3.LE.0.01) Z3=0.01
331. IF (Z4.LE.0.01) Z4=0.01
332. Z1=20.* ALOG10(Z1)
333. Z2=20.* ALOG10(Z2)
334. Z3=20.* ALOG10(Z3)
335. Z4=20.* ALOG10(Z4)
336. 27 WRITE(8,*)Z0,Z1,Z2,Z3,Z4
337. STOP
338. END
339. C ****
340. C * SUBROUTINE OF SELF APERTURE NORMALIZED ADMITTANCE. *
341. C ****
342. SUBROUTINE PSLFR(LN,WN,HN,EPSR,GII,BII)
343. IMPLICIT REAL*8(A-B,D-H,O-Z),COMPLEX*16(C),INTEGER(I-N)
344. REAL*8 LN
345. TOPI=2.*3.14159265D0
346. WK0=TOPI*WN
347. X2=WK0**2
348. GII=0.D0
349. K=1
350. DO 28 I=1,26
351. K=-K
352. J=26-I
353. M=2*J+1
354. MM=(M+1)*(M+2)
355. 28 GII=(GII+K/M)*X2/MM

```

356.      GII=(4/TOPI)*(HN/WN)*GII
357.      EPSE=(EPSR-1)/DSQRT(1.+12.*HN/WN)
358.      EPSE=EPSE+EPSR+1
359.      EPSE=0.5*EPSE
360.      EP=DSQRT(EPSE)
361.      F1=(EPSE+.3)*(.264+WN/HN)
362.      F2=(EPSE-.258)*(.8+WN/HN)
363.      F3=.412*F1/F2
364.      BII=TOPI*HN*F3*EPSE
365.      BII=BII-EP*DCOTAN(TOPI*LN*EP)
366.      RETURN
367.      END
368. C ****
369. C * SUBROUTINE OF MUTUAL INTERACTION ADMITTANCE BETWEEN TWO *
370. C * APERTURES OF THE SAME POLARITY OF ANY TWO PATCHES. *
371. C ****
372.      SUBROUTINE PMUTR(LN,WN,HN,XI,YI,XJ,YJ,PMR,PMI)
373.      IMPLICIT REAL*8(A-B,D-H,O-Z),COMPLEX*16(C),INTEGER(I-N)
374.      REAL*8 LN
375.      N=30
376.      DX=XJ-XI
377.      DY=YJ-YI
378.      DX2=DX*DX
379.      R11=DSQRT(DX2+DY**2)
380.      R12=DSQRT(DX2+(DY+WN)**2)
381.      R21=DSQRT(DX2+(DY-WN)**2)
382.      TOPI=2.*3.14159265D0
383.      X11=TOPI*R11
384.      X12=TOPI*R12
385.      X21=TOPI*R21
386.      PMR=DSIN(X12)/X12+DSIN(X21)/X21-2.*DSIN(X11)/X11
387.      PMI=DCOS(X12)/X12+DCOS(X21)/X21-2.*DCOS(X11)/X11
388.      DEL2=(TOPI*WN/N)**2
389.      SUM1=0.D0
390.      SUM2=0.D0
391.      DO 29 K=1,N
392.      DO 29 L=1,N
393.      RKL=DSQRT(DX2+(DY+(L-K)*WN/N)**2)
394.      RK0=TOPI*RKL
395.      SUM1=SUM1+DSIN(RK0)/RK0
396. 29     SUM2=SUM2+DCOS(RK0)/RK0
397.      SUM1=SUM1*DEL2
398.      SUM2=SUM2*DEL2
399.      PMR=(PMR+SUM1)*(HN/WN)/TOPI
400.      PMI=(PMI+SUM2)*(HN/WN)/TOPI
401.      RETURN
402.      END
403. C ****
404. C * SUBROUTINE OF THE MUTUAL ADMITTANCE BETWEEN TWO APERTURES *
405. C * OF THE SAME PATCH. *
406. C ****
407.      SUBROUTINE QSLFR(LN,WN,HN,EPSR,QSR,QSI)
408.      IMPLICIT REAL*8(A-B,D-H,O-Z),COMPLEX*16(C),INTEGER(I-N)
409.      REAL*8 LN,LN2
410.      N=30
411.      TOPI=2.*3.14159265D0
412.      R11=LN
413.      LN2=LN*LN
414.      R12=DSQRT(LN2+WN**2)
415.      X11=TOPI*R11
416.      X12=TOPI*R12
417.      QSR=2.* (DSIN(X12)/X12-DSIN(X11)/X11)
418.      QSI=2.* (DCOS(X12)/X12-DCOS(X11)/X11)
419.      SUM1=0.D0

```

```
420.      SUM2=0.D0
421.      DO 30 I=1,N
422.      DO 30 J=1,N
423.      RIJ=DSQRT(LN2+((J-I)*WN/N)**2)
424.      XIJ=TOPI*RIJ
425.      SUM1=SUM1+DSIN(XIJ)/XIJ
426. 30    SUM2=SUM2+DCOS(XIJ)/XIJ
427.      DEL2=(TOPI*WN/N)**2
428.      SUM1=SUM1*DEL2
429.      SUM2=SUM2*DEL2
430.      QSR=-(HN/WN)*(QSR+SUM1)/TOPI
431.      QSI=-(HN/WN)*(QSI+SUM2)/TOPI
432.      EPSE=(EPSR-1.)/DSQRT(1.+12.*HN/WN)
433.      EPSE=EPSE+EPSR+1.
434.      EPSE=0.5*EPSE
435.      EP=DSQRT(EPSE)
436.      QSI=QSI+EP/DSIN(TOPI*LN*EP)
437.      RETURN
438.      END
439. C ****
440. C * SUBROUTINE OF THE MUTUAL ADMITTANCE BETWEEN TWO APERTURES *
441. C * OF DIFFERENT POLARITY OF ANY TWO PATCHES. *
442. C ****
443.      SUBROUTINE QMUTR(LN,WN,HN,XI,YI,XJ,YJ,QMR,QMI)
444.      IMPLICIT REAL*8(A-B,D-H,O-Z),COMPLEX*16(C),INTEGER(I-N)
445.      REAL*8 LN
446.      N=30
447.      TOPI=2.*3.14159265D0
448.      DX2=(XJ-XI+LN)**2
449.      DY=YJ-YI
450.      R11=DSQRT(DX2+DY**2)
451.      R12=DSQRT(DX2+(DY+WN)**2)
452.      R21=DSQRT(DX2+(DY-WN)**2)
453.      X11=TOPI*R11
454.      X12=TOPI*R12
455.      X21=TOPI*R21
456.      QMR=DSIN(X12)/X12+DSIN(X21)/X21-2.*DSIN(X11)/X11
457.      QMI=DCOS(X12)/X12+DCOS(X21)/X21-2.*DCOS(X11)/X11
458.      SUM1=0.D0
459.      SUM2=0.D0
460.      DO 31 K=1,N
461.      DO 31 L=1,N
462.      RKL=DSQRT(DX2+(DY+(L-K)*WN/N)**2)
463.      RK0=TOPI*RKL
464.      SUM1=SUM1+DSIN(RK0)/RK0
465. 31    SUM2=SUM2+DCOS(RK0)/RK0
466.      DEL2=(TOPI*WN/N)**2
467.      SUM1=SUM1*DEL2
468.      SUM2=SUM2*DEL2
469.      QMR=-(HN/WN)*(QMR+SUM1)/TOPI
470.      QMI=-(HN/WN)*(QMI+SUM2)/TOPI
471.      RETURN
472.      END
473. $ENTRY
```

100. //PLOT JOB ',,T=3','LAI' - 93 -
101. /*ROUTE PRINT SELF
102. // EXEC FORTXCG,USERLIB='WAILEE.HPPLLOT.LOAD',SIZE=256K
103. //FORT.SYSIN DD *
104. C ****
105. C * A PLOT-PROGRAM OF THE RADIATION PATTERN.
106. C ****
107. DIMENSION IBUF(5000),X(183),EC(183),E(183),HC(183),H(183)
108. CALL PLOTS(IBUF,5000)
109. CALL FACTOR(0.8)
110. CALL PLOT(1.4,3.4,-3)
111. X(182)=-90.0
112. X(183)=30.0
113. EC(182)=-40.0
114. EC(183)=5.0
115. E(182)=-40.0
116. E(183)=5.0
117. HC(182)=-40.0
118. HC(183)=5.0
119. H(182)=-40.0
120. H(183)=5.0
121. CALL AXIS(0.0,0.0,'ANGLE PHI IN DEGREES',-20,6.0,0.0,X(182),
122. * X(183))
123. CALL AXIS(0.0,0.0,'GAIN IN DB',10,8.0,90.0,EC(182),EC(183))
124. CALL LLGRID(0.0,0.0,6.0,8.0,'LIN','LIN',6,8)
125. READ(8,*)(X(I),EC(I),E(I),HC(I),H(I),I=1,181)
126. CALL DASHL(X,EC,181,1,5)
127. CALL LINE(X,E,181,1,0,0)
128. CALL PLOT(12.0,0.0,999)
129. CALL PLOT(12.0,0.0,9999)
130. STOP
131. END
132. //GO.FT08F001 DD DSN=LLAI.DATA,DISP=SHR
133. //GO.FT01F001 DD DSN=LLAI.ZPLOT,DISP=OLD
134. //GO.SYSIN DD DUMMY
135. */

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```
100. //CZIN JOB ',,T=1,L=50,R=128K,I=5','LAI'
101. // EXEC WATFIV,SIZE=256K
102. /*ROUTE PRINT SELF
103. //FT09F001 DD DSN=LLAI.CZIN,DISP=OLD
104. //SYSIN DD *
105. $JOB WATFIV LAI,NOEXT
106. C ****
107. C * INPUT IMPEDANCE OF AN ISOLATED CIRCULAR PATCH. *
108. C ****
109. IMPLICIT REAL*8(A-B,D-H,O-Z),COMPLEX*16(C),INTEGER(I-N)
110. REAL Z0,Z1,Z2
111. NPHI=1
112. B=0.17419D0
113. H=0.01166D0
114. EPSR=2.62D0
115. FN=0.96D0
116. WHILE (FN.LE.1.04D0) DO
117. BN=B*FN
118. HN=H*FN
119. CALL PSLFC (NPHI,BN,HN,EPSR,GIIN,BIIN)
120. CP=DCMPLX(GIIN,BIIN)
121. CZ=120.* (H/B)/CP
122. X1=DREAL(CZ)
123. X2=DAIMAG(CZ)
124. Z0=FN
125. Z1=X1
126. Z2=X2
127. WRITE(9,*)Z0,Z1,Z2
128. FN=FN+0.001D0
129. END WHILE
130. STOP
131. END
132. C ****
133. C * A SUBROUTINE OF FINDING THE SELF WALL ADMITTANCE OF A *
134. C * SINGLE CIRCULAR PATCH. *
135. C ****
136. SUBROUTINE PSLFC(NPHI,BN,HN,EPSR,GIIN,BIIN)
137. IMPLICIT REAL*8(A-B,D-H,O-Z),INTEGER(I-N)
138. REAL*8 B1(50),B2(50)
139. PI=3.14159265D0
140. HB=HN/BN
141. BKO=2.0D0*PI*BN
142. BK=BKO*DSQRT(EPSR)
143. BEK=BK*DSQRT(1.0D0+2.*HB/(PI*EPSR)*(DLOG(PI/2.0D0/HB)+1.7726))
144. JACK=NPHI+1
145. GIIN=0.0D0
146. X2=BKO**2
147. K=-1
148. M=30
149. IF (NPHI.EQ.0) THEN
150. WHILE(M.GE.2)DO
151. K=-K
152. B=K*M*(M-1.0D0)/(2.0D0*M-1.0D0)
153. GIIN=(GIIN+B)*X2/M**2
154. M=M-1
155. END WHILE
156. GIIN=HB*GIIN*X2
157. CALL BGY(BK,1,B1,B2)
158. B1L1=B1(1)
159. B2L1=B2(1)
160. CALL BGY(BEK,2,B1,B2)
161. B1L2=B1(2)
162. B2L2=B2(2)
163. BUP=-2.0D0*B1L2/(PI*BKO*B1L1**2)
```

```
164.      BDI=B2L2-B2L1*B1L2/B1L1
165.      BIIN=BUP/BDI
166.      ELSE
167.      J=1
168.      NF=2*NPHI+1
169.      DO 25 I=1,NF
170. 25    J=J*I
171.      WHILE (M.GE.2) DO
172.      K=-K
173.      TMPN=2.0D0*(M+NPHI)
174.      B=NPHI+M*(M-1.0)/(TMPN-1.0D0)-NPHI**2/(TMPN+1.0D0)
175.      GIIN=(GIIN+K*B)*X2/(M*(M+2.0D0*NPHI))
176.      M=M-1
177.      END WHILE
178.      GIIN=GIIN*X2/J
179.      GIIN=HB*(GIIN-X2*NPHI*(NPHI+3.0D0)/((2.0D0*NPHI+3.0D0)*J)+*
180.      *      NPHI*(NPHI+1.0D0)/J)*X2**NPHI
181.      CALL BJV(BK,JACK,B1,B2)
182.      B1L0=B1(JACK)
183.      B2L0=B2(JACK)
184.      L1=JACK-1
185.      L2=JACK+1
186.      CALL BJV(BEK,L1,B1,B2)
187.      B1L1=B1(L1)
188.      B2L1=B2(L1)
189.      B1L2=B1(L2)
190.      B2L2=B2(L2)
191.      BUP=-2.0D0*(B1L1-B1L2)/(PI*BKO*B1L0**2)
192.      BDI=B2L1-B2L2-B2L0*(B1L1-B1L2)/B1L0
193.      BIIN=BUP/BDI
194.      END IF
195.      RETURN
196.      END
197. C ****
198. C * BESEL FUNCTIONS OF THE FIRST AND SECOND KIND ARE *
199. C * COMPUTED IN THIS SUBROUTINE BJV. *
200. C ****
201.      SUBROUTINE BJV(X,N,JN,YN)
202.      REAL*8 JN(50),YN(50),X,ANORM,PI,SUM,DLOG
203.      INTEGER N,IM,JX,JX1,I,L,AK,K
204.      IM=N+20
205.      JN(IM)=0.0D0
206.      JN(IM-1)=1.0D0
207.      JX=IM-2
208.      JX1=JX+1
209.      DO 27 I=1,JX
210.      L=JX1-I
211.      AK=L-1
212. 27    JN(L)=2.0D0*(AK+1.0)/X*JN(L+1)-JN(L+2)
213.      ANORM=JN(1)
214.      DO 28 L=3,IM,2
215. 28    ANORM=ANORM+2.0D0*JN(L)
216.      DO 29 L=1,IM
217. 29    JN(L)=JN(L)/ANORM
218.      PI=3.14159265D0
219.      SUM=0.0D0
220.      DO 30 L=3,IM,2
221.      K=(L-1)/2
222.      AK=K
223. 30    SUM=SUM+(-1.0)**K*JN(L)/AK
224.      YN(1)=2.0D0/PI*(DLOG(X/2.0D0)+0.577215664)*JN(1)-4./PI*SUM
225.      YN(2)=(JN(2)*YN(1)-2.0D0/(PI*X))/JN(1)
226.      DO 31 L=3,IM
227.      AK=L-1
```

228. 31 YN(L)=2.0D0*(AK-1.0D0)/X*YN(L-1)-YN(L-2)
229. RETURN
230. END
231. \$ENTRY

User Guide to the Numerical Programming

Program function: Radiation pattern of a microstrip planar array.
Patch shape: Circular elements of identical dimensions.

- (1) User must fill in the array dimensions(line 117) with the exact number of radiating patches because of the library subroutine LEQT1C.
- (2) Output data are written into a dataset(allocated by the user, line 103 and 301)from the radiation pattern program. The data are then read by a plot-program from the dataset for plotting.
- (3) Due to different array shape, three cases are introduced to the user for time saving purpose.
Case A: A planar array with arbitrary patch locations.
Case B: A rectangular array shape with a fixed horizontal and vertical separations between adjacent patches.
Case C: An array of Case B above with missing patch(patches).
- (4) Starting from the bottom row, the patches are numbered from left to right.
- (5) Input data in the dimension of length are normalized to the wavelength of the corresponding operating frequency.
- (6) Input data are in double precision and are supplied in free format as follows:

Case A: N,0,0,0,NPHI
 BN,HN,EPSR
 XCR(I),YCR(I),CI(I)
 . . .
 . . .
 XCR(N),YCR(N),CI(N)
 NPDG

Case B: 0,NX,NY,0,NPHI
 BN,HN,EPSR
 HDN,VDN
 CI(1)
 .
 CI(K)
 NPDG

Case C: 0,NX,NY,M,NPHI
 BN,HN,EPSR
 HDN,VDN
 CI(1)
 .
 CI(K)
 MCEL(1),...,MCEL(M)
 NPDG

Variable Dictionary

N	- integer, total number of radiating elements not in rectangular array form.
NX,NY	- integer, number of radiating elements on the x-, and y-axis which are in a rectangular array form.
BN	- real, normalized radius of the patch.
HN	- real, normalized depth of the substrate.
EPSR	- real, relative permittivity of the substrate.
XCR(I),YCR(I)	- real, normalized x and y coordinates of the Ith radiating element.
CI(I)	- complex, normalized feeding current of the Ith cell.
NPHI	- integer, number of phi-variation in the mode.
K	- integer, product of NX and NY.
HDN,VDN	- real, normalized horizontal and vertical distance between any two adjacent patches.
MCEL(I)	- integer, position number of the Ith missing patch in ascending order.
NPDG	- integer, number of points per degree for the radiation pattern plot.

Remark: Subroutine LEQT1C and MMBSJN of the program are the IMSL (International Mathematical and Statistical Library) routines for complex matrix equation solution and Bessel function respectively.

100. //CARRAY JOB ',,T=3M,L=50,R=128K,I=5,','LAI' - 99 -
101. // EXEC WATFIV,SIZE=512K,P=D
102. /*ROUTE PRINT LOCAL
103. //FT08F001 DD DSN=LLAI.DATA,DISP=OLD
104. //SYSIN DD *
105. \$JOB WATFIV LAI,NOEXT
106. C ****
107. C * A PROGRAM OF FINDING THE RADIATION PATTERN OF A CIRCULAR *
108. C * MICROSTRIP PLANAR ARRAY. CASES A, B, AND C ARE INTRODUCED *
109. C * DUE TO DIFFERENT ARRAY SHAPE IN ORDER TO REDUCE THE *
110. C * COMPUTING TIME. *
111. C ****
112. IMPLICIT REAL*8(A-B,D-H,O-Z),COMPLEX*16(C),INTEGER(I-N)
113. C ****
114. C * USER MUST USE A DECLARED DUMMY ARRAYS WITH EXACT NUMBER *
115. C * OF RADIATING PATCHES BECAUSE OF THE SUBROUTINE LEQT1C. *
116. C ****
117. COMPLEX*16 CRP(),CRI(),CRE()
118. COMPLEX*16 CP(100,100),CI(100),CE(100),DCMPLX
119. REAL*8 XCR(100),YCR(100),WA(100),CDABS,B1(50),B2(50)
120. REAL*8 ANG(500),EC(500),E(500),HC(500),H(500)
121. REAL ALOG10,Z0,Z1,Z2,Z3,Z4
122. INTEGER MCEL(20)
123. READ,NR,NX,NY,NMIS,NPHI
124. READ,BN,HN,EPSR
125. IF (NR.EQ.0) GOTO 5
126. C ****
127. C * CASE A *
128. C * A GENERAL PLANAR ARRAY WITH ARBITRARY PATCH LOCATIONS. *
129. C ****
130. DO 1 I=1,NR
131. 1 READ,XCR(I),YCR(I),CI(I)
132. CALL PSLFC(NPHI,BN,HN,EPSR,GIIN,BIIN)
133. DO 2 I=1,NR
134. CE(I)=CI(I)
135. 2 CP(I,I)=DCMPLX(GIIN,BIIN)
136. L=NR-1
137. IF (L.GT.0) THEN
138. DO 3 I=1,L
139. XI=XCR(I)
140. YI=YCR(I)
141. K=I+1
142. DO 3 J=K,NR
143. XJ=XCR(J)
144. YJ=YCR(J)
145. CALL PMUTC(NPHI,BN,HN,EPSR,XI,YI,XJ,YJ,GIJN,BIJN)
146. 3 CP(I,J)=DCMPLX(GIJN,BIJN)
147. DO 4 I=2,NR
148. K=I-1
149. DO 4 J=1,K
150. 4 CP(I,J)=CP(J,I)
151. END IF
152. IJOB=0
153. IA=NR
154. IB=NR
155. M=1
156. N=NR
157. GOTO 21
158. C ****
159. C * CASE B *
160. C * PLANAR ARRAY IN A RECTANGULAR SHAPE WITH A FIXED VERTICAL *
161. C * AND HORIZONTAL SEPARATIONS BETWEEN ADJACENT PATCHES. *
162. C ****
163. 5 NCAL=NX*NY

```
164.      READ,HDN,VDN
165.      DO 6 J=1,NY
166.      DO 6 I=1,NX
167.      XCR(NX*(J-1)+I)=(I-1)*HDN
168. 6     YCR(NX*(J-1)+I)=(J-1)*VDN
169.      CALL PSLFC(NPHI,BN,HN,EPSR,GIIN,BIIN)
170.      DO 7 I=1,NCAL
171. 7     CP(I,I)=DCMPLX(GIIN,BIIN)
172.      XI=XCR(1)
173.      YI=YCR(1)
174.      DO 8 J=2,NCAL
175.      XJ=XCR(J)
176.      YJ=YCR(J)
177.      CALL PMUTC(NPHI,BN,HN,EPSR,XI,YI,XJ,YJ,GIJN,BIJN)
178. 8     CP(1,J)=DCMPLX(GIJN,BIJN)
179.      IF (NX.EQ.1) GOTO 11
180.      DO 10 I=2,NX
181.      DO 10 K=1,NY
182.      MM=(K-1)*NX+I
183.      NN=K*NX
184.      IM1=I-1
185.      DO 9 L=1,IM1
186. 9     CP(I,MM-L)=CP(I-L,MM)
187.      DO 10 J=MM,NN
188. 10    CP(I,J)=CP(I-1,J-1)
189. 11    CONTINUE
190.      IF (NY.EQ.1) GOTO 14
191.      NXP1=NX+1
192.      DO 13 I=NXP1,NCAL
193.      IM1=I-1
194.      DO 12 L=1,IM1
195. 12    CP(I,L)=CP(L,I)
196.      DO 13 J=I,NCAL
197. 13    CP(I,J)=CP(I-NX,J-NX)
198. 14    DO 15 J=1,NCAL
199. 15    READ,CI(J)
200.      NEND=NCAL
201.      IF (NMIS.EQ.0) GOTO 19
202. C ****
203. C * CASE C *
204. C * A PLANAR ARRAY OF CASE B WITH MISSING PATCH(PATCHES). *
205. C ****
206.      READ,(MCEL(I),I=1,NMIS)
207.      K=1
208.      WHILE (K.LE.NMIS) DO
209.      NREV=NMIS+1-K
210.      NSTRT=MCEL(NREV)
211.      NEND=NCAL-K
212.      IF (NSTRT.GT.NEND) GOTO 18
213.      NNI=NEND+1
214.      DO 16 I=1,NNI
215.      DO 16 J=NSTART,NEND
216. 16    CP(I,J)=CP(I,J+1)
217.      DO 17 I=NSTART,NEND
218.      CI(I)=CI(I+1)
219.      XCR(I)=XCR(I+1)
220.      YCR(I)=YCR(I+1)
221.      DO 17 J=1,NEND
222. 17    CP(I,J)=CP(I+1,J)
223. 18    K=K+1
224.      END WHILE
225. 19    IA=NEND
226.      IB=NEND
227.      M=1
```

```
228.      N=NEND
229.      IJOB=0
230.      DO 20 I=1,NEND
231. 20    CE(I)=CI(I)
232. 21    CONTINUE
233.      DO 22 I=1,N
234.      CRI(I)=CI(I)
235.      CRE(I)=CE(I)
236.      DO 22 J=1,N
237. 22    CRP(I,J)=CP(I,J)
238.      CALL LEQT1C(CRP,N,IA,CRE,M,IB,IJOB,WA,IER)
239.      PI=3.14159265D0
240.      BKO=2.0D0*PI*BN
241.      READ,NPDG
242.      DEL=1./NPDG
243.      CJ=(0.0D0,1.0D0)
244.      THETA=-90.0D0
245.      NM=NPHI/2
246.      NM=2*NM
247.      K=1
248.      EMAX=HMAX=0.0D0
249.      WHILE (THETA.LE.90.0D0) DO
250.      SANG=DSIN(PI*THETA/180.0D0)
251.      QANG=DCOS(PI*THETA/180.0D0)
252.      ARG=DABS(BKO*SANG)
253.      NB=NPHI+2
254.      CALL MMBSJN(ARG,NB,B1,IER)
255.      IF (NPHI.EQ.0) THEN
256.      BNPI=-B1(2)
257.      ELSE
258.      BNPI=0.5*(B1(NPHI)-B1(NPHI+2))
259.      END IF
260.      BARG=BNPI+B1(NPHI+2)
261.      CEIX=CIIX=CEIY=CIIY=(0.0D0,0.0D0)
262.      DO 23 I=1,N
263.      CXX=2.0D0*PI*XCR(I)*SANG*CJ
264.      CYY=2.0D0*PI*YCR(I)*SANG*CJ
265.      CEIX=CEIX+CRE(I)*CDEXP(CXX)
266.      CIIX=CIIX+CRI(I)*CDEXP(CXX)
267.      CEIY=CEIY+CRE(I)*CDEXP(CYY)
268. 23    CIIY=CIIY+CRI(I)*CDEXP(CYY)
269.      ABEC=CDABS(BNPI*CEIX)
270.      ABE=CDABS(BNPI*CIIX/CP(1,1))
271.      IF (NM.NE.NPHI) THEN
272.      ABHC=CDABS(BARG*QANG*CEIY)
273.      ABH=CDABS(BARG*QANG*CIIY/CP(1,1))
274.      ELSE
275.      ABHC=CDABS(BNPI*CEIY)
276.      ABH=CDABS(BNPI*CIIY/CP(1,1))
277.      END IF
278.      IF (ABE.GT.EMAX) EMAX=ABE
279.      IF (ABH.GT.HMAX) HMAX=ABH
280.      ANG(K)=THETA
281.      EC(K)=ABEC
282.      E(K)=ABE
283.      HC(K)=ABHC
284.      H(K)=ABH
285.      THETA=THETA+DEL
286.      K=K+1
287.      END WHILE
288.      J=K-1
289.      DO 24 K=1,J
290.      Z0=ANG(K)
291.      Z1=EC(K)/EMAX
```

```
292.      Z2=E(K)/EMAX
293.      Z3=HC(K)/HMAX
294.      Z4=H(K)/HMAX
295.      IF(Z1.LE.0.01) Z1=0.01
296.      IF(Z2.LE.0.01) Z2=0.01
297.      IF(Z3.LE.0.01) Z3=0.01
298.      IF(Z4.LE.0.01) Z4=0.01
299.      Z1=20.* ALOG10(Z1)
300.      Z2=20.* ALOG10(Z2)
301.      Z3=20.* ALOG10(Z3)
302.      Z4=20.* ALOG10(Z4)
303. 24    WRITE(8,*)Z0,Z1,Z2,Z3,Z4
304.      STOP
305.      END
306. C ****
307. C * A SUBROUTINE OF FINDIND THE SELF WALL ADMITTANCE OF A      *
308. C * SINGLE CIRCULAR PATCH.                                *
309. C ****
310.      SUBROUTINE PSLFC(NPHI,BN,HN,EPSR,GIIN,BIIN)
311.      IMPLICIT REAL*8(A-B,D-H,O-Z),INTEGER(I-N)
312.      REAL*8 B1(50),B2(50)
313.      PI=3.14159265D0
314.      HB=HN/BN
315.      BKO=2.0D0*PI*BN
316.      BK=BKO*DSQRT(EPSR)
317.      BEK=BK*DSQRT(1.0D0+2.*HB/(PI*EPSR)*(DLOG(PI/2.0D0/HB)+*
318.      *          1.7726))
319.      JACK=NPHI+1
320.      GIIN=0.0D0
321.      X2=BKO**2
322.      K=-1
323.      M=30
324.      IF (NPHI.EQ.0) THEN
325.      WHILE(M.GE.2)DO
326.      K=-K
327.      B=K*M*(M-1.0D0)/(2.0D0*M-1.0D0)
328.      GIIN=(GIIN+B)*X2/M**2
329.      M=M-1
330.      END WHILE
331.      GIIN=HB*GIIN*X2
332.      CALL BJV(BK,1,B1,B2)
333.      B1L1=B1(1)
334.      B2L1=B2(1)
335.      CALL BJV(BEK,2,B1,B2)
336.      B1L2=B1(2)
337.      B2L2=B2(2)
338.      BUP=-2.0D0*B1L2/(PI*BKO*B1L1**2)
339.      BDI=B2L2-B2L1*B1L2/B1L1
340.      BIIN=BUP/BDI
341.      ELSE
342.      J=1
343.      NF=2*NPHI+1
344.      DO 25 I=1,NF
345. 25    J=J*I
346.      WHILE (M.GE.2) DO
347.      K=-K
348.      TMPN=2.0D0*(M+NPHI)
349.      B=NPHI+M*(M-1.0)/(TMPN-1.0D0)-NPHI**2/(TMPN+1.0D0)
350.      GIIN=(GIIN+K*B)*X2/(M*(M+2.0D0*NPHI))
351.      M=M-1
352.      END WHILE
353.      GIIN=GIIN*X2/J
354.      GIIN=HB*(GIIN-X2*NPHI*(NPHI+3.0D0)/((2.0D0*NPHI+3.0D0)*
355.      *          *J)+NPHI*(NPHI+1.0D0)/J)*X2**NPHI
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356.      CALL BJV(BK,JACK,B1,B2)
357.      B1L0=B1(JACK)
358.      B2L0=B2(JACK)
359.      L1=JACK-1
360.      L2=JACK+1
361.      CALL BJV(BEK,L1,B1,B2)
362.      B1L1=B1(L1)
363.      B2L1=B2(L1)
364.      B1L2=B1(L2)
365.      B2L2=B2(L2)
366.      BUP=-2.0D0*(B1L1-B1L2)/(PI*BKO*B1L0**2)
367.      BDI=B2L1-B2L2-B2L0*(B1L1-B1L2)/B1L0
368.      BIIN=BUP/BDI
369.      END IF
370.      RETURN
371.      END
372. C ****
373. C * A SUBROUTINE OF FINDING THE MUTUAL WALL ADMITTANCE *
374. C * BETWEEN ANY TWO CIRCULAR PATCHES. *
375. C ****
376.      SUBROUTINE PMUTC(NPHI,BN,HN,EPSR,XI,YI,XJ,YJ,GIJN,BIJN)
377.      IMPLICIT REAL*8(A-B,D-H,O-Z),INTEGER(I-N)
378.      PI=3.14159265D0
379.      BKO=2.0D0*PI*BN
380.      DM=DSQRT((XJ-XI)**2+(YJ-YI)**2)
381.      IF (XI.EQ.XJ) THEN
382.      FIM=PI/2.0D0
383.      ELSE
384.      FIM=DATAN((YJ-YI)/(XJ-XI))
385.      END IF
386.      N=72*(1+NPHI)
387.      GIJN=BIJN=0.0D0
388.      FKTR=(HN/BN)*(BKO/N)**2
389.      IF (NPHI.NE.0) FKTR=2.*FKTR
390.      DO 26 I=1,N
391.      FI=(2*I-1)*PI/N
392.      A1=FI-FIM
393.      ELI=DSQRT(DM**2+BN**2+2.*DM*BN*DCOS(A1))
394.      ALFI=DARSIN((BN/ELI)*DSIN(A1))
395.      WI=DCOS(NPHI*FI)
396.      DO 26 J=1,N
397.      FIP=(2*J-1)*PI/N
398.      A2=FIP-FI
399.      A3=FIP-ALFI-FIM
400.      RN=DSQRT(ELI**2+BN**2-2.*ELI*BN*DCOS(A3))
401.      RKO=2.0D0*PI*RN
402.      WJ=DCOS(NPHI*FIP)
403.      BETA=A3+DARSIN(BN/RN*DSIN(A3))
404.      TRS=DSIN(RKO)
405.      TRC=DCOS(RKO)
406.      FG1=(TRS-RKO*TRC)/(RKO**3)
407.      FG2=TRS/RKO
408.      FB1=(TRC+RKO*TRS)/(RKO**3)
409.      FB2=TRC/RKO
410.      TRG1=2.0D0*DCOS(A2)
411.      TRG2=DCOS(BETA)*DCOS(BETA-A2)
412.      FG3=FG1*TRG1+(FG2-3.0D0*FG1)*TRG2
413.      FB3=FB1*TRG1+(FB2-3.0D0*FB1)*TRG2
414.      FG3=WI*WJ*FG3
415.      FB3=WI*WJ*FB3
416.      GIJN=GIJN+FG3
417. 26   BIJN=BIJN+FB3
418.      GIJN=FKTR*GIJN
419.      BIJN=FKTR*BIJN
```

```
420.      RETURN
421.      END
422. C ****BESSEL FUNCTIONS OF THE FIRST AND SECOND KIND ARE ****
423. C * BESSEL FUNCTIONS OF THE FIRST AND SECOND KIND ARE *
424. C * COMPUTED IN THIS SUBROUTINE BJV. *
425. C ****
426.      SUBROUTINE BJV(X,N,JN,YN)
427.      REAL*8 JN(50),YN(50),X,ANORM,PI,SUM,DLOG
428.      INTEGER N,IM,JX,JX1,I,L,AK,K
429.      IM=N+20
430.      JN(IM)=0.0D0
431.      JN(IM-1)=1.0D0
432.      JX=IM-2
433.      JX1=JX+1
434.      DO 27 I=1,JX
435.      L=JX1-I
436.      AK=L-1
437. 27   JN(L)=2.0D0*(AK+1.0)/X*JN(L+1)-JN(L+2)
438.      ANORM=JN(1)
439.      DO 28 L=3,IM,2
440. 28   ANORM=ANORM+2.0D0*JN(L)
441.      DO 29 L=1,IM
442. 29   JN(L)=JN(L)/ANORM
443.      PI=3.14159265D0
444.      SUM=0.0D0
445.      DO 30 L=3,IM,2
446.      K=(L-1)/2
447.      AK=K
448. 30   SUM=SUM+(-1.0)**K*JN(L)/AK
449.      YN(1)=2.0D0/PI*(DLOG(X/2.0D0)+0.577215664)*JN(1)-4./PI*SUM
450.      YN(2)=(JN(2)*YN(1)-2.0D0/(PI*X))/JN(1)
451.      DO 31 L=3,IM
452.      AK=L-1
453. 31   YN(L)=2.0D0*(AK-1.0D0)/X*YN(L-1)-YN(L-2)
454.      RETURN
455.      END
456. $ENTRY
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